1	Experimental study on the electrical conductivity of quartz
2	andesite at high temperature and high pressure: evidence of
3	grain boundary transport
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29	ABSTRACT In this study, the electrical conductivity of quartz andesite was in situ				
30	measured under conditions of 0.5-2.0 GPa and 723-973 K using a YJ-3000t				
31	multi-anvil press and a Solartron-1260 Impedance/Gain-phase Analyzer. Experimental				
32	results indicate that grain interior transport controls the higher frequencies $(10^2 - 10^6)$				
33	Hz), whereas the grain boundary process dominates the lower frequencies $(10^{-1}-10^2)$				
34	Hz). For a given pressure and temperature range, the relationship between $\log \sigma$ and				
35	1/T follows the Arrhenius relation. As temperature increased, both the grain boundary				
36	and grain interior conductivities of quartz andesite increased; however, with				
37	increasing pressure, both the grain boundary and grain interior conductivities of the				
38	sample decreased. By the virtue of the dependence of grain boundary conductivity on				
39	pressure, the activation enthalpy and the activation volume were calculated to be				
40	0.87–0.92 eV and 0.56±0.52 cm^3 /mol, respectively. The small polaron conduction				
41	mechanism for grain interior process and the ion conduction mechanism for grain				
42	boundary process are also discussed.				
43	Keywords electrical conductivity, quartz andesite, grain boundary,				
44	conduction mechanism				
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55 **1. Introduction**

Studies of the electrical conductivity of rocks at high temperatures and high 56 pressures have found that similar to temperature and pressure, grain boundary greatly 57 affecting the electrical properties of rocks. Grain boundary, a general property of 58 rocks, is therefore receiving increasing attention from researchers. Peridotite is the 59 most important rock in the upper mantle, and the influence of grain boundary on its 60 electrical conductivity has been studied in detail (Tyburczy and Roberts, 1990; 61 Roberts and Tyburczy, 1991, 1993, 1994; Xu et al., 1998, 2000; ten Grotenhuis et al., 62 2004; Watson et al., 2010; Wu et al., 2010). However, the relation between the total 63 conductivity, grain boundary and grain interior conductivity for andesite remains 64 unclear till now. 65

Andesite is representative calcium alkaline neutral lava that forms in plate 66 subduction settings, and is thus widely distributed in the orogenic belts bordering the 67 Pacific Ocean. Extensive studies of its electrical conductivity have achieved notable 68 results. Waff and Weill (1975) measured the electrical conductivities of andesite of 69 70 varying components (Na₂O: 4.96–7.83 wt.%; FeO: 4.99–13.7 wt.%) using a direct current (DC) method at room pressure and different oxygen partial pressures of CO_2 71 and H₂. They found that increasing alkali ion content significantly increased the 72 electrical conductivity of andesite whereas oxygen fugacity and iron content had little 73 74 effect (Waff and Weill, 1975). Tyburczy and Waff (1983) employed the alternating current (AC) at pressures of 0–2.55 GPa and temperatures of 1473–1673 K to observe 75 the electrical conductivity of andesite melt from Crater Lake. By combining the 76 electrical conductivity data from andesite melt and tholeiite to model the high 77 78 conductivity zone, they concluded that the electrical conductivity of andesite melts increases with the rise of pressure, and that a minimum melt fraction of 5-10 % can 79 account for the anomalously high electrical conductivity of the upper mantle in typical 80 andesite regions (Tyburczy and Waff, 1983). More recently, Laumonier et al. (2015) 81 82 measured the electrical conductivity of dacitic melts with H₂O contents up to 12 wt.% at pressures of 0.15–3.0 GPa and temperatures of 673–1573 K, and demonstrated that 83

the electrical conductivity is strongly dependent on the water content. Likewise, the influence of pressure on the activation enthalpy is strongly correlated with the sample's water content. By means of T-P-H₂O model, crustal and mantle wedge conductive bodies have been interpreted by the presence of silica-rich, hydrous, partially crystallized magma (Laumonier et al., 2015). However, previous studies mainly focused on the grain interior conductivity of andesite rather than the effect of grain boundary conductivity.

In this study, the grain boundary electrical conductivity of quartz andesite was measured at pressures of 0.5-2.0 GPa and temperatures of 723-973 K within the frequency range of 10^{-1} to 10^{6} Hz. The characteristic parameters of the electrical conductivity of quartz andesite acquired here include the activation enthalpy and the activation volume. These parameters allow discuss of the relationship between the contributions from the grain interior and grain boundary conductivity, and the total conductivity. The conduction mechanism was also discussed.

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99 2. Experimental Procedure

100 **2.1 Sample preparation**

101 Quartz andesite was collected from Shizhu Town, Yongkang City, Zhejiang 102 Province, China. The samples were fresh and unaltered, and appeared pale-yellow in 103 color. According to observation under the optical microscope, the quartz andesite 104 mainly consists of fine-grained plagioclase, amphibole, quartz, and feldspathic matrix, 105 without any accessory mineral.

Before experiment, the samples were cut into cylinders of 6 mm diameter and 6 mm height, and cleaned ultrasonically using deionized water, acetone, and ethanol in turn. Finally, they were placed in an oven at 323 K for 24 h. The chemical composition and mineralogical proportion of the sample (Table 1) were analyzed by X-ray fluorescence spectrometer (XRF) and electron microprobe analysis (EPMA) at the State Key Laboratory of Ore Deposit Geochemistry, Institute of Geochemistry, Chinese Academy of Sciences, Guiyang, China. Table 1

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114 **2.2 High-pressure conductivity cell and impedance measurements**

The electrical conductivity in situ measurements at high pressures and high temperatures were performed in a YJ-3000t multi-anvil apparatus and a Solartron-1260 Impedance/Gain-phase Analyzer at the Key Laboratory of High-Temperature and High-Pressure Study of the Earth's Interior, Institute of Geochemistry, Chinese Academy of Sciences, Guiyang.

The equipment and experimental process are described in detail by Dai et al. 120 (2012) and Hu et al. (2014). A diagram of the cross-section of high-pressure cell 121 assembly is shown in Figure 1. In order to avoid the effect of dehydration on the 122 impedance spectroscopy measurement, a pyrophyllite $(32.5 \times 32.5 \times 32.5 \text{ mm})$ 123 124 pressure-transmitting medium was heated at 1173 K for 12h in a muffle furnace. The heater was composed of three-layer stainless steel sheets (total thickness: 0.5 mm) in 125 the shape of a tube. Similar to previous studies (Dai et al., 2012; Hu et al., 2014), an 126 alumina and magnesia sleeve were used to ensure that the sample was in a relatively 127 insulated environment. A grounded 0.025-mm-thick nickel foil located in the middle 128 of the alumina and magnesia sleeve shielded against external electromagnetic and 129 spurious signal interference. The electrodes were composed of two nickel disks (0.5 130 mm in thickness and 6 mm in diameter). Temperature was monitored using a 131 132 NiCr-NiAl thermocouple in contact with the middle of the sample.

During the experiment, Solartron-1260 Impedance/Gain-phase Analyzer was 133 adopted to collect the impedance spectroscopy with a signal voltage of 1 V and 134 frequency range of 10^{-1} – 10^{6} Hz. To explore the influence of pressure on electrical 135 conductivity, electrical conductivity was conducted in the pressure range of 0.5-2.0 136 GPa. With pressure increased at 1.0 GPa/h to each designated pressure, sample was 137 then heated at 100 K/h, and the complex impedance of quartz andesite was measured 138 at temperature intervals of 50 K. To obtain credible data, the temperature was 139 140 stabilized for several minutes at each step before measurement. Experimental errors in

the temperature and pressure gradients during each measurement were no more than ± 10 K and ± 0.1 GPa, respectively. The obtained impedance spectra were fitted by an equivalent circuit made of a series of R_1 - C_1 and R_2 - C_2 -W (R_1 and C_1 correspond respectively to the resistance and capacitance of grain interior conduction process, and R_2 , C_2 , and W correspond respectively to the resistance, capacitance, and Warburg element of grain boundary conduction process.

Figure 1

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149 **3. Experimental Results**

In this study, the Nyquist and Bode plots, respectively, for the complex 150 impedance of typical quartz andesite were obtained under conditions of 1.0 GPa, 151 723–973 K and 10^{-1} – 10^{6} Hz (Figures 2 and 3). Similar results were also obtained 152 under different pressures. The presence of different relaxation time constant led to the 153 appearance in the Nyquist plot of both a semicircular arc and a 45° slope in the 154 complex impedance plane at the given frequency range. The first semicircle 155 impedance arc $(10^2 - 10^6 \text{ Hz})$ represents the grain interior conduction mechanism; it 156 crosses the origin, and its center lies on the real axis. The 45 ° slope in the complex 157 impedance plane at the end of the first semicircle $(10^{-1}-10^2 \text{ Hz})$ represents grain 158 boundary diffusion. With the rise of temperature, the diameter of impedance arc and 159 value of impedance decreased rapidly; hence, the electrical conductivity increased. 160 The Bode plot (Figure 3) reflects the dependence of modulus (|Z|) and phase angle (θ) 161 on frequency. From high to low frequency, the impedance modulus increased rapidly, 162 and the absolute value of the phase angle tended toward zero. Impedance 163 spectroscopy theory (Nover et al., 1992; Huebner and Dillenburg 1995; Huang et al., 164 2005) relates the real part (Z'), imaginary part(Z''), modulus (|Z|) and phase angle (θ) 165 as follows: $Z' = |Z| \cos \theta$ and $Z'' = |Z| \sin \theta$. According to previous studies (Dai 166 and Karato, 2014a, b), the resistance of the grain interior can be determined by 167 modeling the electrical response with equivalent circuit of resistance and capacitance 168

169 (R_1C_1) . However, the impedance at low frequency is not a wholly semicircular. By 170 combing the 45 ° of slope in the complex impedance at low frequency, a Warburg 171 element was adopted to fit the grain boundary resistance. The equivalent circuit is 172 shown in Figure 2. Another equivalent circuit composed of a resistor and capacitor in 173 parallel was used simultaneously to fit the total resistance. Furthermore, the grain 174 interior, grain boundary and total electrical conductivity are in accordance with the 175 following expression:

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$$\sigma = L/SR \tag{1}$$

where *L* is the sample length (m), *S* is the cross-sectional area of the electrode (m²), and *R* is the resistance for the given conduction process (Ω).

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181 At pressures of 0.5–2.0 GPa and temperatures of 723–973 K, the relationship 182 between the electrical conductivity (σ) of the quartz and esite and reciprocal 183 temperature (1/*T*) was fitted using the Arrhenius relation:

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$$\sigma = \sigma_0 \exp(-\Delta H / kT) \tag{2}$$

185 where σ_0 is the pre-exponential factor (S/m), ΔH is the activation enthalpy (eV), *k* is 186 the Boltzmann constant, and *T* is the absolute temperature (K). The relationship 187 between activation energy ΔU (eV), pressure *P* (GPa) and activation volume ΔV 188 (cm³/mol) is expressed as:

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$$\Delta H = \Delta U + P \times \Delta V \tag{3}$$

The grain interior, grain boundary, and total conductivity at different pressures and temperatures are plotted against reciprocal temperature in Figures 4–6. Figures 4 and 5 show the plots for grain interior and grain boundary conductivity, respectively. The relationship between grain interior, grain boundary and total conductivity at 1.0 GPa is shown in detail in Figure 6, and the value of electrical conductivity under 1.0 GPa is summarized in Table 2. Similar results were obtained under 0.5–2.0 GPa. The ratio of grain boundary (σ_{gb}) to grain interior (σ_{gi}) conductivity represents their respective contributions to total conductivity; it varies with temperature and pressure,
and is plots in the range 0.5–2.0 GPa in Figure 7. Fitting parameters of the grain
interior and grain boundary conductivity are listed in Table 3.

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Figures 4–7 and Tables 2–3

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202 **4. Discussions**

In the present work, the grain interior (σ_{gi}), grain boundary (σ_{gb}) and total electrical conductivity (σ_t) of quartz andesite were in situ measured at the pressures of 0.5–2.0 GPa and temperatures of 723–973 K. With the rise of pressure, the grain boundary conductivity decreases, while the activation enthalpy and pre-exponential factor increase (Figure 4 and Table 3). From Figure 6, it is clear that the grain boundary conductivity was higher than either the grain interior or total conductivity, and the total conductivity was lower than the grain interior conductivity.

The activation energy and activation volume for grain boundary conduction 210 process under the experimental conditions were 0.90 ± 0.10 eV and 0.56 ± 0.52 cm³/mol, 211 respectively. The ratio of grain boundary to grain interior conductivity ($\sigma_{\rm gb}/\sigma_{\rm gi}$) at 212 0.5–2.0 GPa (Figure 7) gradually decreased with increasing temperature and pressure; 213 and thus the contribution of grain boundary conductivity to the total conductivity of 214 215 quartz andesite continually decreased with increasing temperature and pressure. Dai et al. (2008) presented a functional model of the variation of grain boundary 216 conductivity with pressure in which the grain boundary conductivity of peridotite 217 varies with the width of grain boundary, as follows: 218

$$\sigma_{\text{gb-micro}} = \sigma_{\text{gb-bulk}} \left(d \,/\, D \right) \tag{4}$$

where $\sigma_{gb-micro}$ is the microscopic grain boundary conductivity (S/m), $\sigma_{gb-bulk}$ is the bulk grain boundary conductivity (S/m), *d* is the grain boundary width (μ m), and *D* is the grain size (μ m). According to equation (4), the diffusivity of cements between feldspar and amphibole in the quartz andesite increased with the rise of pressure, reducing the grain boundary width along the direction of current transmission and decreasing the grain conductivity accordingly. These results are consistent with those of ten Grotenhuis et al. (2004) and Dai et al. (2008) on the effect of pressure on the grain boundary electrical conductivity of peridotite.

Figure 5 shows that the grain interior conductivity of the quartz and esite 228 decreased with increasing pressure; the activation enthalpy and pre-exponential factor 229 increased accordingly. The variation of grain interior conductivity with pressure 230 observed here is similar to previous studies which were concentrated on the effect of 231 232 partially molten andesite (Waff and Weill, 1975; Tyburczy and Waff, 1983; Laumonier et al., 2015) (Figure 8). The activation enthalpy (0.81–1.05 eV) and activation volume 233 $(4.96\pm0.52 \text{ cm}^3/\text{mol})$ of quartz andesite are within the same range as results for 234 andesite (0.78-1.17 eV and 3.25-17.9 cm³/mol, respectively) from Crater Lake 235 (Tyburczy and Waff, 1983), and are also similar to those of dacitic melts (0.69-1.0 236 eV and 3.9-24.7 cm³/mol, respectively) measured by Laumonier et al. (2015). 237 However, discrepancies in pressure, temperature, melting conditions, and chemical 238 composition of the samples are the important factors that might have led to 1-2 orders 239 240 of magnitudes lower found here, compared with previous studies.

Figure 8

On the other hand, the logarithmic conductivity (Log σ) and reciprocal 241 temperature (T^1) show a strong linear relationship (>99%). On the base of the result, 242 including FeO=5.02 wt.% in the quartz and esite (Table 1), $\Delta H=0.81-1.05$ eV and 243 $\Delta V = 4.96 \pm 0.52$ cm³/mol (Table 3), it imply that there is only one single dominant 244 conduction mechanism for quartz andesite. Numerous studies have reported similar 245 results, indicating that the conduction mechanism is the small polaron (Xu et al., 1998; 246 247 Scarlato et al., 2004; Dai et al., 2008; Yang and Heidelbach, 2011). We consider that the conduction mechanism of grain interior conduction process is the small polaron 248 conduction. The hopping process can be described as follows: 249

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 $Fe^{\times}_{Mg} + h^{\cdot} \rightleftharpoons Fe^{\cdot}_{Mg}$ (5)

In ferromagnesian silicate, the presence of Fe_{Mg} generating an extra positive charge which repulses cations causing the lattice deformation is small polaron (Dai et

al., 2013). The small polaron is an important conduction mechanism at low 253 temperature; it is characterized by the transfer of an electron hole (h) from Fe_{Me} to Fe 254 $_{M_{F}}^{\times}$ (Schmidbauer et al., 2000; Huang et al., 2005; Poe et al., 2008; Dai et al., 2014, 255 2015). In light of the above mentioned results, the low energy barrier for the 256 transmission process resulted in the low activation enthalpy of quartz andesite. Two 257 further factors, oxygen fugacity and iron content, also affect the small polaron 258 conduction of quartz and site. The proportion of ferric iron in the total iron (Fe³⁺/ Σ Fe) 259 increases with increasing oxygen fugacity; with the rise of iron content, the charge 260 carrier concentration also increases. However, the effects of oxygen fugacity and iron 261 content on the grain interior conductivity of quartz andesite require further research. 262

As mentioned above, a Warburg element was adopted to fit the grain boundary 263 resistance; and it indicates that the grain boundary conduction process occurred via 264 ion diffusion. A large quantity of alkali ions are contained in the quartz andesite 265 (Na₂O: 4.98 wt.%; K₂O: 4.16 wt.%), and requires only a low activation energy (Hu et 266 al., 2014). Combing the activation energy of grain boundary conduction process 267 268 (0.87-0.92 eV), we conclude that the grain boundary conduction mechanism for quartz andesite is the ion conduction. However, the effects of alkali iron content on 269 the grain boundary conductivity of quartz andesite require further research. 270

271 **5. Conclusions**

At pressures of 0.5-2.0 GPa and temperatures of 723-973 K, and within the 272 frequency $10^{-1}-10^{6}$ Hz, the grain boundary conductivity of quartz and site ranged 273 from $10^{-4.2}$ – $10^{-2.2}$ S/m; the activation enthalpy and activation volume were 0.87–0.92 274 eV and 0.56±0.52 cm³/mol, respectively. The grain boundary conductivity varied 275 276 greatly with pressure, temperature. Its effect on the total conductivity increased with the rise of temperature. The grain boundary conductivity tended to decrease with 277 increasing pressure. At 0.5-2.0 GPa, the total conductivity of quartz andesite is 278 slightly lower than grain interior conductivity due to the presence of grain boundary. 279 These obtained physical parameters, combined with data on the chemical and 280 mineralogical composition of the andesite, suggest that the conduction mechanism for 281

grain interior of quartz andesite is the small polaron conduction, and for grainboundary is the ion conduction.

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312 **References**

- Dai L.D., Li H.P., Hu H.Y., et al., Experimental study of grain boundary electrical conductivities
 of dry synthetic peridotite under high-temperature, high-pressure, and different oxygen
 fugacity conditions, *J. Geophys. Res.*, 2008, vol. 113, no. B12.
- Dai L.D., Li H.P., Hu H.Y., et al., The effect of chemical composition and oxygen fugacity on the
 electrical conductivity of dry and hydrous garnet at high temperatures and pressures, *Contrib. Mineral. Petrol.*, 2012, vol. 163, no. 4, pp. 689-700.
- Dai L.D., Li H.P., Hu H.Y., et al., Electrical conductivity of Alm₈₂Py₁₅Grs₃ almandine-rich garnet
 determined by impedance spectroscopy at high temperatures and high pressures,
 Tectonophysics, 2013, vol. 608, pp. 1086-1093.
- Dai L.D., Hu H.Y., Li H.P., et al., Influence of temperature, pressure, and chemical composition on
 the electrical conductivity of granite, *Am. Mineral.*, 2014, vol. 99, no. 7, pp. 1420-1428.
- Dai L.D. and Karato S., Influence of FeO and H on the electrical conductivity of olivine, *Phys. Earth Planet. Inter.*, 2014a, vol. 237, pp. 73-79.
- Dai L.D. and Karato S., High and highly anisotropic electrical conductivity of the asthenosphere
 due to hydrogen diffusion in olivine, *Earth Planet Sci. Lett.*, 2014b, vol. 408, pp. 79-86.
- Dai L.D., Li H.P., Hu H.Y., et al., Electrical conductivity of gabbro: the effect of temperature,
 pressure and oxygen fugacity, *Eur. J. Mineral.*, 2015, vol. 27, pp. 215-224.
- Huebner J.S. and Dillenburg R.G., Impedance spectra of hot, dry silicate minerals and
 rock-qualitative interpretation of spectra, *Am. Mineral.*, 1995, vol. 80, nos. 1-2, pp. 46-64.
- Huang X.G., Xu Y.S. and Karato S., Water content in the transition zone from electrical
 conductivity of wadslevite and ringwoodite, *Nature*, 2005, vol. 434, no. 7034, pp. 746-749.
- Huang X.G., Huang X.G. and Bai W.M., Progress of high temperature and high pressure
 experimental study on the electrical conductivity of the minerals and rocks, *Prog. Geophys.*,
 2010, vol. 25, no. 4, pp. 1247-1258. (in Chinese)
- Hu H.Y., Dai L.D., Li H.P., et al., Electrical conductivity of K-feldspar at high temperature and
 high pressure, *Miner. Petrol.*, 2014, vol. 108, no. 5, pp. 609-618.
- 339 Laumonier M., Gaillard F. and Sifre D., The effect of pressure and water concentration on the
- electrical conductivity of dacitic melts: implication for magnetotelluric imaging in subduction
- areas, *Chem. Geol.*, 2015, in press.

- 342 Nover G., Will G. and Waitz R., Pressure-induced phase-transition in Mg₂GeO₄ as determined by
- frequency-dependent complex electrical-resistivity measurements, *Phys. Chem. Minerals*,
 1992, vol. 19, no. 3, pp. 133-139.
- Poe B.T., Romano C., Varchi V., et al., Electrical conductivity of a phonotephrite from Mt.
 Vesuvius: The importance of chemical composition on the electrical conductivity of silicate
 melts, *Chem. Geol.*, 2008, vol. 256, nos. 3-4, pp. 193-202.
- Roberts, J.J. and Tyburczy J.A., Frequency dependent electrical properties of polycrystalline
 olivine compacts, *J. Geophys. Res.*, 1991, vol. 96, no. B10, pp. 16205-16222.
- Roberts J.J. and Tyburczy J.A., Frequency-dependent electrical-properties of dunite as functions of
 temperature and oxygen fugacity, *Phys. Chem. Minerals*, 1993, vol. 19, no. 3, pp. 545-561.
- Roberts J.J. and Tyburczy J.A., Frequency-dependent electrical-properties of minerals and
 partial-melts, *Surv. Geophys*, 1994, vol. 15, no. 2, pp. 239-262.
- Schmidbauer E., Kunzmann T., Fehr T., et al, Electrical resistivity and Fe-57 Mossbauer spectra of
 Fe-bearing calcic amphiboles, *Phys. Chem. Minerals*, 2000, vol. 27, no. 5, pp. 347-356.
- Tyburczy J.A. and Waff H.S., Electrical conductivity of molten basalt and andesite to 25 kilobars
 pressure: Geophysical significance and implications for charge transport and melt structure, *J*.

358 *Geophys. Res.*, 1983, vol. 88, no. B3, pp. 2413.

- Tyburczy, J.A. and Roberts, J.J., Low frequency electrical response of polycrystalline olivine
 compacts: grain boundary transport, *Geophys. Res. Lett.*, 1990, vol. 17, no. 11, pp.
 1985-1988.
- ten Grotenhuis S.M., Drury M.R., Peach C.J., et al., Electrical properties of fine-grained olivine:
 Evidence for grain boundary transport, *J. Geophys. Res.*, 2004, vol. 109, no. B06203.

364 Waff H.S. and Weill D.F., Electrical conductivity of magmatic liquids effects of temperature:

- effects of temperature, oxygen fugacity and composition, *Earth Planet Sci. Lett.*, 1975, vol.
 28, no. 2, pp. 254-260.
- Watson H.C., Roberts J.J. and Tyburczy J.A., Effect of conductive impurities on electrical
 conductivity in polycrystalline olivine, *Geophys. Res. Lett.*, 2010, vol. 37, no. 2, pp. L02302.
- Wu X.P., Zhang B.H., Xu J.S., et al., Electrical conductivity measurements of periclase under high
 pressure and high temperature, *Phys. B*, 2010, vol. 405, no. 1, pp. 53-56.
- 371 Xu Y.S., Poe B.T. Shankland T.J., et al., Electrical conductivity of olivine, wadsleyite, and

- 372 ringwoodite under upper-Mantle conditions, *Science*, 1998, vol. 280, no. 5368, pp.
 373 1415-1418.
- Xu Y.S., Shankland T.J. and Duba A.G., Pressure effect on electrical conductivity of mantle
 olivine, *Phys. Earth Planet. Inter.*, 2000, vol. 118, nos. 1-2, pp. 149-161.
- 376 Yang X.Z. and Heidelbach F., Grain size on the electrical conductivity of clinopyroxene, *Contrib*.
- *Mineral. Petrol.*, 2011, vol. 163, no. 6, pp. 939-947.

- 402 Figure Captions
- 403 Fig.1. Experimental setup of electrical conductivity measurements
- 404 Fig.2. Nyquist plot of the complex impedance of quartz andesite under conditions of
- 405 1.0 GPa and 723–973 K
- 406 Fig.3. Bode plot of dependence of modulus and phase angle on frequency quartz
 407 andesite under conditions of 1.0 GPa and 723–973 K
- 408 Fig.4. The relationship of the logarithmic grain boundary conductivity vs. reciprocal
- temperature under conditions of 0.5–2.0 GPa and 723–973 K
- 410 Fig.5. The relationship of the logarithmic grain interior conductivity vs. reciprocal
- 411 temperature under conditions of 0.5–2.0 GPa and 723–973 K
- Fig.6. The relationship of grain interior, grain boundary and total conductivity under
 conditions of 1.0 GPa and 723–973 K
- 414 Fig.7. Grain boundary/grain interior conductivity (σ_{gb}/σ_{gi}) versus reciprocal
- 415 temperature (1/T) under conditions of 1.0 GPa and 723-973 K. The ratio represents
- the leading role of grain boundary or grain interior conductivity in the conduction
- 417 process.
- 418 Fig.8. A comparison of grain interior conductivity of quartz andesite with previous419 studies









Figure 3





Figure 5







Oridae	XRF for	EPMA for	EPMA for	EPMA for
Oxides	whole rock (wt.%)	Anorthoclase (wt.%)	Albite (wt.%)	Amphibole (wt.%)
SiO ₂	66.47	59.56	65.25	52.47
Al_2O_3	13.57	26.20	19.75	3.61
MgO	0.44	0.15	0.03	15.72
CaO	1.12	2.28	0.62	12.18
Na ₂ O	4.98	6.99	10.55	0.20
K ₂ O	4.16	3.09	0.02	0.18
FeO	5.02	0.90	1.03	11.23
TiO ₂	0.22	0.03	0.01	0.33
Cr_2O_3	0.02	0.02	1.47	0.03
MnO ₂	/	0.04	0.00	0.12
P_2O_5	0.81	/	/	/
L.O.I	2.74	/	/	/
Total	99.55	99.26	98.73	96.07

Table 1. Whole rock analysis by X-ray fluorescence (XRF) and chemical composition of dominant minerals for quartz andesite by the electron microprobe (EPMA)

Table 2. The value of grain interior, grain boundary and total electrical conductivity under 1.0 GPa and 723–973 K. The estimate error for grain interior conductivity is lower than 5%, for grain boundary conductivity is lower than 7%, and for total conductivity is lower than 5%.

	$L_{OG} \sigma \cdot (S/m)$	$\log \sigma \cdot (S/m)$	$I \circ \sigma \sigma (S/m)$
	$LOg O_{gi} (S/III)$	$LOg O_{gb} (S/III)$	$Log O_t (S/III)$
723 K	-3.62	-3.88	-3.65
773 K	-3.21	-3.52	-3.25
823 K	-2.83	-3.22	-2.87
873 K	-2.50	-2.83	-2.539
923 K	-2.21	-2.61	-2.26
973 K	-1.95	-2.29	-1.99

quartz andesite							
	P (GPa)	<i>T</i> (K)	Log σ_0 (S/m)	$\Delta H ({ m eV})$	2	$\Delta U ({ m eV})$	$\Delta V (\text{cm}^3/\text{mol})$
	0.5	723-973	2.36±0.13	0.87±0.02	99.67		
Grain	1.0	723-973	2.24±0.13	0.88±0.02	99.70	0.00.010	0.56±0.52
boundary	1.5	723-973	2.30±0.11	0.92±0.02	99.78	0.90±0.10	
	2.0	723-973	1.94±0.12	0.89±0.02	99.71		
	0.5	723-973	1.63±0.03	0.81±0.01	99.97		
Grain	1.0	723-973	1.90±0.02	0.93±0.01	99.99	0.76±0.06 4.9	4.06 10.52
interior	1.5	723-973	2.41 ±0.01	1.05±0.02	99.81		4.90±0.52
	2.0	723-973	2.03±0.01	1.02±0.03	99.70		

Table 3. Fitted parameters for the grain interior and grain boundary conductivity of