Moment magnitude estimates for Central Anatolian earthquakes using coda waves
 Tuna Eken<sup>1</sup>

3

4

<sup>1</sup>Department of Geophysical Engineering, the Faculty of Mines, Istanbul Technical

5 University, 34469 Maslak, Sariyer, Istanbul, Turkey

6

7 *Abstract* 

Proper estimate of moment magnitude that is a physical measure of the energy released at 8 9 earthquake source is essential for better seismic hazard assessments in tectonically active regions. Here a coda wave modeling approach that enables the source displacement spectrum 10 modeling of examined event was used to estimate moment magnitude of central Anatolia 11 earthquakes. To achieve this aim, three component waveforms of local earthquakes with 12 13 magnitudes 2.0  $\leq$  M<sub>L</sub>  $\leq$  5.2 recorded at 69 seismic stations which have been operated between 2013 and 2015 within the framework of the CD-CAT passive seismic experiment 14 15 were utilized. An inversion on the coda wave traces of each selected single event in our database was performed in five different frequency bands between 0.75 and 12 Hz. Our 16 17 resultant moment magnitudes (M<sub>W</sub>-coda) exhibit a good agreement with routinely reported 18 local magnitude (M<sub>L</sub>) estimates for the study area. Apparent move-out that is, particularly, significant around the scattered variation of M<sub>I</sub>-M<sub>w</sub>-coda data points for small earthquakes 19  $(M_L < 3.5)$  can be explained by possible biases of wrong assumptions to account for anelastic 20 21 attenuation and of seismic recordings with finite sampling interval. Finally, we present an empirical relation between M<sub>W</sub>-coda and M<sub>L</sub> for central Anatolian earthquakes. 22

23

24 Keyword(s): Coda waves modelling, seismic moment, moment magnitude, Radiative Transfer25 Theory

### 26 *1. Introduction*

The robust and stable knowledge of source properties (e.g. moment magnitude estimates) is crucial in seismically active countries such as Turkey for a better evaluation of seismic hazard potential as this highly depends on establishment of reliable seismicity catalogs. Moreover, accurate information on source parameters could be important when developing regional attenuation properties.

32

Conventional type of magnitude scales (M<sub>L</sub>, m<sub>b</sub>, M<sub>S</sub>) as the result of empirically derived using 33 34 direct wave analyses can be biased due to various effects such as source radiation pattern, 35 directivity, and heterogeneities along the path since they may cause drastic changes in direct wave amplitude measurements (e.g., Favreau and Archuleta, 2003). Instead several early 36 studies depending on the analysis of local and/or regional coda envelopes have indicated that 37 38 coda wave amplitudes are significantly less variable by a factor of 3-to-5 compared to direct 39 wave amplitudes (e.g., Mayeda and Walter, 1996; Mayeda et al., 2003; Eken et al., 2004; 40 Malagnini et al., 2004; Gök et al., 2016). In fact local or regional coda waves that are usually 41 considered to be generally composed of scattered waves. These wave trains can be simply explained by the single scattering model of Aki (1969) which have been proven to be 42 virtually insensitive to any source radiation pattern effect in contrast to direct waves due to 43 44 the volume averaging property of the coda waves sampling the entire focal sphere (e.g., Aki and Chouet, 1975; Rautian and Khalturin, 1978). In Sato and Fehler (1998) and Sato et al. 45 (2012) an extensive review study on the theoretical background of coda generation and 46 47 advances of empirical observations and modelling efforts can be found in details.

48

There have been several approaches used for extracting information on earthquake source sizevia coda wave analyses. These approaches can be mainly divided into two groups. The first

group of studies can be considered as the parametric approach and essentially employs coda 51 52 normalization strategy in which measurements require a correction for empirically derived quality factors representing seismic attenuation parameters (e.g. intrinsic and scattering). In 53 54 this case, adjustment of final source properties are achieved with the help of some reference events whose seismic moments are previously estimated based on waveform inversion 55 methods. For forward generation of synthetic coda envelopes, either single-backscattering or 56 57 more advanced multiple-backscattering approximation are used. An example to this group is an empirical method originally developed by Mayeda et al. (2003) to investigate seismic 58 source parameters such as energy, moment, and apparent stress drop in the western United 59 60 States and in Middle East. They corrected observed coda envelopes for various influences, for instance, path effect, S-to-coda transfer function, site effect, and any distance-dependent 61 changes in coda envelope shape. Empirical coda envelope method have been successfully 62 63 applied to different regions with complicated tectonics such as northern Italy (e.g. Morasca et al., 2008), Turkey and Middle East (e.g. Eken et al., 2004; Gök et al. 2016); or Korean 64 65 Peninsula (e.g. Yoo et al., 2013).

66

Second type of approach depends on estimating source and structural properties through a 67 joint inversion technique. This technique employs a simultaneous optimization of source, 68 69 path, and site specific terms via a fitting procedure between physically derived synthetic coda 70 envelope and observed coda envelope within a selected time window that includes both the observed coda and direct-S wave parts. Although the conventional coda-normalization 71 72 method essentially relies on the correction for undesired effects of the source and site amplifications, it may fail for small events with a shorter coda. This mainly stems from 73 74 random seismic noise that dominates the coda, which does not satisfy the requirement of homogeneous distribution of energy in space. In the present study, we avoid this shortcoming 75

by involving source excitation and site amplification terms in the inversion process. To 76 achieve this, the Radiative Transfer Theory (RTT) is employed for analytic expression of 77 synthetic coda wave envelopes. The method was originally developed by Sens-Schönfelder 78 and Wegler (2006) and successfully tested on local and regional earthquakes ( $4 \le Ml \le 6$ ) 79 detected by the German Regional Seismic Network. Further it has been applied to investigate 80 source and frequency dependent attenuation properties of different geological settings, i.e., 81 Upper Rhine Graben and Molasse Basin regions in Germany and western Bohemia/Vogtland 82 in Czechia (Eulenfeld and Wegler, 2016); entire United States (2017); central and western 83 North Anatolian Fault Zone (Gaebler et al., 2018; Izgi et al., 2018). A more realistic earth 84 model in which anisotropic scattering conditions were earlier considered by Gusev and 85 Abubakirov (1987) yielded peak broadening effects of the direct seismic wave arrivals. This 86 approach that examines the propagation of P-wave elastic energy and the effect of conversion 87 88 between P- and S-wave energies was later used in Zeng and Aki (1991), Przybilla and Korn 89 (2008), Gaebler et al. (2015).

90

91 In the current work I present source spectra as the output of a joint inversion of S- and coda waves parts extracted from 487 local earthquakes with magnitudes 2.0 < ML < 4.5 detected in 92 central Anatolia. The approach used here employs isotropic acoustic RTT approach for 93 94 forward calculation of synthetic coda envelopes. Gaebler et al. (2015) have observed that modeling results from isotropic scattering were almost comparable with those inferred from 95 relatively more complex elastic RTT simulations with anisotropic scattering conditions. The 96 97 use of a joint inversion technique is advantageous since it is insensitive to any potential bias, which could be introduced by external information, i.e., source properties of a reference that 98 99 is obtained separately from other methods for calibration. This is mainly because of the fact that we utilize an analytical expression of physical model involving source, and path related 100

parameters to describe the scattering process. Moreover the type of optimization during joint
inversion enables the estimates for source parameters of relatively small sized events
compared to the one used in coda-normalization methods.

104

105 *2. Regional Setting* 

106 Present tectonic setting of Anatolia and surrounding regions have been mainly the outcome of the northward converging movements among Africa, Arab, and Eurasian plates. To the west, 107 108 the subducting African plate with a slab roll-back dynamics beneath Anatolia along Hellenic 109 Trench has led to back-arc extension in the Aegean and western Anatolia, while 110 compressional deformation to the east around the Bitlis-Zagros suture was explained by 111 collisional tectonics (e.g. Taymaz et al., 1990; Bozkurt, 2001) (Fig. 1). Central Anatolia is 112 located between an extensional regime to the west due to the subduction, and a compressional 113 regime to the east due to the collisional tectonics. There are several fault systems responsible 114 for ongoing seismic activity in the region. The major fault zone, the Central Anatolian Fault 115 Zone (CAFZ) (Fig. 2), which primarily represents a transtensional fault structure with a small 116 amount of left-lateral offset during the Miocene (e.g. Koçyiğit and Beyhan, 1998), can be 117 considered as a boundary between the carbonate nappes of the Anatolide-Tauride block and the highly deformed and metamorphosed rocks in the Kırşehir block. To the northwest of the 118 119 CAFZ, Tuz Gölü Fault Zone (TGFZ) (Fig. 2), which is characterized by a right-lateral strike 120 slip motion with a significant oblique-slip normal component, appears to be collocated with 121 the Tuz Gölü Basin sedimentary deposits as well as the crystalline rocks within the Kırşehir Block (e.g. Çemen et al., 1999; Bozkurt et al., 2001; Taymaz et al., 2004; Çubuk et al., 2014). 122 At the southwest tip of the study region, the EAFZ generates large seismic activity that can be 123 124 identified by rather complicated seismotectonic setting: predominantly left-lateral strike-slip motion that is well correlated with the regional deformation pattern and with existing local 125

clusters of thrust and normal faulting events on NS- and EW-trending subsidiary faults, 126 127 respectively (Bulut et al., 2012). Such complicated behavior explains kinematic models (e.g. Riedel shear, anti-Riedel shear models) of the shear deformation zone evolution (Tchalenko, 128 129 1970). It connects to the NAFZ at the Karliova Triple Junction (Bozkurt, 2001) and to the 130 south splits into various segments nearby the Adana Basin (Kaymakci et al., 2006) (Fig. 2). Toward the south, the EAFZ reaches the Dead Sea Fault Zone (DSFZ) that has a key role in 131 accommodating northward relative motions of Arabian and African Plates with respect to 132 133 Eurasia.

134

135 *3. Data* 

The present work utilizes three-component waveforms of local seismic activity detected at 72 136 broadband seismic stations (Fig. 2) that have been operated for 2 years between 2013 and 137 138 2015 within the framework of a temporary passive seismic experiment, the Continental 139 Dynamics-Central Anatolian Tectonics (CD-CAT) (Portner et al., 2018). We benefit from 140 revisited standard earthquake catalogue information that is routinely released by the Kandilli 141 Observatory and Earthquake Research Institute (KOERI) (publicly available at http://www.koeri.boun.edu.tr) to extract waveform data for a total of 2231 examined events 142 with station-event pair distance less than 120 km and focal depths less than 10 km. Most of 143 144 the detected seismic activity in the study area is associated to several fault zones in the region, i.e., the EAFZ, CAFZ, DSFZ, TGFZ, etc. Here we note that the use of only local earthquakes 145 is to exclude possible biases, which may be introduced by Moho boundary guided Sn-waves. 146 147 Upper crustal earthquakes with less than 10 km focal depths are preferred in this study to exclude effect of relatively large-scale heterogeneities on coda wave trains. Additionally, we 148 149 performed a visual inspection over all waveforms to ensure high-quality waveforms. Our final

event number reduced to 1193. Selected station and event distributions can be seen in Figure

151

2.

152

Observed waveforms were prepared at 5 different frequency bands with central frequencies at 153 154 0.75, 1.5, 3.0, 6.0, 12.0 Hz via a Butterworth band-pass filtering process. In the next step, we applied Hilbert transform to filtered waveform data in order to obtain the total energy 155 envelopes. An average crustal velocity model was used to predict P and S wave onsets on 156 157 envelopes and then based on this information: (i) the noise level prior to the P-wave onset was 158 eliminated (ii) S-wave window was determined starting at 3s prior to and 7 s afterwards S-159 wave onset as this allowed to include all direct S-wave energy, (iii) starting at the end of the 160 S-wave window, a coda window of 100s at maximum was determined. Length of coda windows can be shorter when signal-to-noise ratio (SNR) is less than 2.5 or when there are 161 162 coda waves from two earthquakes (e.g. because of an aftershock sequence) within the same 163 analysis window, which can cause another rise instead of a decline in the envelope. We omit 164 the earthquakes with less than 10 s of coda length from our database. Taking into account of 165 these criteria, finally coda waveforms extracted from 6541 source-receiver pairs were used for 166 further data process.

167

### 168 *4. Method*

We adopted an inversion procedure that was originally developed by Sens-Schönfelder and Wegler (2006) and later modified by Eulenfeld and Wegler (2016). The forward part, which involves calculation of energy density for a specific frequency band under assumption of an isotropic source, is expressed in Sens-Schönfelder and Wegler (2006) as follows:

173 
$$E_{mod}(t,r) = WR(r)G(t,r,g)e^{-bt}$$
(1)

where W gives source term and it is frequency dependent. R(r) indicates the energy site amplification factor and b is intrinsic attenuation parameter. G(t, r, g) represents Green's function that includes scattered wave field as well as direct wave and its expression is given by Paasschens (1997) as follows:

$$179 \quad G(t,r,g) = e^{(-v_0 t g_0)} \left[ \frac{\delta(r - v_0 t)}{4\pi r^2} + \left( \frac{4\pi v_0}{3g_0} \right)^{-\frac{3}{2}} t^{-\frac{3}{2}} \times \left( 1 - \frac{r^2}{v_0^2 t^2} \right)^{\frac{1}{8}} K \left( v_0 t g_0 \left( 1 - \frac{r^2}{v_0^2 t^2} \right)^{\frac{3}{4}} \right) H(v_0 t - r) \right]$$

$$180 \quad \frac{r^2}{v_0^2 t^2} \Big)^{\frac{3}{4}} H(v_0 t - r) \left[ \qquad (2) \right]$$

Here the term within Dirac delta function represents direct wave and other term indicates
scattered waves. *v*<sub>0</sub> describes the mean S-wave velocity while *g*<sub>0</sub> is the scattering coefficient.

Possible discrepancy between predicted (Eq. 1) and observed energy densities for each event
at each station with N<sub>ij</sub> time samples (index k) in a specific frequency band can be minimized
using:

186

187 
$$\epsilon(g) = \sum_{i,j,k}^{N_S, N_E, N_{ij}} \left( ln E_{ijk}^{obs} - ln E_{ijk}^{mod}(g) \right)^2 \quad (3)$$

188

Here, the number of stations (index i) and events (index j) are shown by  $N_s$  and  $N_E$ , respectively. Optimization of g will be achieved by fulfilling following equality:

191

192 
$$lnE_{ijk}^{obs} = lnE_{ijk}^{mod}$$
 (4) or

193

194 
$$lnE_{ijk}^{obs} = \ln G(t_{ijk}, r_{ijk}, g) + lnR_i + lnW_j - bt_{ijk}$$
(5)

195 Equation 5 simply define an overdetermined inversion problem with  $\sum_{i,j} N_{ij}$  number equation

196 systems and with  $N_{\rm S} + N_{\rm E} + 1$  variables and thus *b*,  $R_i$ , and  $W_j$  can be solved via a least-197 squares technique.  $\epsilon(g)$  can be defined as sum over the squared residuals of the solution. As 198 can be seen from equation 1 that there is an obvious trade-off between  $R_i$  and  $W_j$ , which we 199 can manage by fixing the geometrical mean of  $R_i$  to 1 ( $\Pi R_i = 1$ ). Equation 1 also implies 200 rather moderate trade-off between  $W_j$  and *b*. Trade-off between g and other inverted 201 parameters are usually small since this parameter is fixed through the energy ratio of the 202 direct S-wave and the level of the coda-waves (Gaebler et al., 2018).

203 Eulenfeld and Wegler (2016) present a simple recipe to perform the inversion:

(i) Calculate Green's functions through the analytic approximation of the solution for 3-D isotropic radiative transfer (e.g. Paasschens 1997; Sens-Schönfelder and Wegler, 2006) by using fixed scattering parameters and minimize equation 5 to solve for b,  $R_i$ , and  $W_j$  via a weighted least-squares approach.

208 (ii) Calculate  $\epsilon(g)$  using equation 3.

209 (iii) Repeat (i) and (ii) by selecting different g to find the optimal parameters g, b,  $R_i$  and  $W_j$ 210 that finally minimize the error function  $\epsilon$ .

In Fig. 3 an example for the minimization process that was applied at five different frequencybands is displayed for one selected event at recorded stations of the CD-CAT project.

213 Minimization described above for different frequencies will yield unknown spectral source 214 energy term,  $W_j$  as well as site response,  $R_i$  and attenuation parameters, b, and g that will 215 satisfy optimal fitting between observed and predicted coda wave envelopes. Example for this 216 fitting can be seen in Figure 4. The present study deals with frequency dependency of  $W_j$ 217 since this information can be later useful to obtain source displacement spectrum and thus seismic moment and moment magnitudes of analyzed earthquakes using the formula of the *S*wave source displacement spectrum for a double-couple source in the far-field, which is given
by Sato et al. (2012):

221 
$$\omega M(f) = \sqrt{\frac{5\rho_0 v_0^5 W}{2\pi f^2}} \quad (6)$$

where *W* indicates the radiated S-wave energy at a center frequency *f* while  $v_0$  and  $\rho_0$ represent the mean S-wave speed and medium density, respectively.

The relation between the obtained source displacement spectrum and seismic moment valuewas earlier described in Abercrombie (1995) by:

226 
$$\omega M(f) = M_0 \left(1 + \left(\frac{f}{f_c}\right)^{\gamma n}\right)^{-\frac{1}{\gamma}} \quad (7)$$

where n is related to the high-frequency fall-off and  $\gamma$  is known as shape parameter that controls the sharpness of spectrum at corner frequency between the constant level M<sub>0</sub> (low frequency part) and the fall-off with f<sup>-n</sup> (high frequency part). Taking the logarithm of equation 7 gives:

231

232 
$$\ln \omega M(f) = \ln M_0 - \frac{1}{\gamma} \ln \left( 1 + \left( \frac{f}{f_c} \right)^{\gamma c} \right)$$
(8)

233

Eq. 8 describes an optimization problem where the observed source displacement spectrum data (left-hand side) can be inverted for four unknown source parameters,  $M_0$ ,  $\gamma$ , n, and  $f_c$ (right-hand side) in a simultaneous least-squares inversion scheme. Finally moment magnitude, M<sub>w</sub> can be calculated from modeled source parameters, seismic moment, M<sub>0</sub> using a formula given by Hanks and Kanamori (1979):

$$M_w = \frac{2}{2} \log_{10} M_0 - 10.7 \quad (9)$$

241

242 *5. Results and Discussions* 

243 5.1 Coda wave source spectra

Figure 5 displays observed values of source spectra established by inserting inverted spectral source energy term W at each frequency in Eq. 6 for all analyzed events. Each curve in this figure represents the model spectrum estimate based on the inversion procedure described in the previous section. Modeled spectrum characteristics computed for 487 local earthquakes whose geographical distribution is presented in Figure 2 suggest, in general, that we were able to obtain typically expected source displacement spectrum with a flat region around the low frequency limit and a decaying behaviour above a corner frequency.

251

Owing to the multiple-scattering process within small scale heterogeneities that makes coda waves gain an averaging nature, the variation in coda amplitudes due to differences in source radiation pattern and path effect are reduced (Walter et al., 1995; Mayeda et al., 2003). Eulenfeld and Wegler (2016) found that radiation pattern would have only a minor influence on the S-wave coda while it might disturb attenuation models inferred from the direct S-wave analyses unless the station distribution relative to the earthquakes indicates a good azimuthal coverage.

259

260 Conventional approaches (e.g. Abercrombie, 1995; Kwiatek et al., 2011) to estimate source 261 parameters such as corner frequency, seismic moment, high-frequency fall-off through fitting 262 of observed displacement spectra observed at a given station in an inversion scheme could be 263 misleading since these methods usually: (i) assume a constant value of attenuation effect (no 264 frequency variation) defined by a factor exp ( $-\pi ftQ^{-1}$ ) over the spectrum, (ii) and assume

omega-square model with a constant high-frequency fall-off parameter, n=2. Following Sens-265 266 Schönfelder and Wegler (2006) and Eulenfeld and Wegler (2016), however, we estimate attenuation parameters (intrinsic and scattering) seperately within a simultaneous inversion 267 268 procedure in which high-frequency fall-off parameter varies. This is fairly consistent with 269 early studies (e.g. Ambeh and Fairhead, 1991; Eulenfeld and Wegler, 2016) where significant 270 deviations from the omega square model (n>3) were reported implying that the omega-square 271 model as a source model for small earthquakes must be reconsidered in its general 272 acceptance. Earlier it has been well-observed that the source spectra, especially, for large 273 earthquakes could be better explained by models of two corner frequencies (e.g., 274 Papageorgiou and Aki, 1983; Joyner, 1984; Atkinson, 1990). Recently, Denolle et al. (2016) 275 observed that conventional spectral model of a single-corner frequency and high-frequency fall-off rate could not explain P wave source spectra of thrust earthquakes with magnitude 276 277 Mw 5.5 and above. Instead, they suggested the double-corner-frequency model for large 278 global thrust earthquakes with a lower corner frequency related to source duration and with an 279 upper corner frequency suggesting a shorter time scale unrelated to source duration, which 280 exhibits its own scaling relation. Uchide and Imanishi (2016) reported similar differences 281 from the omega-square model would be valid also for smaller earthquakes by using spectral 282 ratio technique that involves empirical Green's function (EGF) events to avoid having a complete knowledge of path and site effects for shallow target earthquakes ( $M_W$  3.2–4.0) in 283 284 Japan. The source spectra for many of the target events in their study suggested a remarkable 285 discrepancy from the omega-square model for relatively small earthquakes. They explained 286 such differences by incoherent rupture due to heterogeneities in fault properties and applied stress, the double-corner-frequency model, and possibility of a high-frequency falloff 287 288 exponent value slightly higher than 2. In our case, the smallest event was with M<sub>W</sub>-coda 289 larger than 2.0, thus we had no chance to make a similar compared to that of Eulenfeld and

Wegler (2016). However, high-frequency fall-off parameters varied from n=0.5 to n=4. A 290 notable observation in the distribution of n was n=2 or n=2.5 would be better explained for 291 292 earthquakes with  $M_W$ -coda >4.0 whereas the smaller magnitudes exhibited more scattered pattern of variation in n (Figure 7). Eulenfeld and Wegler (2016) claimed that the use of 293 separate estimates of the attenuation or correction for path effect via emprically determined 294 295 Green's function would be better strategy in order to invert station displacement spectra for 296 source parameters. This is mainly because smaller earthquakes (with n>2), in particular, 297 assuming omega-square model can distort the estimates of corner frequency and even seismic 298 moment especially in regions where Q is strongly frequency dependent. Thus, independent estimates of Q during station displacement spectra inversions for source parameters must be 299 300 taken into account or the influence of path such as attenuation must be removed via 301 empirically determined Green's functions (Eulenfeld and Wegler, 2016).

302

# 303 5.2 Coda wave-derived magnitude vs. $M_L$ catalogue magnitude

A scatter plot between catalogue magnitudes based on local magnitudes ( $M_L$ ) and our codaderived magnitudes ( $M_W$ -coda) that are inferred from resultant frequency dependent source displacement spectra and thus seismic moment (e.g. Eq. 9) is shown in Fig. 6. Such comparison suggests an overall coherency between both types of magnitudes. This implies that a very simple model of a first-order approximation for S-wave scattering with isotropic acoustic radiative transfer approach can be efficient to link the amplitude and decaying character of coda wave envelopes to the seismic moment of the source.

311

312 In the present study, a linear regression analyses performed between  $M_W$ -coda and  $M_L$ 313 magnitudes (Fig. 5) resulted in an empirical formula that can be employed to convert local 314 magnitudes into coda-derived moment magnitude calculation of local earthquakes in this 315 region:

316

317 
$$M_{W-coda} = 1.1655 \pm 0.0337 \times M_L - 0.7085 \pm 0.0128$$
(10)

318

Bakun and Lindh (1977) empirically described the linear log seismic moment-local 319 320 magnitude relation between seismic moments (Mo) and local magnitudes (ML) for earthquakes near Oroville, California. Beside this several other studies investigated to find an 321 optimum relation between M<sub>W</sub> and M<sub>L</sub> by implementing linear and/or non-linear curve-fitting 322 323 approaches. Malagnini and Munafò (2018) proposed two different linear fits separated by a 324 crossover M<sub>L</sub>=4.31 could represent M<sub>L</sub>-M<sub>W</sub> data points obtained from earthquakes of the 325 central and northern Apennines, Italy. Several coefficient of regression analyses in their fits 326 account for the combined effects of source scaling and crustal attenuation as well as regional 327 attenuation, focal depth, and rigidity at source. Goertz-Allmann et al. (2011), for instance, 328 introduced hybrid type of scaling relation that is linear below M<sub>L</sub> 2 and above M<sub>L</sub> 4 and a quadratic relation in between  $(2 \le M_L \le 4)$  for earthquakes in Switzerland detected between 329 330 1998 and 2009. Edwards and Rietbrock (2009) employed a second-order polynomial equation 331 to relate local magnitudes routinely reported in the Japan Meteorological Agency (JMA) magnitude and moment magnitude. More recently, using multiple spectral ratio analyses 332 Uchide and Imanishi (2018) estimated relative moment magnitudes for the Fukushima 333 334 Hamadori and the northern Ibaraki prefecture areas of Japan and reported a quadratic form of correlation between JMA magnitudes and moment magnitudes. Resultant empirical curve in 335 Uchide and Imanishi (2018) implied a considerable discrepancy between the moment 336 magnitudes and the JMA magnitudes, with a slope of 1/2 for microearthquakes suggesting 337

possible biases introduced by anelastic attenuation and the recording by a finite samplinginterval.

340

Apparent move-out in Fig. 5 and Eq. 10, presumably stems from the use of different 341 magnitude scales for comparison. Conventional magnitudes scales such as M<sub>L</sub>, mb inferred 342 from phase amplitude measurements are seemingly sensitive to attenuation and 2D variation 343 344 along the path (Pasyanos et al., 2016). Unlike local magnitude scales, seismic moment-based 345 moment magnitude (M<sub>W</sub>) essentially represents a direct measure of the strength of an earthquake caused by fault slip and is estimated from relatively flat portion of source spectra 346 347 at lower frequencies that can be less sensitive to the near surface attenuation effects. The consistency between coda-derived moment magnitude and local magnitude scales for the 348 earthquakes with  $M_W$ -coda > 3.0 indicates that our non-empirical approach successfully 349 350 worked in this tectonically complex region. This observation is anticipated, for relatively 351 large earthquakes, since more energy will be characteristic at lower frequencies. We 352 observed similar type of consistency in early studies that investigate source properties of local 353 and regional earthquakes based on emprical coda methods with simple 1-D radially symmetric path correction (e.g. Eken et al., 2004; Gök et al., 2016). Coda waves-derived 354 source parameters were obtained with high-precision in Mayeda et al. (2005), Phillips et al. 355 356 (2014), Pasyanos et al. (2016) following the use of 2-D path-corrected station techniques to 357 consider the amplitude-distance relationships. Observable outliers in Figure 5, for the events 358 with less than Mw 3.5, however, can be attributed to the either possible biases on local magnitude values taken from the catalogue or small biases on our intrinsic  $(Q_i^{-1})$  and 359 scattering  $(Q_s^{-1})$  attenuation terms. One another possible contribution to such mismatch might 360 be associated to the influences of mode conversions between body and surface waves or 361

362 surface-to-surface wave scattering that are not restricted to low frequencies (<1Hz) (Sens-</li>
363 Schönfelder and Wegler, 2006).

364

### 365 *6.* Conclusions

366 This study provides moment magnitude estimates as a direct physical measure of the seismic 367 energy for local earthquakes with magnitudes  $2.0 \le M_L \le 5.2$  recorded at 69 seismic stations 368 in central Anatolia. The source displacement spectra were obtained following the application 369 of a coda wave modeling procedure that employs a simultaneous optimization of source, path, 370 and site specific terms by fitting physically derived synthetic coda envelope and observed 371 coda envelopes. The Radiative Transfer Theory was used for analytic expression of synthetic 372 coda wave envelopes. Overall consistency between Mw-coda and ML suggests that our non-373 empirical approach successfully worked in this tectonically complex region. Variation of 374 high-frequency fall-off parameter indicated that for smaller earthquakes (n>2) assuming 375 omega-square model can distort the estimates of corner frequency and even seismic moment 376 especially in regions where Q is strongly frequency dependent. Since the present study mainly focuses on source properties of local earthquakes in the study area, scattering and intrinsic 377 378 attenuation properties that are other products of our coda envelope fitting procedure will be 379 examined in details within a future work. Finally, a linear regression analysis resulted in an 380 empirical relation developed between Mw-coda and ML, which will be a useful tool in the 381 future to quickly convert catalogue magnitudes into moment magnitudes for local earthquakes 382 in the study area.

383

384 *Data and resources* 

385 The python code used for carrying out the inverse modeling is available under the permissive386 MIT license and is distributed at https://github.com/trichter/qopen. We are grateful to the IRIS

387 Data Management Center for maintaining, archiving and making the continuous broadband
388 data used in this study open to the international scientific community. The KOERI is specially
389 thanked for providing publicly open local seismicity catalogues.

390

391 *Acknowledgement* 

392 The facilities of IRIS Data Services, and specifically the IRIS Data Management Center, were 393 used for access to waveforms, related metadata, and/or derived products used in this study. 394 IRIS Data Services are funded through the Seismological Facilities for the Advancement of 395 Geoscience and EarthScope (SAGE) Proposal of the National Science Foundation under 396 Cooperative Agreement EAR-1261681. Data for the **CD-CAT** experiment (https://doi.org/10.7914/SN/YB 2013) are available from the IRIS Data Management Center 397 at http://www.iris.edu/hg/. Tuna Eken acknowledge financial support from Alexander von 398 399 Humboldt Foundation (AvH) towards computational and peripherals resources. I am grateful 400 to the Topical Editor Charlotte Krawczyk for handling the revision process and Takahiko 401 Uchide and Ludovic Margerin for their valuable opinions on the improvement of manuscript.

## 402 *References*

403 Abercrombie, R.E.: Earthquake source scaling relationships from -1 to 5 ML using
404 seismograms recorded at 2.5-km depth, J. geophys. Res., 100(B12), 24 015–24 036,
405 1995.

406 Aki, K., and Chouet., B.: Origin of coda waves: Source, attenuation, and scattering effects, J.
407 Geophys. Res. 80, 3322–3342, 1975.

408 Atkinson, G. M.: A comparison of eastern North American ground motion observations with
409 theoretical predictions, Seismol. Res. Lett. 61, 171–180, 1990.

410 Bakun, W.H. and Lindh, A.G.: Local Magnitudes, Seismic Moments, and Coda Durations for

411 Earthquakes Near Oroville, California, Bulletin of the
412 SeismologicalSocietyofAmerica.Vol.67,No.3, pp. 615-629, 1977.

413 Bozkurt, E.: Neotectonics of Turkey—A synthesis: Geodinamica Acta, v. 14, p. 3–30, 2001.

414 Bulut, F., Bohnhoff, M., Eken, T., Janssen, C., Kılıç, T., and Dresen, G.: The East Anatolian

415 fault zone: Seismotectonic setting and spatiotemporal characteristics of seismicity based

416 on precise earthquake locations: Journal of Geophysical Research, v. 117, B07304,
417 https://doi.org/10.1029/2011JB008966, 2012.

418 Çemen, I., Göncüoglu, M.C., and Dirik, K.: Structural evolution of the Tuz Gölü basin in central

419 Anatolia, Turkey: Journal of Geology, v. 107, p. 693–706, https://doi.org/10.1086
420 /314379, 1999.

421 Çubuk Y, Yolsal-Çevikbilen S, Taymaz, T.: Source parameters of the 20052008 Bal<sup>a</sup>Sirapinar
422 (central Turkey) earthquakes: Implications for the internal deformation of the Anatolian
423 plate. Tectonophysics 635(Supplement C) :125 – 153, 2014.

424 Denolle, M. A., and Shearer, P.M.: New perspectives on self-similarity for shallow thrust 425 earthquakes, *J. Geophys. Res. Solid Earth*, *121*, 6533–6565, 2016.

426 Edwards, B., & Rietbrock, A.: A comparative study on attenuation and source-scaling relations
427 in the Kantō, Tokai, and Chubu regions of Japan, using data from Hi-net and KiK-net.
428 Bulletin of the Seismological Society of America, 99, 2435–2460, 2009.

429 Eken, T., Mayeda, K., Hofstetter, A., Gök, R., Orgülü, G. and Turkelli, N.: An application of the
430 coda methodology for moment-rate spectra using broadband stations in Turkey.
431 Geophys. Res. Lett, 31, L11609, 2004.

432 Eulenfeld, T. and Wegler, U.: Measurement of intrinsic and scattering attenuation of shear
433 waves in two sedimentary basins and comparison to crystalline sites in Germany,
434 Geophys J Int., 205(2):744-757, 2016.

- 435 Eulenfeld, T. and Wegler, U.: Crustal intrinsic and scattering attenuation of high-frequency
  436 shear waves in the contiguous United States. J Geophys., Res, 122, 2017.
- 437 Favreau, P., and Archuleta, R.J.: Direct seismic energy modelling and application to the 1979
  438 Imperial Valley earthquake, Geophys. Res. Lett., 30, 1198, 2003.
- Gaebler, P.J., Eulenfeld, T. & Wegler, U.: Seismic scattering and absorption parameters in the
  W-Bohemia/Vogtland region from elastic and acoustic radiative transfer theory,
  Geophys. J. Int., 203(3), 1471–1481, 2015.
- Gaebler, P.J., Eken, T., Bektaş, H.Ö, Eulenfeld, T., Wegler, U., Taymaz, T.: Imaging of Shear
  Wave Attenuation Along the Central Part of the North Anatolian Fault Zone, Turkey,
  submitted to the Journal of Seismology, 2018.
- 445 Goertz-Allmann, B. P., Edwards, B., Bethmann, F., Deichmann, N., Clinton, J., Fäh, D., &
- Giardini, D.: A new empirical magnitude scaling relation for Switzerland. Bulletin of the
  Seismological Society of America, 101, 3088–3095, 2011.
- 448 Gök, R., Kaviani, A., Matzel, E. M., Pasyanos, M. E., Mayeda, K., Yetirmishli, G., El-Hussain,
- 449 I., Al-Amri, A., Al-Jeri, F., Godoladze, T., Kalafat, D., Sandvol, E. A., and Walter,
- W.R.: Moment Magnitudes of Local/Regional Events from 1D Coda Calibrations in the
  Broader Middle East Region. Bull Seismol Soc Am., 106(5):1926-1938, 2016.
- 452 Gusev, A.A. & Abubakirov, I.R.: Simulated envelopes of non-isotropically scattered body
  453 waves as compared to observed ones: another manifestation of fractal heterogeneity,
  454 Geophys. J. Int., 127(1), 49–60, 1996.
- 455 Hanks, T.C. and Kanamori, H.: A moment magnitude scale, J. Geophys., Res., 84, 2348–2350,456 1979.

- 457 Izgi, G., Eken, T., Gaebler, P., and Taymaz, T.: Frequency-Dependent Shear Wave Attenuation
  Along the Western Part of the North Anatolian Fault Zone, Geophysical Research
  Abstracts, Vol. 20, EGU2018-629-2, 2018.
- 460 Kaymakci, N. Inceöz, M. Ertepinar, P.: 3D architecture and Neogene evolution of the Malatya
  461 Basin: inferences for the kinematics of the Malatya and Ovacik Fault Zones. Turkish
  462 Journal of Earth Sciences, 15, 123-154, 2006.
- 463 Kwiatek, G., Plenkers, K. & Dresen, G.: 2011. Source parameters of pico-seismicity recorded at
  464 Mponeng Deep Gold Mine, South Africa: implications for scaling relations, Bull. seism.
  465 Soc. Am., 101(6), 2592–2608, 2011.
- 466 Malagnini, L., Mayeda, K., Akinci, A., and Bragato, P. L.: Estimating absolute site effects,
  467 Bull. Seismol. Soc. Am. 94, no. 4, 1343–1352, 2004.
- Malagnini, L., and Munafò., I.: On the Relationship between M<sub>L</sub> and Mw in a Broad Range: An
  Example from the Apennines, Italy, Bulletin of the Seismological Society of America,
  Vol. 108, No. 2, pp. 1018–1024, 2018.
- 471 Mayeda, K., and Walter, W.R.: Moment, energy, stress drop, and source spectra of western
  472 United States earthquakes from regional coda envelopes, J. Geophys. Res. 101, 11,195–
  473 11,208, 1996.
- 474 Mayeda, K., Hofstetter, A., O'Boyle, J.L., and Walter, W.R.: Stable and transportable regional
  475 magnitudes based on coda-derived moment-rate spectra, Bull. Seismol. Soc. Am. 93,
  476 224–239: 2003.
- 477 Mayeda, K., Malagnini, L., Phillips, W. S., Walter, W. R., and Dreger, D.: 2D or not 2D, that is
  478 the question: A Northern California Test. Geophys- ical Research Letters, 32(12), 2005.
- 479 Morasca, P., Mayeda, K., Malagnini, L. and Walter, W.R.: Coda and direct-wave attenuation
  480 tomography in northern Italy, Bull Seismol Soc Am., v. 98, pages, 1936-1946, 2004.

- 481 Morasca, P., Mayeda, K., Gök, R., Phillips, W.S., and Malagnini, L.: Coda-derived source
  482 spectra, moment magnitudes and energy-moment scaling in the western Alps, Geophys.
  483 J. Int., 160, 263–275, 2008.
- 484 Paasschens, J.: Solution of the time-dependent Boltzmann equation, Phys. Rev. E, 56(1), 1135–
  485 1141, 1997.
- Papageorgiou, A., and Aki, K.: A specific barrier model for the quantitative description of
  inhomogeneous faulting and the prediction of strong ground motion I: Description of the
  model, Bull. Seismol. Soc. Am., 73(3), 693–722, 1983.
- 489 Pasyanos, M. E., R. Gök, and Walter, W.R.: 2-D variations in coda amplitudes in the Middle
  490 East. Bull. Seismol. Soc. Am. 106, no. 5, 2016.
- 491 Phillips, W. S., Mayeda, K. M., and Malagnini, L.: How to invert multi-band, regional phase
  amplitudes for 2-d attenuation and source parameters: Tests using the usarray. Pure and
  Applied Geophysics, 171(3):469-484, 2014.
- 494 Portner, D.E., Delph, J.R., Biryol, C.B., Beck, S.L., Zandt, G., Özacar, A.A., Sandvol, E., and
  495 Türkelli, N.: Subduction termination through progressive slab deformation across
  496 Eastern Mediterranean subduction zones from updated P-wave tomography beneath
  497 Anatolia, Geosphere, 14(3): 907-925, 2018.
- 498 Przybilla, J. and Korn, M.: Monte Carlo simulation of radiative energy transfer in continuous
  elastic random mediathree-component envelopes and numerical validation. Geophys J
  500 Int , 173(2):566-576, 2008.
- 501 Rautian, T.G. & Khalturin, V.I.: The use of the coda for determination of the earthquake source
  502 spectrum, Bull. Seism. Soc. Am., 68(4), 923–948, 1978.
- 503 Sato, H. and Fehler, M.C.: Seismic Wave Propagation and Scattering in the Heterogeneous
  504 Earth, Springer-Verlag, New York, 1998.

505 Sato, H., Fehler, M.C. & Maeda, T. Seismic Wave Propagation, and Scattering in the 506 Heterogeneous Earth, 2nd edn, Springer: 2012.

- 507 Sens-Schönfelder, C. and Wegler, U.: Radiative transfer theory for estimation of the seismic
  508 moment. Geophys J Int, 167(3):1363-1372.
- 509 Taymaz, T., Jackson, J., Westaway, R.: Earthquake mechanisms in the Hellenic Trench near
  510 Crete. Geophys. J. Int.102, 695–731, 1990.
- 511 Taymaz, T., Westaway, R., Reilinger, R.: Active faulting and crustal deformation in the eastern
  512 Mediterranean Region. Spec. Issue Tectonophys. 391 (1-4), 1–9. http://
  513 dx.doi.org/10.1016/j.tecto.2004.07.005, 2004.
- 514 Tchalenko, J. S.: Similarities between shear zones of different magnitudes. Geol. Soc. Am.
  515 Bull., 81, 1625–1640, 1970.
- 516 Uchide, T., & Imanishi, K.: Small earthquakes deviate from the omega-square model as revealed
  517 by multiple spectral ratio analysis. Bulletin of the Seismological Society of America,
  518 106(3), 1357–1363, 2016.

519 Uchide, T., & Imanishi, K.: Underestimation of microearthquake size by the magnitude scale of

- the Japan Meteorological Agency: Influence on earthquake statistics. Journal of
  Geophysical Research: Solid Earth, 123, 606–620, 2018.
- 522 Yoo, S.-H., Rhie, J., Choi, H.-S., and Mayeda, K.: Coda-derived source parameters of
  earthquakes and their scaling relationships in the Korean Peninsula, Bull. Seismol. Soc.
  Am., 101, 2388–2398, 2011.
- 525 Zeng, Y., Su, F. and Aki, K.: Scattering wave energy propagation in a random isotropic
  526 scattering medium: 1. Theory, J. Geophys. Res., 96(B1), 607–619, 1991.
- 527
- 528
- 529

530 *Figure Captions* 

Figure 1. Major tectonic features of Turkey and its adjacent. The plate boundary data used
here is taken from Bird (2003). Subduction zones are black, continental transform faults are
red, continental rift boundaries are green, and spreading ridges boundaries are yellow. NAFZ,
EAFZ, and DSFZ are the North Anatolian Fault, East Anatolian Fault, and the Dead Sea fault,
respectively.

536

Figure 2. Epicentral distribution of all local events selected from the study area in the KOERI catalogue. Gray circles represent earthquakes with poor quality that are not considered for the current study while black indicates the location of local events with good quality. Red circles among these events are 487 events used in coda wave inversion since they are successful at passing quality criteria of further pre-processing procedure.

542

543 Figure 3. An example from the inversion procedure explained in chapter 3. Here coda 544 envelope fitting optimization is performed on band-pass filtered (4-8Hz) digital recordings of 545 an earthquake (2014 April 09, M<sub>W</sub>-coda3.2) extracted for 7 seismic stations that operated within the CD-CAT array. Large panel at the lower left-hand side displays the error function  $\varepsilon$ 546 as a function of  $g_0$ . Thick blue cross here represent the optimal value of  $g = g_0$ . Other small 547 548 panels at upper and right-hand side show the least- squares solution of the weighted linear equation system for the first 6 guesses and optimal guess for  $g_0$ . The dots and gray curves 549 indicate the ratio between energy (E<sup>obs</sup>) and the Green's function (G) obtained for direct S-550 551 waves and observed envelopes at various stations, respectively (Please notice that during this optimization process envelopes are corrected for the obtained site corrections R<sub>i</sub>). The slope 552 of linear curve at each small panel yields –b in relation to the intrinsic attenuation. The linear 553

554 curve has an intercept of W representing source related terms at the right-hand side of 555 equation 5.

556

Figure 4. a) Results of the inversion of the 2014-April-09, M<sub>w</sub>-coda3.2 earthquake: Sample 557 558 fits between observed and calculated energy densities in the frequency band 0.5–1.0 Hz are 559 given for 6 different stations (see upper right corner for event ID, station name, and distance to hypocenter). Note that light blue curves represent observed envelope. Smoothed observed 560 561 calculated envelopes in each panel are presented by blue and red curves, respectively. Blue 562 and red dots exhibit location of the average value for observed and calculated envelopes 563 within the S-wave window, respectively. b) The same as in (a) obtained in the frequency band 564 4.0-8.0 Hz.

565

Figure 5. All individual observed (black squares) and predicted (gray curve) source
displacement spectra observed at 72 stations from 487 local earthquakes in central Anatolia.

Figure 6: Scatter plot between local magnitudes  $(M_L)$  of analyzed events with coda wavesderived magnitudes  $(M_W$ -coda) of the same events. The outcome of a linear regression analysis yielded an emprical formula (e.g. Eq. 10) to identify the overall agreement represented by gray straight line. Yellow and red dashed lines indicate upper and lower limit of linearly fitting to that scatter.

574

Figure 7: Same scatter plot displayed in Fig. 6. Here color code indicates estimated high-frequency fall-off parameter for each inverted event.

577









Figure 2. 

599

601



610 Figure 3.



- 622 623



Frequency (Hz)

625 626 636 637 

Figure 5.

650





