

Analysing stress field conditions of the Colima Volcanic Complex (Mexico) by integrating FEM simulations and geological data

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Abstract

In the last decades numerical methods have become very popular tools in volcanological studies, since capable of considering many relevant parameters in their calculations, such as the presence of multiple reservoirs, topography, and heterogeneous distribution of host rock mechanical properties. Although the widespread availability of geodetic data is keep growing, the influence of geological data on the numerical simulations is still poorly considered. In this work a 2D Finite Element Modelling is provided by using the LInear Static Analysis (LISA) software, in order to investigate the stress field conditions occurring around the Colima Volcanic Complex (CVC, Mexico) at increasing the details of geological and geophysical input data. By integrating the published geophysical, volcanological, and petrological data, we provide a first-order description of the domain of the CVC feeding system, considering either one or two magma chambers connected to the surface via dykes or isolated (not connected) in the elastic host rocks. We test the methodology by using a gravitational modelling with different geometrical configurations and constraints (i.e. magma chamber dimensions, depth, overpressure). Our results suggest that an appropriate set of geological data is of pivotal importance for improving the mesh generation procedures and the degree of accuracy of numerical outputs, aimed to more reliable physics-based representations of the natural systems.

1 Introduction

Magmatism and tectonism in volcanic active areas are strongly related to the regional and local stress fields, affecting both the orientation of faults and the location of volcanic vents, two fundamental

38 aspects when interpreting volcanic unrest and forecasting volcanic eruptions (Geyer et al., 2016). The
39 stress field around a magmatic source originates from three main contributions: (1) the background
40 stress, composed of a vertical gravitational load and a lateral horizontal load corresponding to
41 lithostatic confinement and tectonic regimes; (2) the stress field caused by the loading of the volcano
42 edifice; and (3) the stress field generated by the magmatic overpressure in the chamber system (e.g.
43 Martí and Geyer, 2009; Currenti and Williams et al., 2014). In recent years, a large number of semi-
44 analytical and numerical solutions for the stress field state of geological and volcanological systems
45 have been proposed (e.g. Cayol and Cornet, 1998; Simms and Garven, 2004; Manconi et al., 2007;
46 Long and Grosfils, 2009; Currenti et al., 2010; Currenti and Williams et al., 2014; Zehner et al.,
47 2015), taking into account the static elastic deformation in a multi-layered half-space (e.g. Dieterich
48 and Decker, 1975; Bonafede et al., 2002; Wang et al., 2003; Gudmundsson and Brenner, 2004; Zhao
49 et al., 2004; Pritchard and Simons, 2004; Gottsmann et al., 2006; Geyer and Gottsmann, 2010; Zhong
50 et al., 2019). Following the successful application in mechanical engineering, the use of Finite
51 Element Method (FEM) has been extensively introduced in Earth Sciences in order to investigate the
52 effects of topography, lithologic heterogeneities, tectonic stresses and the gravity field on the Earth's
53 surface deformation (e.g. Cailleau et al., 2003; 2005; Buchmann and Conolly 2007; Manconi et al.,
54 2009; Masterlak et al., 2012), including volcanoes (e.g. Fujita et al., 2013; Carcho and Gàlan del
55 Sastre, 2014; Bunney, 2014; Ronchin et al., 2015; Hickey et al., 2015; Cabaniss et al., 2019; Rivalta
56 et al., 2019).

57 The use of FEM in volcanic areas has several examples, which vary from the influence of layered
58 materials on the surface deformation process during volcanic inflation (e.g. Darwin volcano,
59 Galapagos Islands; Manconi et al., 2007; Albino et al., 2010) to processes affecting chamber rupture
60 (e.g. Grosfils, 2007; Long and Grosfils, 2009). FEM is also used in fluid dynamics and
61 thermodynamics (e.g., Gutiérrez and Parada, 2010; Gelman et al., 2013) for solving issues related to
62 motion of fluids and heat transfer.

63 The local stress around a volcanic feeding system strongly depends on the magma chamber geometry
64 and on the mechanical properties of the layered host rock around it (e.g. Martí and Geyer, 2009),
65 mainly due to broad changes in Young's modulus (e.g. Gudmundsson et al., 2011; Jeanne et al.,
66 2017; Heap et al., 2020). For instance, limestones, lava flows, welded pyroclastic units and intrusive

67 rocks can be very stiff (high Young's modulus; from ca. 1.7 to 27 GPa for limestones, Touloukian,
68 1981; ca. 5.4 GPa for volcanic rocks, Heap et al., 2020), whereas young and non-welded pyroclastic
69 units may be very soft (low Young's modulus; ca. 1.7 – 3.1 GPa, Margottini et al., 2013).
70 Consequently, the local stress may change abruptly from one layer to another (e.g., Gudmundsson,
71 2006). Irrespective of the scope of the numerical investigation, the importance of applying accurate
72 rheological constraints to FEM modelling was discussed in many studies (e.g., Folch et al., 2000;
73 Newman et al., 2001; Fernandez et al., 2001; Currenti et al., 2010; Geshi et al., 2012). This implies
74 that geology of the volcanic area needs to be considered as more accurate as possible. However, few
75 investigations have been carried out to assess the influence of the amount and quality of geological
76 data into FEM computations (Kinvig et al., 2009; Norini et al., 2010, 2019; Cianetti et al., 2012;
77 Ronchin et al., 2013; Chaput et al., 2014). To bridge this gap, in this work we use the Linear Static
78 Analysis (LISA) software (version 8.0; www.lisafea.com) to study the subsurface stress behaviour in
79 an elastic domain at Colima Volcanic Complex (CVC, Mexico) when improving the description of
80 geological constraints.

81 The CVC area is a good candidate for testing the response of FEM software to different geological
82 conditions, being constituted by a large volcanic complex (significant topographic load; Lungarini et
83 al., 2005), a well-defined feeding system inferred from geophysical and petrological data (e.g. Spica
84 et al., 2017; Massaro et al. 2018, 2019), and growth within a tectonic graben (bordered by normal
85 faults; Fig. 1a) infilled by volcanoclastic material (variability of rock mechanical characteristics;
86 Norini et al., 2010, 2019).

87 In this light, the present study proposes a contribution to a more proper use of FEM models for
88 assessing the stress state pattern in volcanic areas at different levels of description of the geological
89 features. In particular, we focus on the CVC by using the available published data of the inferred
90 feeding system structure, in order to assess how the addition of geological and volcanological
91 constraints (i.e. stratigraphy, geometry of the plumbing system, extensional tectonic regime, local
92 fault systems) may, and at what extent, affect the model outputs (Fig 1b). Beside and beyond the
93 evaluation of geological details on FEM outputs, we also obtained a picture of the large-scale stress
94 distribution in the CVC subsurface.

95

96 **2 The Colima Volcanic Complex (Mexico)**

97 *2.1 Geological framework*

98 The Pleistocene-Holocene CVC is one of the most prominent volcanic edifices within the Trans-
99 Mexican Volcanic Belt (TMVB) (Macías et al., 2006; Capra et al., 2016; Norini et al., 2019; Fig. 1a).
100 In this area, the Rivera microplate and the Cocos plate subduct beneath the North America plate
101 along the Middle American Trench, producing great deformation and fragmentation of the
102 continental plate (Stock and Lee, 1994), and forming a triple junction that delimits the tectonic units
103 known as the Jalisco Block (JB) and the Michoacán Block (MB) (Luhr et al., 1985; Allan, 1986;
104 Rosas-Elguera et al., 1996; Rosas-Elguera et al., 1997; Ferrari and Rosas- Elguera, 1999; Rosas-
105 Elguera et al., 2003; Frey et al., 2007). The three rifts of this system are the Tepic-Zacoalco (TZR),
106 the Chapala-Tula (CTR), and the Colima Rift (CR) where the CVC is emplaced (Allan, 1986;
107 Escudero and Bandy, 2017). The still active NS trending Colima Rift (CR) was formed during an
108 extensional phase occurred after the Late Cretaceous–Paleogene compressive and transpressive phase
109 (Allan, 1986; Serpa et al., 1992; Bandy et al., 1995; Cortés et al., 2010). The rifting phase deformed
110 Cretaceous marine limestones, Jurassic–Tertiary metamorphosed clastic and volcanoclastic sediments,
111 Cretaceous–Tertiary intrusive rocks and Tertiary-Quaternary volcanic deposits along sub-vertical
112 crustal faults. While opening, CR was gradually filled with Pliocene–Quaternary lacustrine sediments,
113 alluvium and colluvium (e.g. Allan, 1986; Allan et al., 1991; Norini et al., 2010). The geometry,
114 kinematics and dynamics of the CR have been studied on the basis of field, seismic, and geodetic
115 data, mainly collected in its northern and central sectors (see Fig. 1 in Norini et al., 2010).

116 The amount of vertical displacement of the northern and central sectors is estimated to be at least 2.5
117 km by adding the topographic relief of the bounding fault scarps (1.5–1.6 km) to the calculated
118 sediment depth (Allan, 1985; Serpa et al., 1992). Field data and focal mechanism solutions are
119 consistent with a direction of opening of the northern and central sectors oriented from E-W to NW-
120 SE, with a mainly normal and minor right-lateral displacements of the bounding faults (Barrier et al.,
121 1990; Suárez et al., 1994; Rosas-Elguera et al., 1996; Garduño-Monroy et al., 1998; Norini et al.,
122 2010, 2019). In contrast to field and seismic evidence of long-term slightly dextral oblique extension,
123 recent GPS geodetic measurements suggest a possible sinistral oblique extension of the CR (Selvans
124 et al., 2011). In both cases, the stress regime is mainly extensional, with an approximately E-W

125 orientation of the minimum horizontal stress in the basement of the CVC (Barrier et al., 1990; Suárez
126 et al., 1994; Rosas-Elguera et al., 1996; Selvans et al., 2011; Norini et al., 2010, 2019).

127 The CVC stands within the central sector of the CR, on top of the Cretaceous limestones, Late
128 Miocene-Pleistocene volcanic rocks, and Pliocene-Holocene lacustrine sediments, alluvium, and
129 colluvium (Allan, 1985, 1986, 1991; Cortès, 2005; Norini et al., 2010). The volcanic complex is
130 affected and displaced by the N-S/NNE-SSW-trending recent-active crustal faults of the CR,
131 controlling the geometry and location of the volcano feeding system (Fig. 1a). Indeed, the CVC was
132 formed by three andesitic stratovolcanoes aligned parallel to the CR bounding faults: the northern
133 inactive Cantaro volcano (2900 m a.s.l.), followed by the inactive Nevado de Colima (4255 m a.s.l.)
134 and, in the southern part, the youngest and active Volcán de Colima (3763 m a.s.l.) (Norini et al.,
135 2019 and reference therein).

136

137 *2.2 Eruptive activity*

138 The eruptive history of the CVC started in the northeast area with the formation of Cantaro volcano
139 at ca. 1-1.5 Ma. The volcanic activity of the Nevado de Colima started at ca. 0.53 Ma. It is composed
140 of voluminous andesitic lava domes and flows and pyroclastic deposits associated with caldera
141 forming eruptions and numerous partial sector collapses (Robin et al., 1987; Roverato et al., 2011;
142 Roverato and Capra, 2013; Cortès et al., 2019). The youngest Volcán de Colima, now considered one
143 of the most active volcanoes of the world, consists of the Paleofuego edifice that suffered several
144 sector collapses, with the formation of a horseshoe-shaped depression where the new active cone
145 (also known Volcán de Fuego) grew up, through Merapi and Soufrière type dome collapses,
146 extrusion of lava flows, Vulcanian and occasionally sub-Plinian explosive eruptions (Saucedo et al.,
147 2010; Massaro et al., 2018, 2019). The activity of both Nevado and Volcán de Colima volcanoes also
148 included several sector collapses, occurred frequently in the Upper Pleistocene and Holocene,
149 repeatedly devastating the floor of the Colima Rift down to the Pacific Ocean (Robin et al., 1987;
150 Luhr and Prestegard, 1988; Stoopes and Sheridan, 1992; Capra and Macias, 2002; Cortès, 2005;
151 Roverato et al., 2011).

152

153 *2.3 The CVC plumbing system*

154 Spica et al. (2017) indicate a 15 km-deep low velocity body (LVB) as the CVC deep magma
155 reservoir. Its horizontal extension seems to be delimited by the borders of the CR, suggesting a
156 structural control of the normal fault system on it (Spica et al., 2014). The LVB has an extent of ca.
157 55 km × 30 km in the N-S and E-W directions respectively, showing a mean thickness < 8 km.
158 Escudero and Bandy (2017) obtained a higher resolution tomographic image of the subsurface in the
159 CVC area, showing that the most active magma generation zone is presently under the Fuego de
160 Colima edifice. Here, the ambient seismic noise tomographic study proposed by Spica et al. (2014)
161 confirmed the presence of a shallow magma chamber above ca. 7 km depth, as also demonstrated by
162 petrological studies (Medina-Martinez et al., 1996; Luhr, 2002; Zobin et al., 2002; López-Loera et al.,
163 2011; Reubi et al., 2013, 2019; Macías et al., 2017). Cabrera-Gutiérrez and Espíndola (2010)
164 suggested the shallow active magma storage has a volume of ca. 30 km³. The shallow magma
165 chamber is connected to the surface by a dyke/conduit system, whose path is facilitated by the
166 presence of the CR fault zone, which provides a natural pathway for fluids (e.g., Allan, 1986; Norini
167 et al., 2010, 2019). The arrangement of dykes and the alignment of volcanic centres of CVC suggest
168 that the dykes swarm draining the magma chambers developed along the NNE-SSW-trending, steep,
169 eastward dipping normal fault exposed on the northern CVC flank (Norini et al., 2010, 2019).

170 Taking into account the previous information, Massaro et al. (2018) provided a first-order
171 geometrical reconstruction of the Fuego de Colima feeding system during the 1913 sub-Plinian
172 eruption, by using volcanological data (Saucedo et al., 2010, 2011; Bonasia et al., 2011) as input and
173 constraints for numerical simulations. Results showed good matches for a hybrid configuration of the
174 shallow conduit-feeding system (i.e., dyke developing into a shallower cylindrical conduit). The best-
175 fit dyke geometry has width in the range from 200 m to 2000 m and thickness of ca. 40 m, with the
176 cylindrical conduit diameter similar to the dyke thickness. The shallow magma chamber top was set
177 at 6 km of depth, and dyke-cylinder transition at 500 m below the summit, as also inferred from
178 geophysical data (Salzer et al., 2014; Aràmbula et al., 2018).

179

180 **3 Methods**

181 In this study, we used the commercial 8.0 version of LISA software (www.lisafea.com). LISA is a
182 general-purpose Finite Element Analysis (FEA) software developed in the '90s based on the
183 formulations proposed by Rao (1989). Since then, formulations from many other sources were also
184 integrated (Bathe, 1990; Michaeli, 1991; Schwarz, 1991; Babuska et al., 1995). Despite FEA was
185 originally used for structural analysis (Rao, 1989; 2013), it is also able to successfully predict the
186 stress-strain behaviour of rock masses accounting for elastic models, in particular the deformation
187 and failure mechanisms even in layered rock masses (Gabrieli et al., 2015).

188 Simplifying techniques in structural FEA can give valuable insights into local stresses more rapidly
189 and efficiently than a full 3D model. Here we considered a 2D model throughout a complex structure
190 (i.e. dual magma chamber feeding system, rift system, rock layering, and faults), in order to
191 investigate the stress behaviour induced in the host rocks in response to the increasing detail of
192 geological data used to constrain the model.

193

194 *3.1 Modelling approach*

195 Taking into account the works of Norini et al. (2010, 2019), we simulated the stress field of the CVC
196 plumbing system considering an E-W cross-section, which is parallel to the extension associated to
197 the active Colima Rift (Norini et al., 2010), shown in Figure 1a-b (a-a').

198 Since the extent of the CVC magma chambers in the NNE-SSW direction is typically much longer
199 than the dimensions of the E-W cross section (Spica et al., 2017), 2D solutions of either numerical or
200 analytical models describing E-W elongated magma chambers in the crust can be reasonably adopted
201 (Jaeger et al., 2009; Costa et al., 2011). A topographic profile and 2D plane along the chosen E-W
202 cross-section of the CVC area was obtained in ESRI ArcGIS from a Digital Elevation Model (DEM;
203 resolution 50 m; Instituto Nacional de Estadística y Geografía - INEGI <https://en.www.inegi.org.mx/>).

204 This cross section was imported into Autodesk Auto-Cad R13 and approximated to a third-degree
205 spline. Finally, the IGES file was imported into LISA, where the mesh discretization was performed.
206 The domain was discretized by three and four-node finite elements (Table 1; Fig. 1c). The volcanic

207 area domain extends 60 km horizontally and 30 km below the surface set in an x - z Cartesian
208 Coordinate System. Zero normal displacements are assigned at the bottom and the lateral boundaries
209 of the domain, while the upper boundary representing the ground surface is stress free (Fig. 1c). The
210 analysis is carried out by using a plane strain approximation, implying that the deformation in the
211 third direction is assumed to be negligible.

212 FEM of geological structures requires accurate discretization of the computational domain such that
213 geological units are represented correctly. Zehner et al. (2015) reported that the unstructured
214 tetrahedral meshes on a complex geological model has to fulfil the following requirements: i)
215 sufficient mesh quality: the tetrahedrons should not be too acute-angled, since numerical instabilities
216 can occur, ii) incorporation of geometry for defining boundary conditions and constraints, iii) local
217 adaption, which is a refinement of the mesh in the vicinity of physical sources in order to avoid
218 numerical errors during the simulation. Considering these requirements, in this work we adopt as the
219 best discretization a mesh with 4660 plane continuum elements for the E-W cross-section. The size of
220 finite elements was refined in the regions with higher gradients, especially near the contours of the
221 magmatic feeding systems.

222 In our simulations, the extent of the rock layers (Table 2) is referred to the model of Norini et al.
223 (2010, 2019). Magma chambers and dykes are considered as pressurized finite-size bodies in an
224 elastic crustal segment, acting as fluid-filled holes. The boundary condition (pressurization) is
225 provided by applying internal forces that act on the walls. This approach has been extensively used in
226 several analytical and numerical models that treat magma reservoirs as internally pressurized
227 ellipsoidal cavities within an elastic half space, in order to gain insight into the behaviour of magma
228 plumbing systems (Pinel and Jaupart, 2004; Gudmundsson, 2006; Grosfils, 2007; Andrew and
229 Gudmundsson, 2008; Hautmann et al., 2013; Currenti and Williams, 2014; Zhong et al., 2019).

230 The geometrical configuration set for the CVC feeding system (i.e. the shape and dimensions of the
231 magmatic chambers) derives from the literature (Spica et al., 2014, 2017; Massaro et al., 2018, 2019)
232 and it is simplified in Figure 1d. The overpressure in magma chambers may be produced by a variety
233 of processes, including fractional crystallization, volatile exsolution and magma recharge, leading to
234 deviatoric stresses in the country rock that may be tens of MPa in magnitude (Jellinek and DePaolo,
235 2003; Karlstrom et al., 2010).

236 Previously published studies indicate that differences between, and problems with, elastic models
237 derive principally from the key role played by gravity (e.g. Albino et al., 2018). Gravity plays a first
238 order role on bedrock failure conditions (Gerbault, 2012), on the geometry of magma propagation
239 with respect to an edifice load and on buoyancy contrasts driving magma upward (Lister and Kerr,
240 1991; Watanabe et al., 2002). However, in a wide variety of simulations of natural phenomena the
241 gravitational effects are often incorporated either incorrectly or incompletely (e.g. Grosfils, 2007).
242 Some authors argued on whether it is appropriate or not to account for the gravity body force in
243 numerical models of volcanic inflation (e.g. Currenti and Williams, 2014; Grosfils et al., 2015).
244 When the gravitational loading is not included in the model, the volcanic deformation results from a
245 change with respect to a stage previously at equilibrium (e.g. Gerbault et al. 2018). In this work, we
246 carried out simulations considering the effect of the gravitational loading. Gravity in the host rock is
247 implemented via body forces. The model initial condition has a pre-assigned lithostatic stress, whose
248 computation, in presence of topography and material heterogeneities, is not trivial because it requires
249 applying the gravity load preserving the original not deformed geometry of the mesh (Cianetti et al.,
250 2012). Since the presence of a lithostatic stress field, the load applied at the reservoir boundaries
251 represents a superposition of the magmatic overpressure and lithostatic component.

252 We also took into account the effect of the existing faults of the Colima Graben (CG) system even if
253 LISA cannot include a frictional law to represent the fault movement (i.e. Chaput et al., 2014). As
254 reported in Jeanne et al. (2017 and reference therein) the damage induced by faults increases from the
255 host rocks to the fault core implying the reduction in the effective elastic moduli represented by a
256 progressive decrease in Young's Modulus. Considering the evaluation of fault zone elastic properties
257 provided by Jeanne et al. (2017), we represented the faults bordering the CG as two damage zones
258 inclined of ca. 70° and with a thickness of ca. 1 km, showing reduced elastic properties with respect
259 to the surrounding host rocks down to 10 km in depth.

260 It is important to note that we chose to represent the different simulations using different colour
261 scales. Although such a choice makes more difficult a visual comparison of the simulation outputs
262 and it needs to be kept in mind looking at the different figures, it preserves the necessary details of

263 stress distribution, which would have been lost using a common colour scale for all the figures in
264 LISA.

265

266 **4 Geological data**

267 In this work, we used geological information available in literature as input data, in order to estimate
268 the stress variations around the CVC magmatic plumbing system. Here we briefly describe the main
269 geological features taken into account in LISA simulations.

270 *4.1 Stratigraphy*

271 Four units forming the CVC system were defined from the available geological data (Table 2): i)
272 Basement (Unit B): cretaceous limestones and intrusive rocks forming the bed-rock underlying the
273 CVC; ii) Graben fill deposits (Unit GF): Quaternary alluvial, colluvial, and lacustrine deposits filling
274 the graben; iii) Fuego de Colima deposits (Unit FC): andesitic lavas and pyroclastic deposits forming
275 the Paleofuego-Fuego de Colima edifices; and iv) Volcaniclastic deposits (Unit VD): volcaniclastic
276 deposits covering the southern flank of the CVC (e.g. Cortés et al. 2010; Norini et al., 2010, 2019).
277 Being the area interested by FEM extended down to 30 km, it is evident how Unit B is dominant with
278 respect to the others, which occupy only few km in the upper part of the simulated domain. We
279 assumed constant mechanical characteristics within each Unit (Table 2). In particular, Unit B was
280 considered mechanically homogeneous with elastic properties of a carbonate, due to the lack more
281 detailed information of deeper lithologies (Norini et al., 2019).

282 Deformation within the brittle upper crust is described by elastic material behaviour. For each Unit
283 we fixed typical rock mass properties, density (ρ), Young's Modulus (E) and Poisson's Ratio (ν)
284 (Table 2). The rock masses are considered dry, in order (eventual) pore pressure to be neglected.
285 Only for Unit GF a higher value for the Poisson's Ratio was used close to the surface in order to
286 mimic high water content in the graben sediments. The maximum thickness of the graben fill (about
287 1 km) was assumed from the literature (Allan, 1985; Serpa et al., 1992; Norini et al., 2010, 2019). For
288 Units B and GF rock mass proprieties were derived from Hoek and Brown (1997) and Marinos and
289 Hoek (2000), while for volcanic materials (units FC and VD; Table 2) were estimated according to

290 the approach proposed by Del Potro and Hürlimann (2008). This information allowed Norini et al.
291 (2019) to derive the equivalent Mohr-Coulomb properties for the stress ranges expected in the
292 different sectors of the CVC. In addition, in order to describe the effects of the CG faults on stress
293 field distribution, the mechanical properties were locally degraded in proximity of the faults
294 themselves.

295 *4.2 The geometry of the plumbing system*

296 The geometry of the E-W cross-section of the CVC plumbing system was modelled taking into
297 account the previous subsurface information described in Section 4.1. In our 2D model, we assumed
298 the CVC composed of a two magma chambers connected by dykes and to the surface by a conduit
299 (Fig. 1d). The shape of the magma chambers and dykes are represented by elliptical cross-sections
300 with the major (*2a*) and minor (*2b*) axes.

301 Generally, the magma chambers have a sill-like shape that is often imaged in seismic studies of
302 volcanoes and rift zones (Macdonald, 1982; Sinton and Detrick, 1992; Mutter et al., 1995; MacLeod
303 and Yaouancq, 2000; Singh et al., 2006; Canales et al., 2009). Most of them are not totally molten but
304 rather a mixture of melt and crystal mush (i.e. Parfitt and Wilson, 2008). Various estimates have been
305 made to infer the actual amount of melt in a magmatic body, showing that it is only ca. 10% of the
306 total chamber volume (Gudmundsson et al., 2012 and reference therein).

307 Spica et al. (2017) described a 15 km-deep low velocity body (LVB) with its top at ca. 15 km of
308 depth and with an estimated volume of ca. 7000 km^3 , representing the deep magmatic reservoir of
309 CVC. Assuming the melt as 10%, the deep magma chamber volume would be ca. 700 km^3 .
310 Simplifying this volume in an elliptical sill-like geometry, the dimensions (i.e. *2a*, *2b*, *2c* axes) have
311 to be scaled according to those of LVB ($55 \times 30 \times 8 \text{ km}$; Spica et al., 2017). We therefore fixed *2a* =
312 14 km , *2b* = 3.6 km , *2c* = 26 km as the dimensions of the deep magma chamber, being *2c* elongated
313 in NW-SE direction.

314 For the shallow part of the feeding system, we have no detailed geophysical constraints. However,
315 Massaro et al. (2019) reproduced through numerical modelling the nonlinear cyclic eruptive activity
316 at Fuego de Colima in the last 20 years, using a shallow magma chamber volume in the range of 20-
317 50 km^3 , also according to the estimation of Cabrera-Gutiérrez and Espindola (2010). Assuming a

318 volume of 30 km^3 , we fixed $2a = 3.5 \text{ km}$, $2b = 2 \text{ km}$, $2c = 8 \text{ km}$ as dimensions of the shallow magma
319 chamber.

320 Numerous theoretical and field studies have established that host rock stresses dictate the magma
321 pathways (e.g. Maccaferri et al., 2011; Gudmundsson, 2011). During ascent to the surface, the dykes
322 align themselves with the most energy-efficient orientation, which is roughly perpendicular to the
323 least compressive principal stress axis σ_3 (e.g. Gonnermann and Taisne, 2015; Rivalta et al., 2019),
324 providing the magma driving pressure remains small compared to the deviatoric stress (Pinel et al.,
325 2017; Maccaferri et al., 2019). This behaviour, however, can be modulated in the presence of
326 significant variations in fracture toughness of the surrounding rock due to stratification (Maccaferri et
327 al., 2010) or to old and inactive fracture systems (Norini et al., 2019). Although for oblate magma
328 chambers the propagation of dykes is most probable from the tip areas, in our simulations the
329 orientation of dykes is assumed vertical, because of the preferential pathways represented by the CR
330 fault planes (Spica et al., 2017).

331 Although, for decades, magma conduits were modelled as cylinders, because of easiness of their
332 mathematical treatment, geophysical data and field observations highlight the importance and
333 peculiarities of dykes in magma transport and hence the need to adopt more realistic geometries
334 (Costa et al., 2009; Hautmann et al., 2013; Tibaldi, 2015). It is important to stress that although all
335 cavities/inclusions in a medium modify the local stress field and concentrate stresses, the induced
336 perturbation depends mainly on the geometry of the cavity/inclusion (Savin, 1961; Boresi et al.,
337 1985; Tan, 1994; Saada, 2009). We set the dimensions of feeder dykes in agreement with Massaro et
338 al. (2018): deep dyke $2ad = 2 \text{ km}$; shallow dyke $2a$ varies from 1 km at bottom to 500 m in the upper
339 part of the volcano; width of both deep and shallow dyke $2bd = 2b = 100 \text{ m}$ (Fig. 1d), although the
340 exact value of the latter is not crucial for the purposes of this study. Moreover, it is worth noting that
341 it is not the aim of this work to provide the conditions for the magma chamber rupture, being LISA
342 accounting only for the elastic regime. For these reasons, the selected magma overpressures (ΔP)
343 acting on the magma reservoirs and dykes have to be less than the tensile strength of the rocks. We
344 therefore fixed ΔP at 10 MPa and 20 MPa for the 15 km -deep chamber, and 5 MPa for the 6 km -deep
345 one. For the dykes and conduit, the magmatic overpressure is fixed at 10 MPa in the deeper dyke and
346 5 MPa in the shallower dyke, except for the upper 500 m of the shallower conduit where overpressure

347 is set at 0.4 MPa.

348 To take into account the effect of both far field extensive regime and CG around the magma feeding
349 system, we applied a uniform extension at the lateral boundaries of the domain (as reported in Martì
350 and Geyer, 2009) of 5 MPa and included two damage zones with reduced rock elastic moduli and
351 density (i.e. $E = 1 \text{ GPa}$, $\nu = 0.20$; Jeanne et al., 2017; $\rho = 1850 \text{ kg/m}^3$).

352

353

354 **5 Results**

355 The first part of this section is focused on a sensitivity analysis of Young modulus variation, aimed to
356 quantify the numerical effects of approximation of this important rock property on FEM outputs. The
357 second part of this section describes the model outputs when adding complexity to the input
358 geological/geophysical data.

359 Considering the E-W cross-section (a-a'; Fig. 1a), we provided six domain configurations with
360 increasing geological complexity: i) "*homogeneous lithology model*" in which the volcanic domain is
361 only composed of andesite rocks; ii) "*not homogeneous lithology model*" where different geological
362 units are considered; iii) "*single magma chamber model*" composed of a not homogeneous lithology
363 and a 15 km-deep magma chamber; iv) "*dual magma chamber model*" composed of a not
364 homogeneous and 6 km- and 15 km-deep magma chambers; v) "*conduit feeding system model*"
365 composed of not homogeneous lithology, 6 km- and 15 km-deep magma chambers connected
366 through a deep-dyke, and a shallow conduit connecting to the surface; vi) "*extensional model*", in
367 which we added a 5 MPa horizontal extensional stress (far field) and, vii) "*faulted model*", in which
368 two damaged zones mimicking the CG faults were added to the "*extensional model*" (local stress)
369 (Fig. 1b).

370 The number of nodes in the *only substratum* and *single magma chamber* models is set at 4426, for the
371 dual magma chamber model is set at 4161, and at 3737 for the *conduit feeding system* and *faulted*
372 models.

373 5.1 Sensitivity analysis on selected input parameters

374 In order to quantify the influence of Young Modulus selection on the model outputs, we performed a
375 sensitivity test using the single magma chamber model as reference case. We evaluated the influence
376 of varying the Young Modulus in each geological Units on the principal stresses σ_1 and σ_3 . Taking
377 into account the material properties used in the simulations (Norini et al., 2010, 2019; Table 2) as
378 reference values, we compared the stress state of the computational domain at changing (\pm) Young
379 Modulus by an order of magnitude. This variation has been separately applied to each Unit, in order
380 to assess what is the effect of changing material properties on model outputs. This sensitivity analysis,
381 although incomplete, may lead to raise awareness on the selection of input data when running a FEM.
382 The sensitivity analysis was carried out on a reduced simulation domain (the x -axis was set to 35 km)
383 in order to diminish the influence of binding effects that are present along domain borders.
384 We used the Euclidean norm (L2) method for illustrating the results of the sensitivity analysis. The
385 L2 norm applied on a vector space x (having components $i = 1, \dots, n$) is strongly related with the
386 Euclidean distance from its origin, and is equal to:

387

$$388 \quad ||x||_2 = \sqrt{\sum_i^n xi^2} \quad (1)$$

389

390 In our case, the vector space x is composed of all nodes of the computational domain (Table 1). We
391 defined x_{ref} the vector containing the results for the maximum and minimum principal stress when
392 using the selected values of material properties (Table 1) and $x(-)$, $x(+)$ the vectors at varying the
393 Young Modulus of one order of magnitude in each Unit.

394 We evaluated the global variation of stress in the proposed geometrical configurations of the domain
395 (i.e. not homogeneous lithology, single magma chamber, dual magma chamber, and dual magma
396 chamber with conduits models) calculating the global relative variation in L2 as follow:

397

$$398 \quad L2(-) = \frac{||x_{ref} - x(-)||_2}{||x_{ref}||_2} \quad (2)$$

$$399 \quad L2(+) = \frac{||x_{ref} - x(+)||_2}{||x_{ref}||_2} \quad (3)$$

400 In Figure 2 are reported the global relative variations in L2 of the principal maximum stress σ_1 and
401 principal minimum stress σ_3 caused by the variation of Young's Modulus in each Unit. All the
402 geometric configurations show variability less than 15%, with few exceptions within Unit B that have
403 variability over 30% (Fig. 2). It is worth noting that the spatial distribution of the major variations
404 seems to not significantly affect the final stress distributions, because: i) they are located near the
405 mesh borders (Fig. 3a, b); and, ii) when not at the mesh borders, the variations are limited to few %
406 (Fig. 3c, d). It means that changing the Young's Modulus of one order of magnitude produces
407 variation in FEM outputs distributed over a large domain and the change affecting the single nodes is
408 limited to few %.

409

410 *5.2 Homogeneous and not homogeneous lithology*

411 We carried out LISA simulations considering the effect of the gravitational loading on the
412 homogeneous and not homogeneous lithology on FEM outputs. In Figure 4 we reported a gravity
413 loading model for E-W cross-section of the CVC system. We first considered the homogeneous rock
414 composition composed by only andesitic lavas (Fig. 4a) and then by carbonates (Unit B), alluvial,
415 volcanoclastic and pyroclastic deposits (Units GF and VD; Fig. 4b). We analysed the principal
416 stresses σ_1 and σ_3 acting on the system, which correspond to the maximum and minimum stress at
417 a point, respectively.

418 Figure 4 shows the patterns of the minimum principal stress σ_3 (panels i-ii) and of the maximum
419 principal stress σ_1 (panels iii-vi), highlighting very slight differences between the homogeneous and
420 not homogeneous lithology cases. It is very important to stress that the x - z zero displacement
421 assigned at the bottom and the lateral boundaries of the domain created substantial artefacts in the
422 results (i.e. curved patterns of stress). The artefacts are also evident when considering σ_3 (panels i-ii)
423 where the boundary effect on x -axis is amplified by the presence of the upper free surface. For this
424 reason, the only area to be considered as unperturbed is the central part of the entire domain, and it
425 extends ca. 30 km horizontally and ca. 15 km vertically (within the blue contour in Fig. 4).

426

427 *5.3 Gravitational modelling using the inferred feeding system geometry*

428 We progressively add the elements of the conduit/feeding system of the CVC to FEM under the
429 effect of the gravitational loading. Three cross-section profiles (Figs. 5, 6) show increasing
430 complexity of the feeding system starting from a single magma chamber, passing to two magma
431 chambers, then adding the conduits, and, finally, considering the effects of faults.

432 Figure 5a describes the distribution of the minimum principal stress σ_3 (panel i) and the maximum
433 principal stress σ_1 (panel ii) at magma chamber overpressure of 10 MPa, showing how the insertion
434 of the pressurized magma chamber modifies the lithostatic stress. No significant differences in
435 magnitude and pattern of stresses are visible when having a magma chamber overpressure of 20 MPa
436 (Appendix 1a).

437 The addition of the shallow magma chamber significantly changes the values and pattern of both σ_3
438 and σ_1 (Fig. 5b). In particular, σ_3 and σ_1 stresses describe a typical inflation pattern produced by
439 overpressurised magma chamber(s) (Anderson, 1936; Gudmundsson, 2006), producing well-defined
440 stress arches of σ_3 (red dotted lines in Figs. 5bi) and divergent strong gradients of σ_1 , well developed
441 around the larger magma chamber (Fig. 5bii). Stress arch is a common phenomenon occurring in
442 continuous materials as response to applied pressure. It has been proved to have great influences on
443 the self-stabilization of soils or rock masses (Huang and Zhang, 2012), and may influence
444 mechanisms of caldera collapse (Holohan et al., 2015). Very slight differences in magnitude and
445 pattern of stresses appear when using 10 MPa (Fig. 5b) or 20 MPa of deep magma chamber
446 overpressure (Appendix 1b).

447 Figure 6 shows the effect of adding two conduits connecting the deep and shallow magma chambers.
448 It is evident how the insertion of the conduits in the feeding system of CVC dramatically changes the
449 stress distribution, with disappearance of the stress arch and an almost constant stress in the
450 computational domain except than on the tips of the deep magma chamber.

451

452 *5.4 Extensional field stress*

453 In order to explore the influence of the extensional far field stress on stress patterns (Fig. 1a), we run
454 simulations applying 5 MPa of extensional stress to the FEM domain, which is a typical low value for

455 rift zones (Turcotte and Schubert, 2002; Moeck et al., 2009; Maccaferri et al., 2014; Sulpizio and
456 Massaro, 2017; Fig. 7).

457 In the case of a single magma chamber (with 10 MPa overpressure; Fig. 7, panels i-ii), the addition of
458 extensional far field stress reduces the confinement effect due to the no displacement condition
459 imposed along the x - z directions (plane strain approximation). The effect of the extensional field
460 stress on double magma chamber configuration (with 10 MPa overpressure in the deep chamber and
461 5 MPa in the shallower one) produces slight changes in stress magnitude and pattern for both σ_3 and
462 σ_1 (Fig. 7, panels iii-iv) with respect to Figure 5b. The same applies also for the complete feeding
463 system configuration, in which the attrition of the far field stress slightly changes the intensity of the
464 stresses and patterns (Fig. 7, panels v-vi). Using 20 MPa overpressure in the deep magma chamber
465 does not significantly affect the model outputs (Appendix 2).

466

467 *5.5 Faults bordering the Colima Rift*

468 In order to reproduce the effect of faults bordering the Colima Rift on the final feeding system
469 configuration, we added two damage zones by degrading the elastic properties of a volume of rock
470 mass.-The insertion of the two zones of weakness does not alter significantly the stress distribution
471 observed in Figures 7v and 7vi, with only reduction of both σ_1 and σ_3 values in the surroundings of
472 the damage zones (Figs. 7vii and 7viii). The different distance of the two damage zones to the
473 feeding system (especially the deep magma chamber) produces a small asymmetry in both σ_1 and σ_3
474 patterns with respect to simulations without damage zones (Figs. 7v-viii).

475

476

477 **6 Discussions**

478

479 *6.1 FEM analysis at increasing geological details*

480 The presented FEM model of the CVC highlighted some important characteristics of crustal stress
481 distribution at changing geological constraints used as input conditions (Spica et al., 2014, 2017;
482 Massaro et al., 2018). Although the results have to be considered as a first order approximation, the

483 changes in stress distribution are evident and useful for the understanding of limitations and
484 advantages of FEM.

485 Under the assumptions of plane strain, gravitational loading, and overpressured magma chambers and
486 dykes, the use of homogeneous or not homogeneous lithology for FEM provides negligible effects in
487 stress intensity and pattern (Fig. 4). This is because the upper Units (Units FC, VD, GF; Table 2)
488 represent only a limited part of the simulated domain, which in the remaining part results entirely
489 composed of the assumed homogeneous basement (Unit B; Table 2). This does not mean that the
490 influence of the upper Units may be still negligible using smaller scales of the simulated domain.

491 Analysing the FEM outputs with the single magma chamber, it emerges how the overpressures, ΔP ,
492 only limited the effects of gravitational loading. The use of a dual magma chamber geometry better
493 describes the inflation induced by overpressure within magma chambers, with the formation of the
494 stress arch in the minimum compressive stress σ_3 plot. It is important to highlight that for both single
495 and dual magma chamber models, the change of internal overpressure from 10 to 20 MPa slightly
496 changes the magnitude of the stress but not their general patterns (Appendix 1-2).

497 The presence of dykes in the magma feeding system dramatically change the σ_3 and σ_1 patterns (Fig.
498 6). Indeed, they become quite homogeneous throughout the computational domain, with the only
499 exception of sidewall effects induced by the zero displacement conditions, already discussed in
500 Figure 4.

501 The addition of extensional field stress of 5 MPa reduces the sidewall effects and produces an almost
502 homogeneous stress distribution in the upper part of the computational domain, above the top of the
503 deep magma chamber. This, along with the additional inclusion of the damage zones introduced to
504 mimic the effects of CG faults, describes a close to equilibrium volcanic system, in which volcanic
505 overpressure and lithostatic stress almost equilibrate each other (Sulpizio et al., 2016).

506

507 *6.2 Some implication of the stress state of the CVC inferred from FEM*

508 The results obtained with the insertion of the full feeding system and far field stress on the FEM
509 highlight an almost homogeneous stress distribution in the CVC area. This means that the shape of
510 the dual magma chamber feeding system model and far field stress provide a stable geometry, which
511 limits the stress changes to few MPa. All the large stress variations are located at the tips of the

512 magma chambers, as expected for pressurized or under-pressurized cavities in the lithosphere (Martì
513 and Geyer, 2009). This means that the whole feeding system is in a quasi-equilibrium state, and, as
514 an example, any overpressure created by input of new magma is adjusted by increasing the magma
515 chamber volume or erupting at the surface. Even if we consider the scenario of complete emptying
516 the upper conduit and part of the shallow magma chamber, as occasionally occurred during the past
517 sub-Plinian and Plinian eruptions (Luhr et al., 2002; Saucedo et al., 2010; Massaro et al., 2018), this
518 would result in the restoration of the stress arch, which is still a stable stress configuration. Even the
519 complete emptying of the shallow magma chamber probably would be ineffective for triggering a
520 large collapse (caldera forming) of the feeding system. This latter event would be possible only if a
521 large depressurization of the deeper magma chamber would occur, but it implies the eruption of tens
522 to hundreds of km³ of magma, which seems not very likely provided the current stress distribution in
523 CVC.

524 Beside and beyond the limitations due to the first order approximation of the FEM analysis, other
525 sources of uncertainties in the discussion about present and future stress state of the CVC come from
526 not considering gravity-driven processes, such as volcano spreading due to plastic deformation of the
527 GF Unit (Norini et al., 2010, 2019) or pressurization of the shallower conduit (Massaro et al., 2018),
528 and detailed regional tectonics (Norini et al., 2010, 2019). The effect of the two fault systems
529 bordering the Colima Rift were simulated by degrading the mechanic properties of rocks in an area of
530 about 1 km width up to a depth of 10 km. Although the effects are negligible at the scale of the
531 computational domain, it cannot be excluded some local significant effects that cannot be resolved
532 using the described approach.

533

534 **7 Conclusions**

535 The increasing details of geological and geophysical data to FEM simulation at Colima Volcanic
536 Complex (Mexico) showed the importance of using the most accurate input data in order to have
537 reliable outputs. In particular, the data here presented highlighted how the use of simplified models
538 produces unreliable outputs of the stress state of the volcano subsurface.

539 Beside and beyond the results obtained by analysing the influence of detailing geological and
540 geophysical data, the FEM of CVC confirms the close to equilibrium state of the volcano, which is
541 the expected stress distribution induced by a feeding system directly connected to the surface.
542 This means that any overpressure created by input of new magma is adjusted within the feeding
543 system, sometimes triggering eruptions. The complete emptying the upper conduit and part of the
544 shallow magma chamber, as occasionally occurred in the past, originating sub-Plinian and Plinian
545 eruptions, would result in the restoration of the stress arch, which is still a stable stress configuration.
546 Descends that large magnitude, caldera forming eruptions are possible only if the bigger deep magma
547 chamber is also involved and significantly emptied during an eruption.

548

549 **Appendices**

550

551 **Appendix 1**

552 E-W gravitational modelling of the CVC domain (stratified lithology) for all configurations
553 investigated. The magnitude and pattern of the principal stress account for a) single magma chamber
554 model (number of nodes: 4426); b) dual magma chamber model (number of nodes: 4161); c) dual
555 magma chamber with conduits model (number of nodes: 3737). The dimension of the deep magma
556 chamber: $2a = 14$ km and $2b = 3.6$ km at 15 km of depth; shallow magma chamber: $2a = 3.5$ km and
557 $2b = 2$ km at 6 km. The magmatic overpressure is 20 MPa for the deep chamber, and 5 MPa for the
558 shallower. Black dotted lines highlight the passage from different stress values. Note that the scale of
559 stress values are different for each panel in order to maximise the simulation details.

560

561 **Appendix 2**

562 E-W gravitational modelling of the CVC domain (stratified lithology) considering a far extensional
563 stress field of 5 MPa for all configurations investigated. The magnitude and pattern of the principal
564 stress account for a) single magma chamber model model (number of nodes: 4426); b) dual magma
565 chamber model (number of nodes: 4161); c) dual magma chamber with conduits model (number of
566 elements: 3737). The dimension of the deep magma chamber: $2a = 14$ km and $2b = 3.6$ km at 15 km
567 of depth; shallow magma chamber: $2a = 3.5$ km and $2b = 2$ km at 6 km. The magmatic overpressure
568 is 20 MPa for the deep chamber, and 5 MPa for the shallower. Black dotted lines highlight the
569 passage from different stress values. The red arrows indicate the direction of the applied far field

570 stress. Note that the scale of stress values are different for each panel in order to maximise the
571 simulation details.

572

573 **Code/Data Availability**

574 The LISA code is available at <https://lisafea.com/>.

575

576 **Author's contribution**

577 SM, RS, AC, GN and GG conceived the study. SM and RS wrote the bulk of the manuscript with the
578 input of all the co-authors. SM and GL compiled the numerical simulations and formulated the
579 adopted methodology. MP and SM carried out the sensitivity analysis. RS, AC, SM, GN, GG, LC,
580 GL, MP and AG worked on the interpretation of the results.

581

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L048 **Table 1** - Element types used in LISA analysis considering the final conduit feeding system
L049 configuration – Fig.1d, panel vi)

L050	<i>E-W cross-section (a-a')</i>		Element Type	Elements	Nodes
L051	FC	Fuego de Colima	quad4-tri3	372	384
L052	VD	Volcanic Deposits	quad4-tri3	245	273
L053	GF	Graben Fill	quad4-tri3	456	338
L054	B	Basament	quad4-tri3	3088	2907
L055	CG	Colima graben	quad4-tri3	48	71
L056	Total Elements: 4209				

L057 **Table 2** - Rock mass and mechanical properties of the geological Units used in the finite-element
L058 model (from Norini et al., 2010, 2019).
L059

Acronym	Model Unit	Rock Type	Density (kg/m ³)	Young's Modulus (MPa)	Poisson's ratio ν
FC	Fuego de Colima	Andesitic lavas and pyroclastic deposits forming the Paleofuego-Fuego de Colima	2242	1.4×10^3	0.30

		volcano			
VD	Volcaniclastic deposits	Pyroclastic and epiclastic deposits covering the southern flank of the CVC	1539	1.7×10^3	0.32
GF	Graben Fill	Quaternary alluvial, colluvial, lacustrine deposits filling the graben	1834	1.5×10^3	0.35
B	Basement	Cretaceous limestones and intrusive rocks forming the bed-rock underlying the CVC	2650	3.6×10^4	0.30

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L062 Figures Captions

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L064 **Fig. 1** (a) Morphotectonic map of the Colima Volcanic Complex (NC=Nevado de Colima volcano;
L065 FC=Fuego de Colima volcano) and Colima Rift with the main tectonic and volcano-tectonic
L066 structures (NCG =Northen Colima Graben; CCG= Central Colima Graben, from Norini et al., 2019).
L067 In the inset, the location of the Colima Volcanic Complex (CVC) within the Trans-Mexican Volcanic
L068 Belt (TMVB) is shown in the frame of the subduction-type geodynamic setting of Central America
L069 (from Davila et al., 2019); (b) general sketch of the geometrical configurations used in LISA; (c)
L070 example of mesh of the investigated area for the dual magma chamber model with conduits (case v in
L071 panel (b), considering zero-displacement along the bottom and left and right sides. Note that for case
L072 (vi) in panel (b) the zero-displacement is removed from the lateral sides; (d) sketch of the Fuego de
L073 Colima feeding system composed of a 15 km-deep magma chamber connected to surface via a 6 km-
L074 deep magma chamber and dykes. ΔP_{chs} and ΔP_{chd} are the magmatic overpressures in the shallow
L075 and deep chambers, respectively (modified from Massaro et al., 2019).

L076

L077 **Fig. 2** Results of the sensitivity analysis carried out on the Young's Modulus variations within each
L078 rock layer of the domain considering different configurations (stratified substratum model – nodes:
L079 4426; single magma chamber model – nodes: 4426; dual magma chamber model – nodes: 4161; dual
L080 magma chamber with conduits model – nodes: 3737). For each geological Unit (B, FC, GF, VD), the
L081 relative global variation in L_2 (%) is provided for σ_1 and σ_3 . The $x(-)$ and $x(+)$ vectors indicate the
L082 Young's Modulus variation by an order of magnitude with respect to x_{ref} vector, containing the stress
L083 values calculated by using the values of material's properties indicated in Table 2.

L084

L085 **Fig. 3** Spatial variation (%) of the L_2 norm's components at varying Young's Modulus for selected
L086 cases of Units B and VD: (a) Unit B in the stratified substratum model (nodes: 4426); (b) Unit B in
L087 the single magma chamber model (nodes: 4426); (c) Unit B in the dual magma chamber model

L088 (nodes: 4161); (d) Unit VD in the dual magma chamber with conduits model (nodes: 3737). Symbols
L089 $x(-)$ and $x(+)$ have the same meaning of Figure 2.

L090

L091 **Fig. 4** E-W gravitational modelling of the CVC domain. The scale of the mesh is expressed in Unit
L092 of Design (1 UD = 1 km). The domain extends 60 km along the x -axis, and 30 km along the z -axis.
L093 The number of nodes used in the mesh is set to 4426. The magnitude and pattern of the principal
L094 stresses (dotted black lines) are reported for (a) the homogeneous stratigraphy (Unit FC = andesitic
L095 lavas and pyroclastic deposits) and for (b) the not homogeneous stratigraphy (Unit FC; Unit B =
L096 Cretaceous limestones and intrusive rocks forming the bed-rock underlying the CVC; Unit GF =
L097 Quaternary alluvial, colluvial, and lacustrine deposits filling the graben; Unit VD = volcanoclastic
L098 deposits covering the southern flank of the CVC). The blue line contours the unperturbed part of the
L099 domain, which extends ca. 30 km horizontally and ca. 25 km vertically. Note that the scale of stress
L100 values is the same for the all simulations.

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L103 **Fig. 5** E-W gravitational modelling of the CVC domain with a not homogeneous stratigraphy. The
L104 magnitude and pattern of the principal stresses are reported for (a) the single magma chamber
L105 represented by a magma chamber ($2a = 14$ km and $2b = 3.6$ km) at 15 km of depth, and (b) the dual
L106 magma chamber model composed of a 15 km-deep magma chamber ($2a = 14$ km and $2b = 3.6$ km)
L107 and a shallow 6 km-deep one ($2a = 3.5$ km and $2b = 2$ km). The magma chambers are not connected.
L108 The magmatic overpressures are set to 10 and 5 MPa for the 15 km-deep and 6 km-deep magma
L109 chambers, respectively. The number of nodes is set to 4426 and 4161 for the single and dual magma
L110 chamber models, respectively. Black dotted lines highlight the passage from different stress values.
L111 The red dotted line in panel (b-i) indicates the formation of the stress arch. Note that the scale of
L112 stress values are different for each panel in order to maximise the simulation details.

L113

L114 **Fig. 6** E-W gravitational modelling of the CVC domain with a not homogeneous stratigraphy
L115 accounted for a dual magma chamber system connected by dykes via surface (deep magma chamber,
L116 $2a = 14$ km and $2b = 3.6$ km at 15 km of depth; shallow magma chamber, $2a = 3.5$ km and $2b = 2$ km
L117 at 6 km of depth). The magnitude and pattern of the principal stresses are shown. The number of
L118 nodes used is set to 3737. The magmatic overpressures are set to 10 and 5 MPa for the 15 km-deep
L119 and 6 km-deep magma chambers, respectively. The black dotted lines in panel (ii) highlight the
L120 passage from different stress values. Note that the scale of stress values are different for each panel in
L121 order to maximise the simulation details.

L122

L123 **Fig. 7** E-W gravitational modelling of the CVC domain with a not homogeneous stratigraphy

l124 considering the extensional field stress. The magnitude and pattern of the principal stresses are shown
l125 for the single magma chamber model (panels i-ii), the dual magma chamber model (panels iii-iv), the
l126 dual magma chamber with conduits model (panels v-vi-vii-viii). Note that in panel vii-viii the faults
l127 bordering the Colima graben are shown. For all configurations an extensive far-field stress of 5 MPa
l128 is applied at the lateral boundaries of the domain. In panels vii-viii the additional effect of the local
l129 extensive field is simulated using a reduced values of material's properties (Table 2). The magmatic
l130 overpressures are set to 10 and 5 MPa for the 15 km-deep and 6 km-deep magma chambers,
l131 respectively. Black dotted lines highlight the passage from different stress values. The red arrows
l132 indicate the direction of the applied far field stress. Note that the scale of stress values are different
l133 for each panel in order to maximise the simulation details.

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l142 **Figures**

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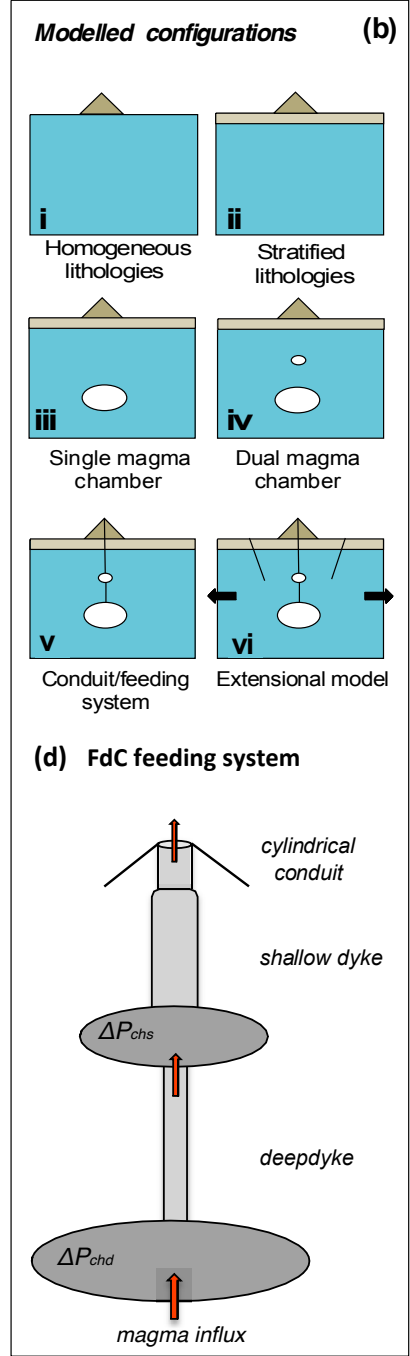
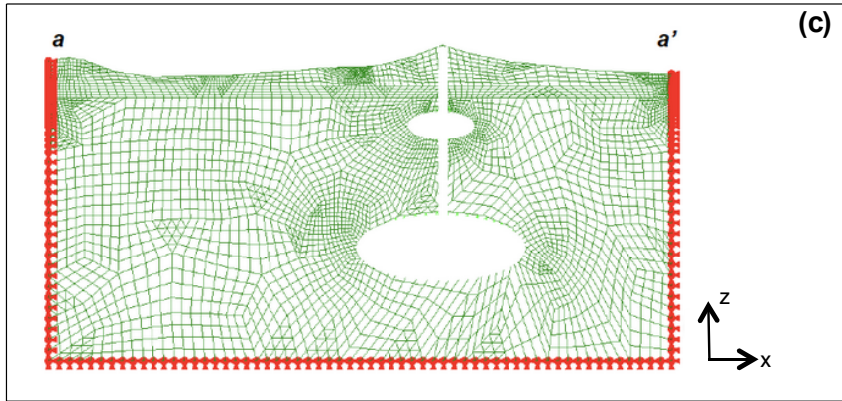
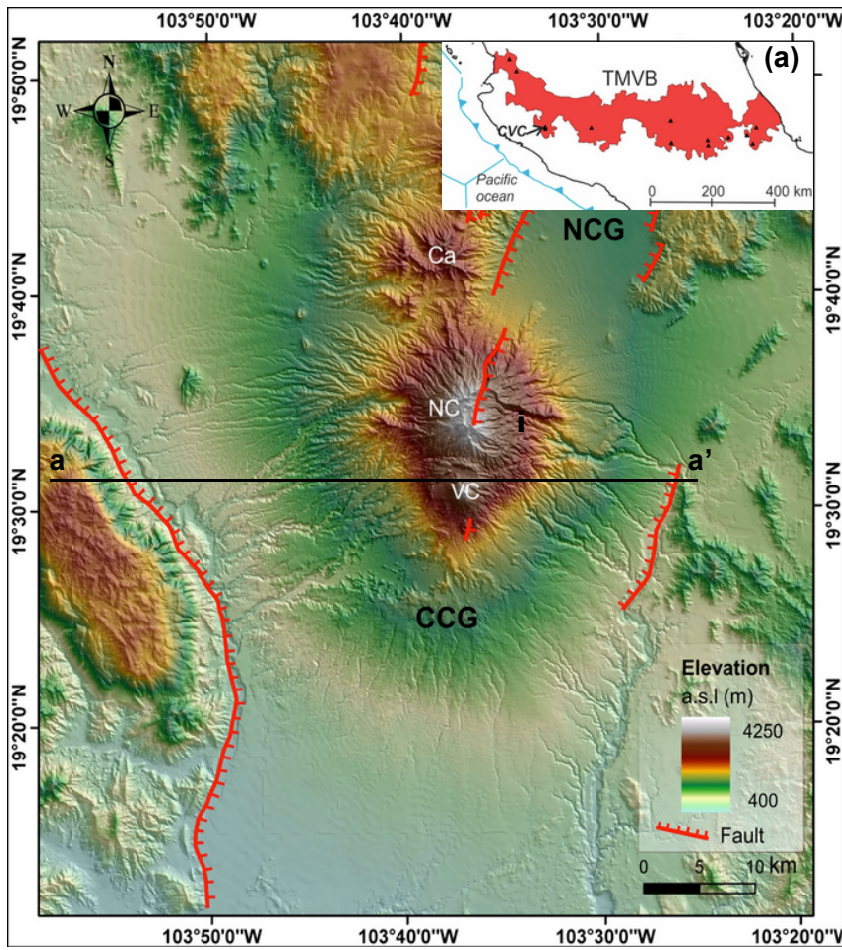
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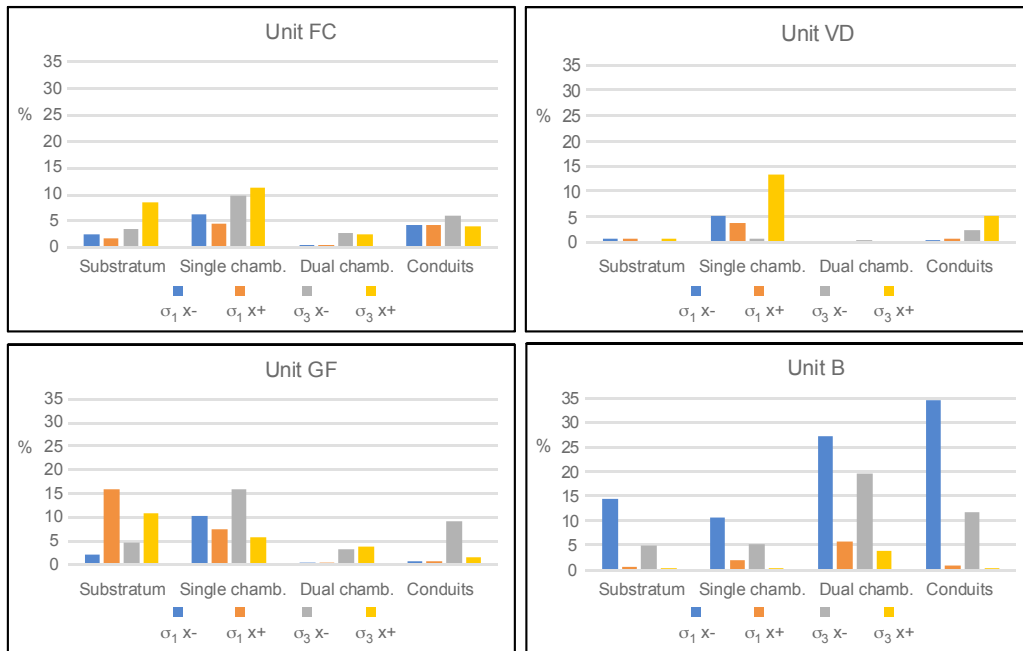
L152 Figure 1



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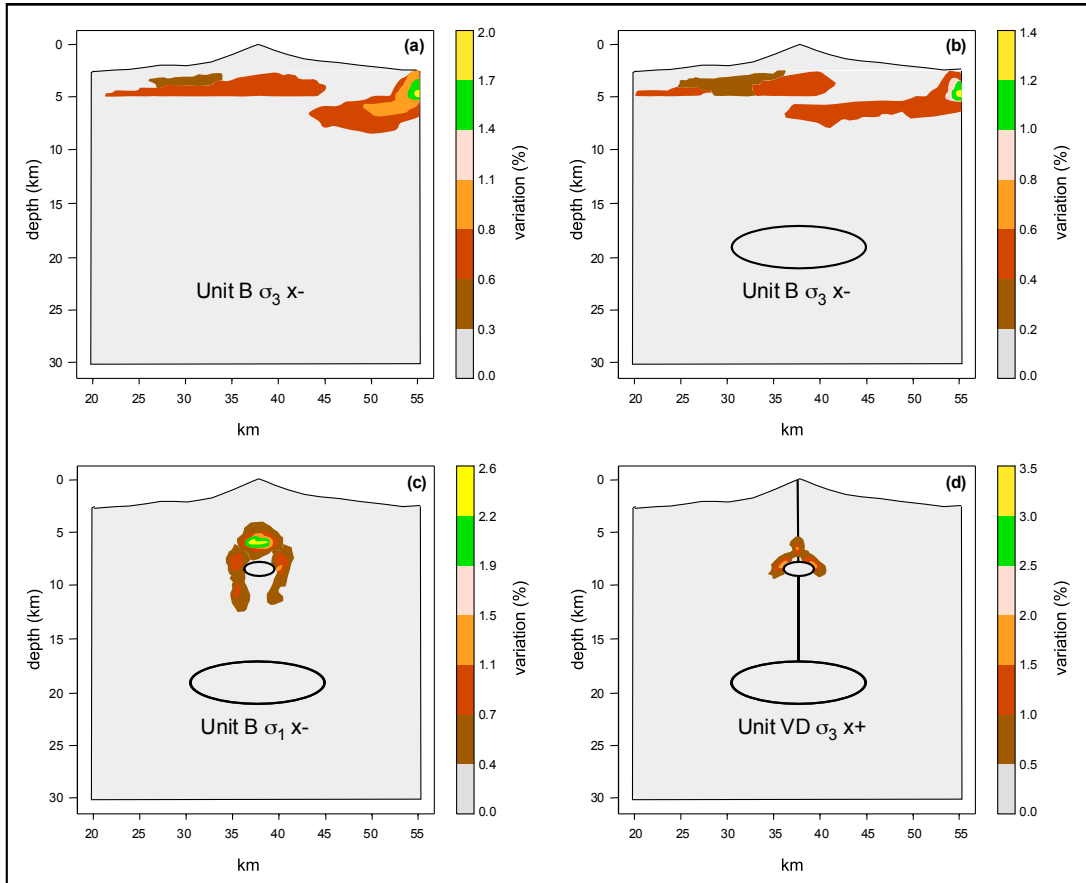
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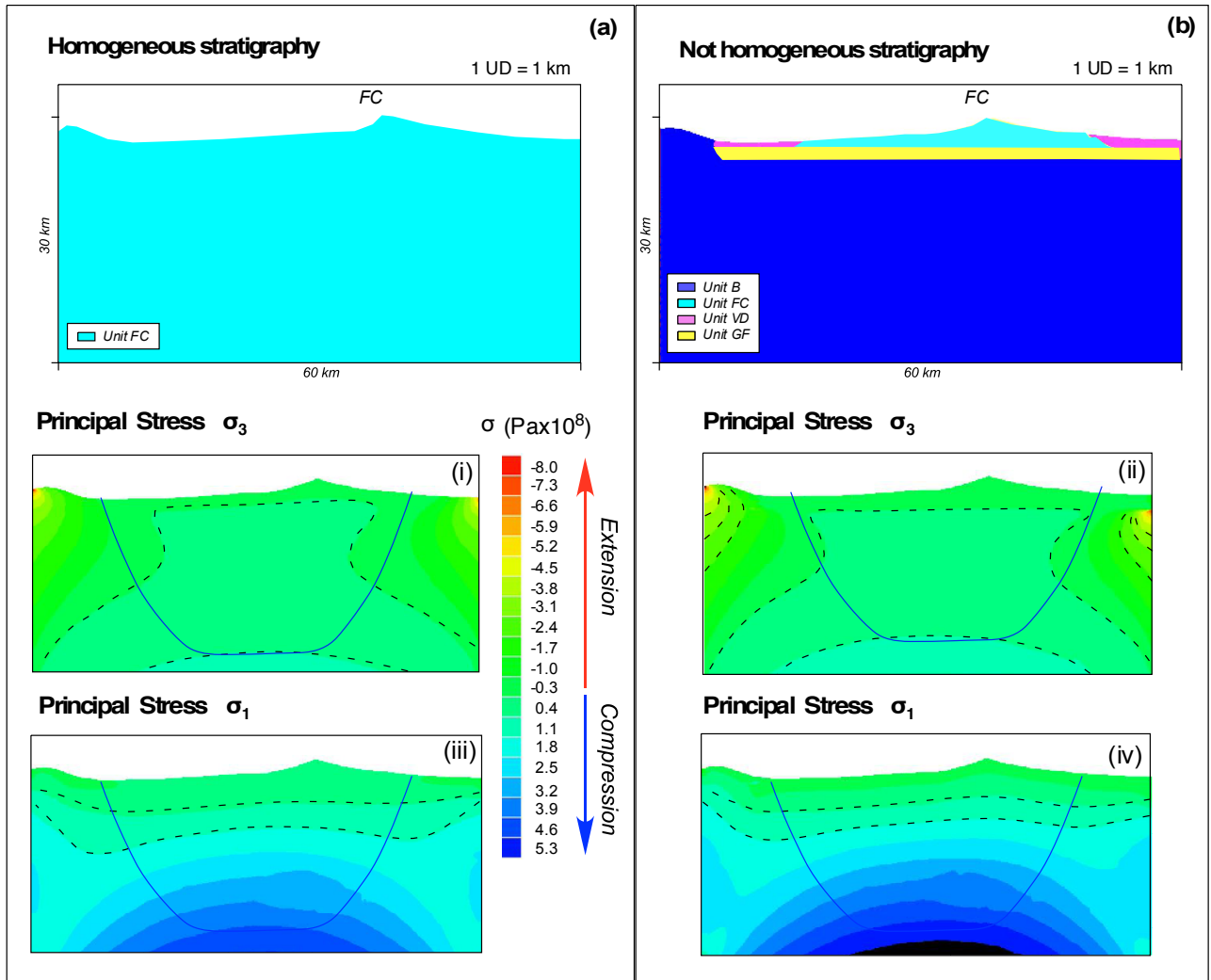


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L161 Figure 4

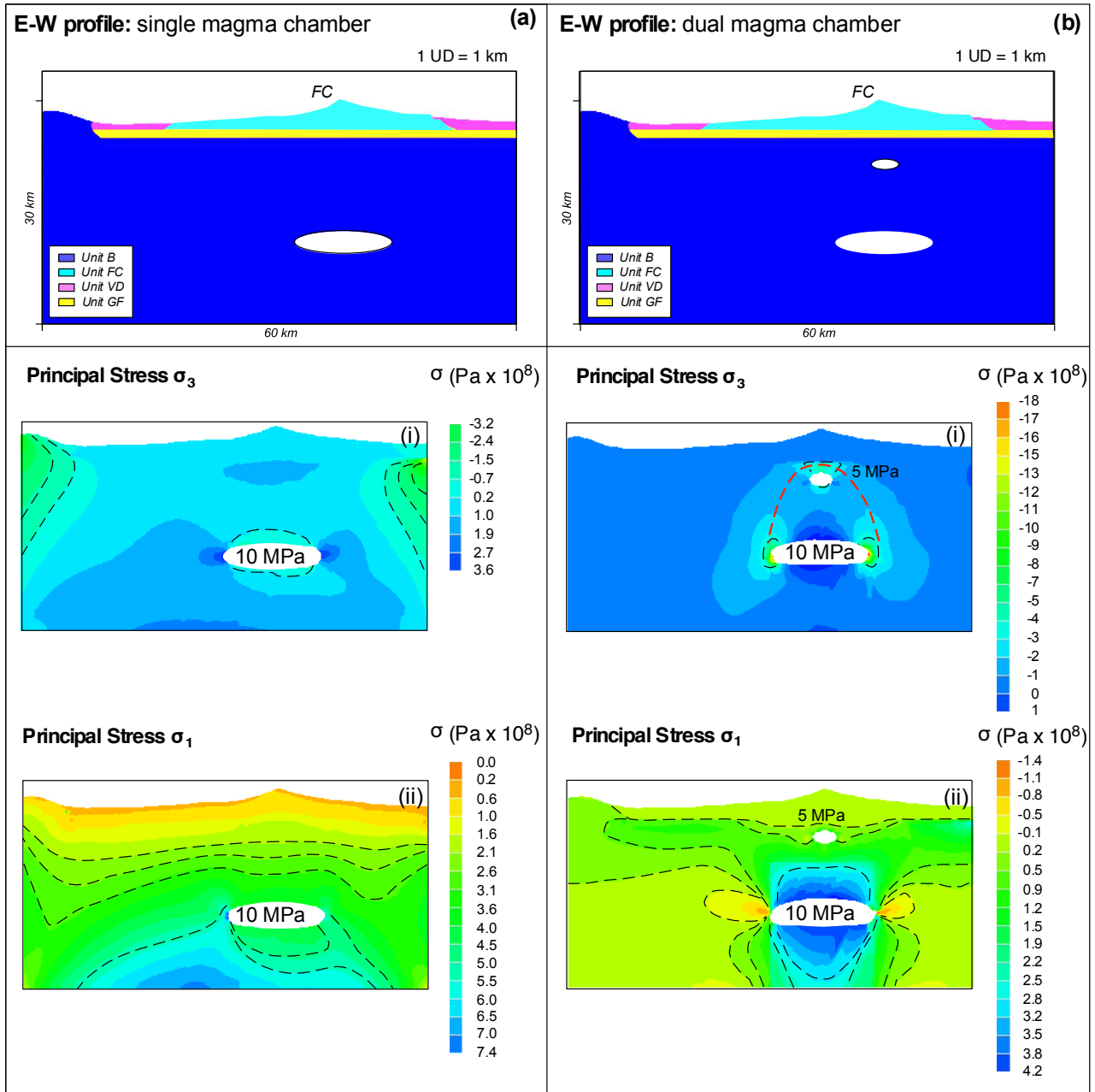


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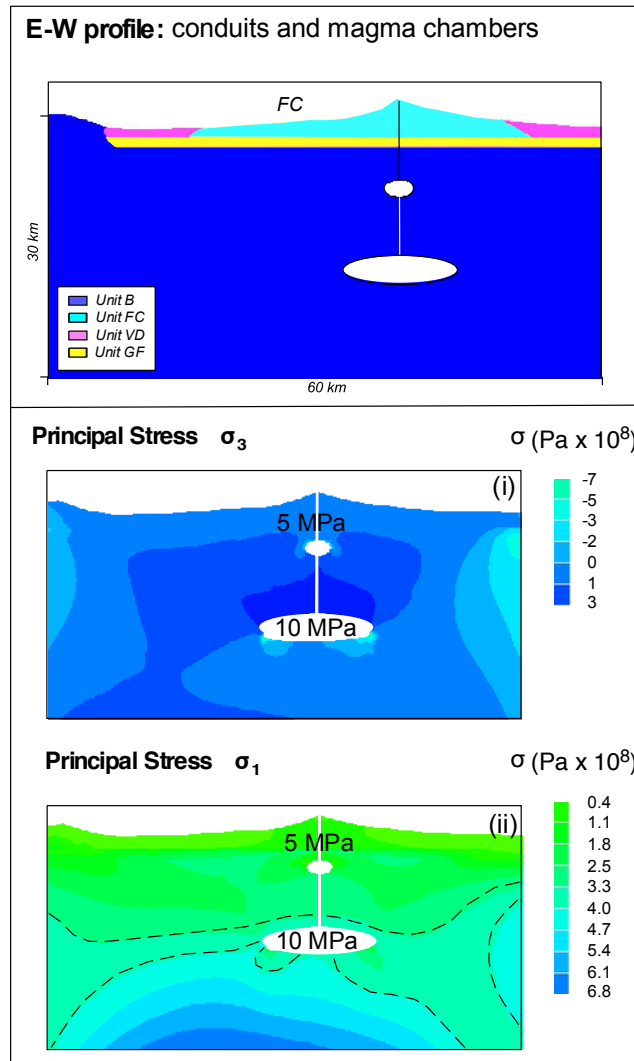
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L171 Figure 6



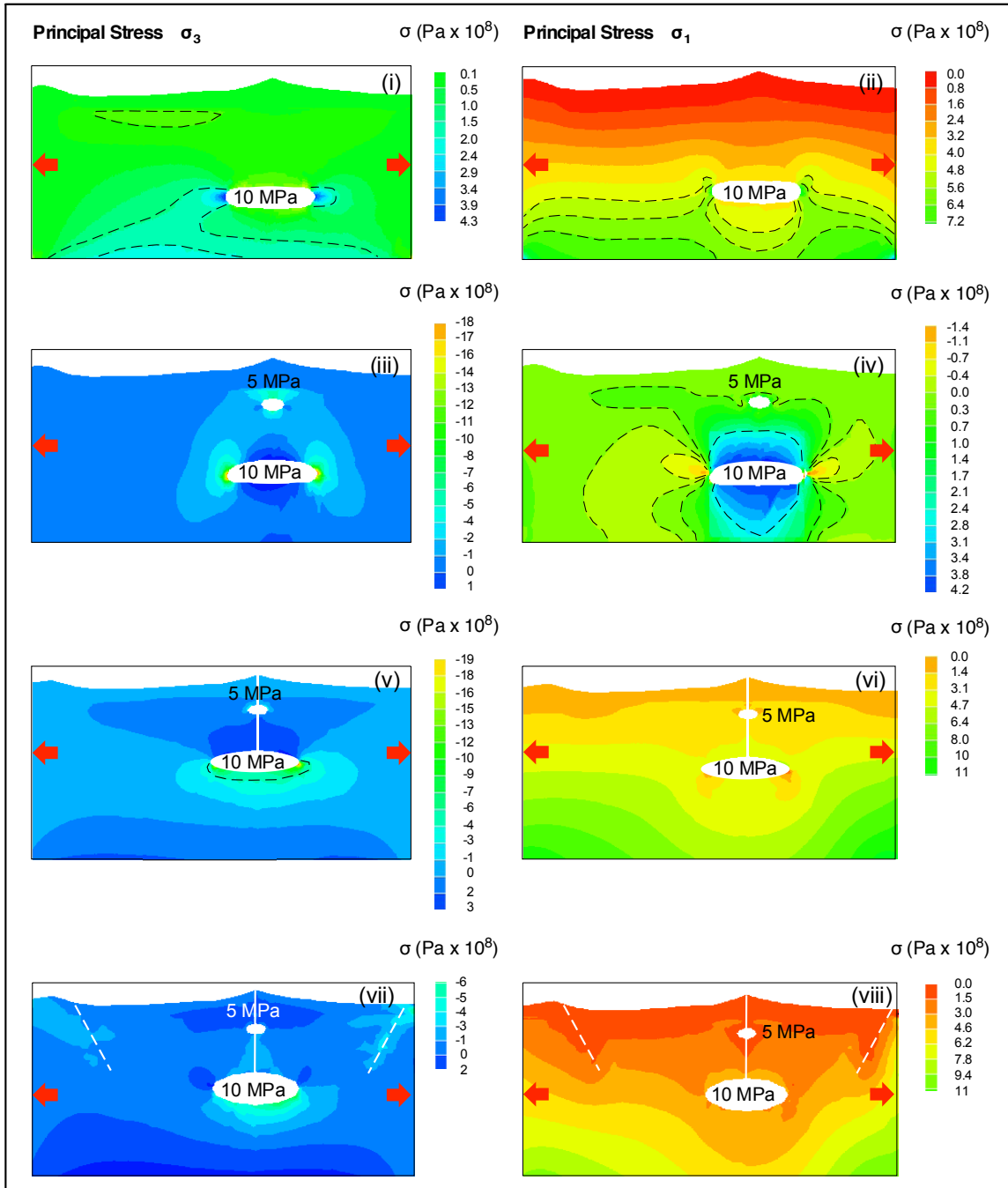
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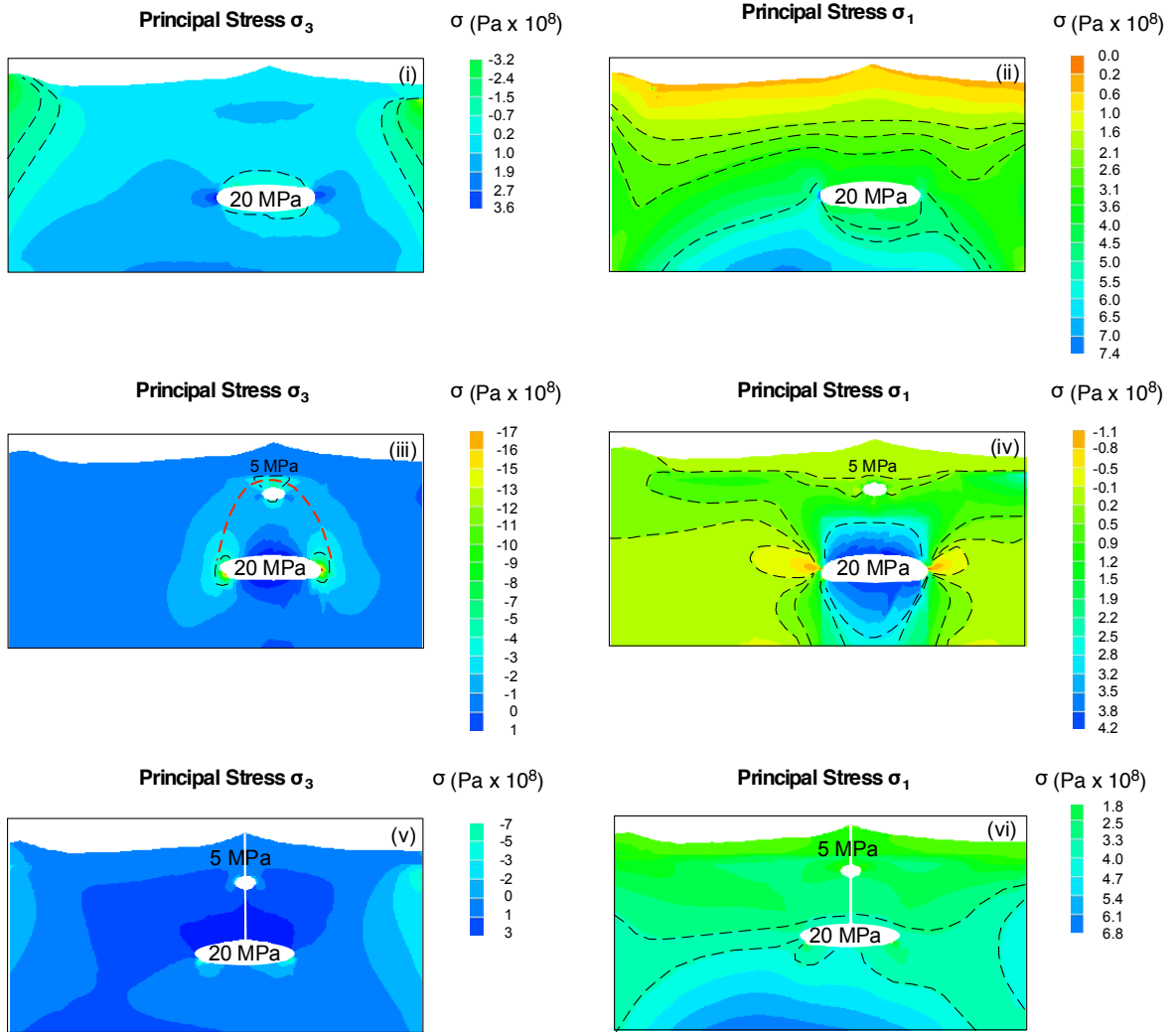
L176 Figure 7



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Appendix 1



Appendix 2

