



1 Transient conduit permeability controlled by a shift between compactant shear and dilatant

2 rupture at Unzen volcano (Japan)

3

- 4 Lavallée Yan^{1*}, Miwa Takahiro², Ashworth James D.¹, Wallace Paul A.^{1,3}, Kendrick Jackie E.^{1,4}, Coats
- 5 Rebecca¹, Lamur Anthony¹, Hornby Adrian⁵, Hess Kai-Uwe.⁶, Matsushima Takeshi⁷, Nakada Setsuya⁸,
- 6 Shimizu Hiroshi⁸, Ruthensteiner Bernhard⁹, Tuffen Hugh¹⁰

7

- 8 ¹Earth, Ocean and Ecological Sciences, University of Liverpool, Liverpool, United Kingdom
- ² Earthquake Research Department, National Research Institute for Earth Science and Disaster
 Resilience (NIED), Tsukuba, Japan
- ³ Department of Geosciences, Environment and Society, Université Libre de Bruxelles, Brussels,
 Belgium
- 13 ⁴ Geosciences, University of Edinburgh, Edinburgh, United Kingdom
- 14 ⁵ Earth and Atmospheric Sciences, Cornell University, United States of America
- 15 ⁶ Earth and Environmental Sciences, Ludwig-Maximilians University of Munich, Germany
- 16 ⁷ Institute of Seismology and Volcanology, Faculty of Sciences, Kyushu University, Shimabara,
- 17 Nagasaki, Japan
- ⁸ National Research Institute for Earth Science and Disaster Resilience, Tennodai, Tsukuba, 3050006, Japan
- ⁹ Staatliche Naturwissenschaftliche Sammlungen Bayerns (SNSB), Zoologische Staatssammlung
 München, München, Germany.
- 22 ¹⁰ Earth Sciences, University of Lancaster, United Kingdom
- 23 *ylava@liverpool.ac.uk

24

25 ABSTRACT

26 The permeability of magma in shallow volcanic conduits controls the fluid flow and pore pressure 27 development that regulates gas emissions and the style of volcanic eruptions. The architecture of the 28 permeable porous structure is subject to changes as magma deforms and outgasses during ascent. Here, 29 we present a high-resolution study of the permeability distribution across two conduit shear zones (marginal and central) developed in the dacitic spine that extruded towards the closing stages of the 30 31 1991-1995 eruption at Unzen volcano, Japan. The marginal shear zone is approximately 3.2 m wide and exhibits a 2-m wide, moderate shear zone with porosity and permeability similar to the conduit 32 33 core, transitioning into a ~1-m wide, highly-sheared region with relatively low porosity and 34 permeability, and an outer 20-cm wide cataclastic fault zone. The low porosity, highly-sheared rock 35 further exhibits an anisotropic permeability network with slightly higher permeability along the shear 36 plane (parallel to the conduit margin) and is locally overprinted by oblique dilational Riedel fractures. 37 The central shear zone is defined by a 3-m long by ~9-cm wide fracture ending bluntly and bordered





by a 15-40 cm wide damage zone with an increased permeability of ~3 orders of magnitude; directional
 permeability and resultant anisotropy could not be measured from this exposure.

40 We interpret the permeability and porosity of the marginal shear zone to reflect the evolution of 41 compactional (i.e., ductile) shear during ascent up to the point of rupture, estimated by Umakoshi et al. (2008), at ~500 m depth. At this point the compactional shear zone would have been locally overprinted 42 43 by brittle rupture, promoting the development of a shear fault and dilational Riedel fractures during 44 repeating phases of increased magma ascent rate, enhancing anisotropic permeability that channels fluid 45 flow into, and along, the conduit margin. In contrast, we interpret the central shear zone as a shallow, 46 late-stage dilational structure, which partially tore the spine core with slight displacement. We explore 47 constraints from monitored seismicity and stick-slip behaviour to evaluate the rheological controls, 48 which accompanied the upward shift from compactional toward dilational shear as magma approached 49 the surface, and discuss their importance in controlling the permeability development of magma 50 evolving from overall ductile to increasingly brittle behaviour during ascent and eruption.

51

52 1. Introduction

53 1.1. Outgassing pathways and volcanic eruptions

54 The style and timing of activity exhibited during a volcanic eruption are strongly influenced by 55 the presence and mobility of volatiles in magma (Sparks, 1997; Woods and Koyaguchi, 1994) and 56 surrounding conduit wallrock (Jaupart and Allègre, 1991). During magma ascent, volatiles are exsolved 57 into gas bubbles (Navon et al., 1998; Sparks, 2003) as their solubility decreases with decompression 58 (Liu et al., 2005), crystallisation (Tait et al., 1989), and heat generated by crystallisation (Blundy et al., 59 2006) and shear (Lavallée et al., 2015). This causes the accumulation of pressurised fluids in vesicles 60 that charges ascending magma, which, if sufficient may lead to fragmentation (Mueller et al., 2008; 61 Alidibirov and Dingwell, 1996) and an explosive eruption (Sahagian, 1999). The development of a permeable network governs outgassing (Edmonds et al., 2003), pore pressure release (Mueller et al., 62 63 2005), and eruptive cyclicity (Michaut et al., 2013), thereby reducing the potential for explosive activity (Klug and Cashman, 1996) and encouraging effusion (Edmonds and Herd, 2007; Eichelberger et al., 64 65 1986; Degruyter et al., 2012). Lava dome eruptions-the topic of this study-commonly switch between effusive and explosive modes of activity due to this competition between permeability, pore 66 67 fluid pressure and the structural integrity of magma (Melnik and Sparks, 1999; Calder et al., 2015; 68 Cashman et al., 2000; Castro and Gardner, 2008; Edmonds et al., 2003; Lavallée et al., 2013; Lavallée 69 et al., 2012; Sparks, 1997; Holland et al., 2011; Kendrick et al., 2016; Platz et al., 2012). Considering 70 the water solubility-pressure relationships in magmas (Zhang, 1999), permeability-porosity 71 relationships in magma (Westrich and Eichelberger, 1994) and eruptive patterns (Edmonds et al., 2003), 72 it has been suggested that much of the outgassing during lava dome eruptions occurs in the upper few 73 kilometres of the conduit (Westrich and Eichelberger, 1994; Edmonds et al., 2003). This observation is 74 corroborated by rapid shallowing of seismicity leading to explosions (e.g., Rohnacher et al., 2021) and 75 the existence of shallow long-period seismic signals resulting from resonance in fractures and faults 76 (Chouet, 1996; Matoza and Chouet, 2010) as fluids are channelled to the surface (Holland et al., 2011; Kendrick et al., 2016; Gaunt et al., 2014; Nakada et al., 1995; Newhall and Melson, 1983; Pallister et 77 78 al., 2013b; Sahetapy-Engel and Harris, 2009; Sparks, 1997; Sparks et al., 2000; Edmonds et al., 2003; 79 Varley and Taran, 2003; Stix et al., 2003). Therefore understanding the evolution of the permeable network during eruptive shearing is central to constrain the evolution of the magmatic system in the 80 81 shallow crust (Blower, 2001).





82 Close examination of the architecture of shallow dissected conduits and structures in vent-83 proximal silicic lava exposes complex shearing histories that would impact the permeable porous 84 network of erupting magma. These structures reveal porosity contrasts through the lavas, and strain 85 localisation near the conduit margins is commonly identified via the presence of flow bands and variably 86 porous shear zones with a spectrum of configurations (Gaunt et al., 2014; Kendrick et al., 2012; 87 Kennedy and Russell, 2012; Pallister et al., 2013a; Smith et al., 2001; Stasiuk et al., 1996; Tuffen and 88 Dingwell, 2005); features that are preserved to differing extents in crystal-poor and crystal-rich magmas 89 (Calder et al., 2015; Lavallée and Kendrick, 2021). For example, crystal-poor obsidian in dissected 90 conduits and dykes commonly exhibits marginal flow bands, showing alternation between glassy, finely 91 crystalline and microporous bands (Gonnermann and Manga, 2007). Flow bands also occur as variably 92 sintered, cataclastic breccia layers, resulting from fracture and healing cycles (Tuffen and Dingwell, 93 2005; Tuffen et al., 2003), and as variably sintered tuffisite layers, resulting from fragmentation and 94 entrapment of fragments into narrow fractures (Castro et al., 2012; Heiken et al., 1988; Kendrick et al., 95 2016; Kolzenburg et al., 2012). Exposed crystal-poor conduits, dykes and domes are commonly dense, 96 as the porous network may easily collapse (unlike crystal-rich lavas; e.g., Ashwell et al., 2015). The 97 collapse of the porous network occurs as eruptions wane and pore pressure is insufficient to counteract 98 surface tension and local magmastatic and lithostatic stresses (Wadsworth et al., 2016a; Kennedy et al., 99 2016), a process which hinders interpretation of the syn-eruptive permeable structure of crystal-poor 100 magma from the study of large-scale relict formations. Studies of erupted crystal-poor pumices (which 101 quench rapidly) help provide constraints on the extent of magma permeability at the point of 102 fragmentation (Wright et al., 2006), but the task of reconstructing the permeable architecture of an entire 103 conduit from these pyroclasts is challenging (Dingwell et al., 2016), further complicated by postfragmentation vesiculation (Browning et al., 2020) and vesicle relaxation (Rust and Manga, 2002), and 104 105 so remains to be attempted systematically.

106 Crystal-rich volcanic rocks (the subject of this study) expose a wider range of permeable porous 107 structures (Farquharson et al., 2015; Mueller et al., 2005; Klug and Cashman, 1996; Lamur et al., 2017; 108 Kushnir et al., 2016). For instance, dacitic volcanic spines extruded in 2004-08 at Mount St. Helens 109 (USA) and in 1994-95 Unzen volcano (Japan) reveal the presence of a complex 'damage halo' near the 110 conduit margin (Calder et al., 2015; Gaunt et al., 2014; Pallister et al., 2013a; Smith et al., 2001; 111 Kendrick et al., 2012; Wallace et al., 2019). These structures frequently share common characteristics, 112 with magma being increasingly sheared and damaged near the conduit margin, defined by a cataclastic 113 fault zone, adjacent to a brecciated wall-rock. A permeability study of the shear zone at Mount St. 114 Helens showed increased porosity and permeability, and the development of permeability anisotropy 115 towards the conduit margin (Gaunt et al., 2014), thus describing a scenario where shearing of dense, 116 crystal-rich magma induced dilation. In the case of Mount St. Helens' Spine 7, the fault zone is further 117 defined by the presence of a pseudotachylyte (Kendrick et al., 2012), a feature which can decrease the 118 permeability of fault zones (Kendrick et al., 2014a). At Unzen volcano, Smith et al. (2001) qualitatively 119 described the character of the shear zone developed in the centre of the lava spine at Mount Unzen, 120 highlighting the presence of a dilational cavity associated with shearing in the core of the magmatic 121 column. However, they did not quantify any porosity-permeability relationships. The cavity (hereafter termed "central shear zone") was defined by an area in which the groundmass was torn, producing pore 122 123 spaces in the shadow of phenocrysts. The margin of the Unzen spine also hosts a spectrum of shear 124 textures (Hornby et al., 2015; Wallace et al., 2019), and significant low-frequency seismicity during the 125 eruption indicated flushing of fluids in the marginal fault zone (Lamb et al., 2015). Thus, the study of 126 evolving monitored signals and eruptive products at Unzen depicts a wide range of outgassing 127 pathways, which evolve during the course of magma ascent and lava dome eruptions.

128





129 1.2. The permeability of magmas and rocks

Several studies have explored the permeability evolution of volcanic materials, but due to the 130 131 occurrence of many influential structural and petrological processes in shallow volcanic conduits, no solutions yet encompass the complete history of magma permeability during volcanic eruptions: 132 especially its time- and strain-dependent evolution. Following nucleation and growth, bubbles interact 133 134 and coalesce beyond a certain vesicularity, termed the percolation threshold, promoting the onset of 135 fluid flow through a connected bubble network (Baker et al., 2012; Eichelberger et al., 1986; Rust and 136 Cashman, 2004; Burgisser et al., 2017). The porosity of the percolation threshold varies widely 137 (between ~30 vol. % and 78 vol. % bubbles) depending on the size and geometry distributions of the 138 bubble population (Colombier et al., 2017; Rust and Cashman, 2004; Burgisser et al., 2017). 139 Vesiculation experiments have shown that permeability remains low in isotropically vesiculated 140 (aphyric and crystal-bearing) magmas as percolation initiates at vesicularities higher than those 141 theoretically predicted (Okumura et al., 2012; Okumura et al., 2009). Yet, bubble coalescence may be 142 accentuated by transport processes such as the thinning or draining of melt along the bubble wall (Castro 143 et al., 2012), deformation (Ashwell et al., 2015; Kennedy et al., 2016; Okumura et al., 2010; Okumura 144 et al., 2006; Okumura et al., 2008; Wadsworth et al., 2017; Shields et al., 2014; Farquharson et al., 145 2016b; Kendrick et al., 2013), and rupture (Lamur et al., 2017; Lavallée et al., 2013; Heap and Kennedy, 146 2016; Okumura and Sasaki, 2014; Heap et al., 2015a; Laumonier et al., 2011), or lessened by fracture 147 infill (Kendrick et al., 2014a; Kendrick et al., 2016; Wadsworth et al., 2016b), all of which influence 148 the permeability of magma and promote permeability anisotropy (Farquharson et al., 2016c) during its 149 prolonged ascent to the Earth's surface.

150 In recent decades, laboratory measurements have helped us gain a first order constraint on the 151 permeability-porosity relationships of volcanic products (Eggertsson et al., 2018; Mueller et al., 2005; 152 Acocella, 2010; Rust and Cashman, 2011; Colombier et al., 2017; Farquharson et al., 2015; Klug and 153 Cashman, 1996). These suggest a non-linear increase of permeability with porosity; yet, depending on 154 the nature of the porous network, influenced by eruptive history, the permeability of rocks with a given 155 porosity may range by up to 4-5 orders of magnitude. Controlled laboratory experiments have given us 156 insights on probable permeability trends of magma subjected to different stress, strain, and temperature 157 conditions (Ashwell et al., 2015; Kendrick et al., 2013; Lavallée et al., 2013; Okumura et al., 2012; 158 Okumura et al., 2006; Shields et al., 2014), but a complete description of the dynamic permeability of 159 deforming magma requires in-operando determination under controlled conditions, which remain 160 scarce (Gaunt et al., 2016; Kushnir et al., 2017; Wadsworth et al., 2017; Wadsworth et al., 2021); these 161 studies have shown that surface tension and/or low-strain rate conditions under positive effective 162 pressure (i.e., confining pressure greater than pore pressure) promote compaction and reduce 163 permeability. These informative descriptions require further inputs to enable robust relationships with 164 magma rheology, influenced by the presence and configuration of bubbles. Shallow magmas contain 165 bubbles and crystals and exhibit a non-Newtonian rheology (Caricchi et al., 2007; Lavallée et al., 2007; 166 Lejeune et al., 1999; Lejeune and Richet, 1995; Kendrick et al., 2013; Coats et al., 2018) that favours 167 the development of strain localisation, in particular, by preferentially deforming pore space (Kendrick 168 et al., 2013; Okumura et al., 2010; Shields et al., 2014; Pistone et al., 2012; Mader et al., 2013) As 169 magma shears, the porous network adopts a new configuration reflecting the stress conditions and 170 magma viscosity (Rust et al., 2003; Wright and Weinberg, 2009), which influences the permeability 171 (Ashwell et al., 2015; Kendrick et al., 2013; Okumura et al., 2010; Okumura et al., 2009; Okumura et al., 2006; Okumura et al., 2008; Okumura et al., 2013). Shearing may increase or decrease the porosity 172 173 and permeability depending on the applied stress, strain and porosity of the deforming material and 174 direction of the permeability measurement due to the development of anisotropy (Ashwell et al., 2015; 175 Kendrick et al., 2013). In cases of extreme shear, magma may rupture, thereby increasing pore





176 connectivity and permeability (Laumonier et al., 2011; Lavallée et al., 2013; Okumura et al., 2013) until the fracture heals via diffusion (Okumura and Sasaki, 2014; Tuffen et al., 2003; Lamur et al., 2019; 177 178 Yoshimura and Nakamura, 2010), seals via secondary mineralisation (Heap et al., 2019; Ball et al., 179 2015), or infills with tuffisitic material (Castro et al., 2012; Kendrick et al., 2016; Kolzenburg et al., 2012; Tuffen and Dingwell, 2005), which may densify through time (Kendrick et al., 2016; Vasseur et 180 181 al., 2013; Wadsworth et al., 2014; Farguharson et al., 2017). The densification of magma under isotropic 182 stresses (due to surface tension) has been reconstructed using high-resolution x-ray computed 183 tomography from synchrotron imaging, providing us with a first complete description of magma permeability evolution as a function of porosity. This indicates that densification intrinsically relates to 184 185 the evolution of the size distribution and surface area of the connected pore space (Wadsworth et al., 186 2017; Wadsworth et al., 2021). Nonetheless, a time- and strain-dependent description of the 187 development of the porous network of shearing magma remains incomplete, and information must be 188 sourced from our understanding of permeability evolution in deforming rocks.

189 In rock physics, the evolution of the porous network in deforming rocks has been extensively 190 studied. In its simplest description, the modes of deformation differ at low and high effective pressures 191 as rocks adopt brittle or ductile behaviour, respectively. These are defined as a macroscopic behaviour 192 (not a mechanistic description), whereby 'brittle' refers to the localisation of deformation leading to 193 rupture, and 'ductile' refers to the inability for rocks to localise strain during deformation (e.g., Rutter, 194 1986); see Lavallée and Kendrick (2020) and Heap and Violay (2021) for reviews of brittle and ductile 195 deformation in volcanic materials. The key distinction between these two deformation modes is that 196 brittle failure results in dilation (i.e. the creation of porosity), whereas ductile deformation results in compaction of the porous network (Heap et al., 2015a). As a result, brittle (dilational) failure generally 197 198 enhances the permeability of rocks (Heap and Kennedy, 2016; Lamur et al., 2017; Farquharson et al., 199 2016b), whereas ductile (compactional) deformation generally causes reduction in permeability (Heap 200 et al., 2015a; Loaiza et al., 2012). Despite its crucial role in defining deformation mode in rock, the role 201 of effective pressure in dictating the ductile and brittle modes of deformation has not been 202 systematically mapped out for multiphase magma; instead, we generally consider the effects of 203 temperature and applied stress or strain rate (e.g., Lavallée et al., 2008) over that of stress distribution, 204 as the deformability of magma imparts technical challenges to classic rock mechanic tests and 205 permeability determination (Kushnir et al., 2017). We may thus anticipate some similarities between 206 rock and magma deformation modes, whereby: At high effective pressure, ductile deformation is 207 favoured via compactant viscous flow or even cataclastic flow (if strain rates is relatively high to cause 208 pervasive fracturing of bubble walls), causing porosity and permeability decrease; at low effective 209 pressure, viscous flow may promote compaction at low strain rates whereas dilation may ensue if strain 210 rate favours localised rupture (Lavallée and Kendrick, 2020). Across this transition, magma rupture 211 may be partial and end abruptly, leaving a blunt fracture tip (Hornby et al., 2019). Most, if not all, of 212 the features observed in experimentally deformed rocks and lavas should be observable in a shallow 213 magmatic system hinging on a delicate balance between ductile and brittle deformation regimes, 214 promoted by outgassing which induces temporal and spatial variations in effective pressure. In this 215 study, we examine the well-preserved, dacitic lava spine erupted in 1994-95 at Unzen volcano to constrain the permeability of dilational and compactional shear zones that developed in the shallow 216 217 volcanic conduit.

218

219 1.3. 1990-1995 eruption of Unzen volcano

220

Unzen volcano is a stratovolcano located near the city of Shimabara on the island of Kyushu,





221 Japan (Fig. 1). The volcano underwent a 5-year period of protracted dome growth which threatened the surrounding population with the occurrence of several thousand rockfalls and many pyroclastic flows, 222 such as the destructive event on 3rd June 1991 that caused 43 fatalities. Activity initiated in early 1990 223 224 with a series of phreatic explosions and brief extrusion of a spine on 19th May; this was swiftly followed 225 by continuous growth of a lava dome until early 1995 (Nakada et al., 1995). Between October 1994 and 226 January 1995, the eruption concluded with the extrusion of a spine through the dome surface (Fig. 1c). 227 At the dome surface, gas emissions focused along the spine marginal faults (Ohba et al., 2008). The 228 dome products have a dacitic composition and contain euhedral phenocrysts of plagioclase and amphibole in a groundmass containing microlites of plagioclase, amphibole, pyroxene and iron oxides 229 230 (Nakada et al., 1995; Wallace et al., 2019). Petrological constraints suggest that degassing initiated at a 231 pressure of approximately 70-100 MPa; i.e., in the upper ~3-4 km depth (Nakada et al., 1995).

232 Dome growth occurred in stages, forming thirteen discrete lobes until mid-July 1994. Growth 233 was observed to be typically exogenous when effusion rates were high, and endogenous at effusion 234 rates lower than 2.0 x 10^5 m³d⁻¹ (Nakada et al., 1999). In five years, the eruption generated 2.1 x 10^8 m³ of lava at an average ascent rate estimated at 13-20 md⁻¹ (Nakada et al., 1995); the final spine extruded 235 from late-1994 to early-1995 at a rate of approximately 0.8 md⁻¹ (Yamashina et al., 1999). The rheology 236 237 of the erupted dome lavas has been sourced of debate (Goto et al., 2020; Sato et al., 2021), as it is challenging to precisely reconstruct the physico-chemical, petrological and structural parameters which 238 239 control rheology as a function of depth during eruption. For the late-stage spine, Nakada and Motomura 240 (1999) proposed that it formed due to a lower effusion rate, which resulted in extensive magma degassing and crystallisation, and thus high viscosity, which promoted rupture and exogenic growth at 241 242 relatively low strain rates (e.g., Hale and Wadge, 2008; Goto, 1999). Extrusion occurred through pulsatory magma ascent, accompanied by ~40 h inflation/deflation cycles (Yamashina et al., 1999) and 243 244 a rhythmic pattern of summit earthquakes, interpreted to result from magma rupture in the top 0.5 245 kilometre of the conduit (Lamb et al., 2015; Umakoshi et al., 2008); waveform correlation of the seismic 246 record revealed rhythmic seismicity grouped into two primary clusters (Lamb et al., 2015). Hornby et 247 al. (2015) statistically analysed the slip duration of seismic events in the clusters, defining a mode and 248 mean of 0.1 s. As magma ascent occurred through an inclined conduit (Umakoshi et al., 2008), the spine extruded at an inclined angle of \sim 45° towards the ESE (Fig. 2a) and increasingly leaned against the 249 250 lower fault zone as extrusion rate waned, causing the shallowing of seismogenic magma rupture in this 251 area (Lamb et al., 2015). In contrast, the upper fault zones may have opened up as the spine settled, thus 252 triggering rupture at increasing depth and promoting preferential pathways for fluid flow (Lamb et al., 253 2015). By the end of the eruption, the spine achieved approximate dimensions of 150 m length, 30 m 254 width and 60 m in height (Nakada and Motomura, 1999; Nakada et al., 1999); it is complemented by 255 multiple fragments of spines, extruded earlier in the eruptive phase, which we examine in this study. 256 Unfortunately, the lower and upper fault zones are not observable in the spine exposures, but the 257 northern lateral conduit margin contains well-defined shear zones (Smith, 2002; Smith et al., 2001), 258 which are revisited here and augmented by structural and microtextural descriptions as well as porosity 259 and permeability constraints. Our study of the spine sheds new light on the permeability evolution of its shear zones, and thus the nature of outgassing during the waning phase of the 1990-1995 eruption. 260

261

262 2. Materials and Methods

263 2.1 Localities and sample collection





264 The 1994-95 lava spine was investigated during two field campaigns, in November 2013 and 265 May 2016. Close structural examination at different scales forms the basis of this study along with 266 porosity and permeability measurements, using field and laboratory equipment. Owing to the inclination of the spine (extruded towards the east), large blocks ranging from 5 to 20 m-wide are dislocated from 267 268 the front of the in situ western main spine structure (Fig. 2a, b). Here, we investigated two blocks that 269 reveal a central shear zone (CSZ) and marginal shear zone (MSZ) that developed in the spine. These 270 detached, yet fully intact, spine blocks were selected owing to their contrasting shear textures that would 271 have represented different positions within the volcanic conduit during magma ascent and extrusion (i.e., central vs. marginal), thus allowing assessment of syn-eruptive outgassing pathways. The marginal 272 273 shear zone (MSZ) block, located ~60 m east of the main spine (Latitude: 32.76131° Longitude: 274 130.29983°), was carefully sampled to quantify the spatial distribution of permeability across the spine 275 margin (samples A-H; Fig. 2c). The central shear zone (CSZ) block, located centrally between the main spine and MSZ (Latitude: 32.761271° Longitude: 130.299472), features the dilatational cavity 276 277 (described in Smith et al., 2001) and was also studied in situ, using non-destructive methods to preserve 278 the integrity of this exemplary feature. The main spine and CSZ are protected by UNESCO heritage 279 site regulations (Figs. 1c, 2a), thus only permitting *in situ* sample collection from the MSZ.

280

281 2.2 Sample preparation

282 Samples collected from the marginal shear zone were cut and cored parallel to the shear 283 direction and perpendicular to the shear plane in order to constrain the anisotropy developed in shear 284 zones. A total of eight thin sections (fluorescent dyed) were prepared for microtextural analysis (labelled 285 A-H). For the largest samples (A, B, C, E, H; see Fig. 2c-d) a set of 2-3 cylindrical cores (two parallel 286 and one perpendicular to shear plane) were prepared with a diameter of 26 mm and a length of 30 or 13 287 mm, depending on the size of the sample. Within the highly sheared sample B (Fig. 2c-d), which is 288 directly adjacent to the fault and gouge zone, multiple sets of cores of 20 mm diameter were prepared, 289 closely spaced, to obtain porosity/permeability determinations at a higher resolution across this defining 290 part of the shear zone.

291 2.3 Microstructural analysis in 2D and 3D

2D analysis of the microstructures exhibited across the shear zones was carried out using a
 Leica DM2500P optical microscope in plane polarised and ultraviolet (UV) light, as well as a Philips
 XL30 scanning electron microscope (SEM) in backscattered electron (BSE) mode, set at 20 kV and 10
 mm working distance. For this purpose, representative features were imaged for each sample across the
 shear zone (Fig. 3).

297 To further evaluate the architecture of the porous network in three dimensions (3D), four 298 samples collected across the shear zone were scanned using a phoenix nanoton® m x-ray computed 299 tomography scanner to produce high-resolution reconstructions with a voxel size of 11.111 µm. For each sample we acquired 1440 radiographs, scanning 360°, under the following conditions: exposure 300 301 time of 1000 ms; voltage of 80 kV; current of 120 µA; 0.2 mm aluminium filter. The radiographs were 302 then reconstructed using the inverse Radon transformation (Radon, 1986), resulting in a 3D image of 303 the sample. These files were processed in FEI Avizo and ImageJ/Fiji software to illuminate the 304 permeable, porous network.

305

306 2.4. Porosity measurement in the laboratory





307 Each core was dried in an oven at 50 °C overnight, then kept in a desiccator (for thermal equilibration to ambient conditions) before being weighed and loaded in a pycnometer. The fraction of 308 connected pores (which controls permeability; Colombier et al., 2017) was determined using a 309 Micromeritics AccuPyc II 1340 helium pycnometer. The porosity determination first necessitated 310 311 measurement of the geometric volume of the sample (Vsample). Then, once inserted in the specimen 312 chamber of the pycnometer, helium gas was injected in the chamber to estimate the volume taken up by the solid fraction of the sample, thus providing the skeletal volume $(V_{skeletal})$ of the rock. The 313 314 fraction of connected pores ($\phi_{connected}$) in a sample was then calculated via:

315
$$\phi_{connected} = \frac{(v_{sample} - v_{skeletal})}{v_{sample}}$$
 (1).

316

317 2.5. Permeability determination in the laboratory

318 The prepared cores were jacketed with a VitonTM tube and inserted in a hydrostatic cell from Sanchez technologies to measure permeability and pore volume as a function of pressure. The jacketed 319 320 samples were externally loaded using a Maximator \mathbb{R} oil pump to various confining pressures (P_c) and internally loaded using water to an average pore pressure (P_p) of 1.25 MPa, in order to obtain a range 321 322 of effective pressures $(P_{eff} = P_c - P_p)$ from 5 to 100 MPa. Each time the sample was loaded to new confining pressure increment, the volume of water expelled from the pores in a given sample (due to 323 324 compaction) was monitored to constrain pore volume change due to crack closure as a function of 325 pressure (Lamur et al., 2017). Steady-state flow permeability (k) was measured by applying low pore pressure gradients (ΔP) of 0.5 and 1.5 MPa to ensure laminar flow with no slip conditions (Heap et al., 326 327 2017a) to satisfy Darcy's Law:

$$k = \frac{Q\eta L}{A(\Delta P)}$$
(2),

where Q is the flow rate monitored through the sample $(m^3 s^{-1})$, η is the viscosity of the water in pores (Pas), L is the length of the sample (m), and A is the cross-sectional area of the sample (m^2) .

331

332 2.6 In-situ permeability measurements in the field

333 To measure the permeability of rocks in the central shear zone (CSZ; Fig. 1c) that could not be 334 sampled for laboratory testing due to preservation restrictions, we used a non-destructive, portable, air 335 permeameter (TinyPerm II) from New England Research, which estimates permeability by monitoring 336 pressure recovery rate from a vacuum, based on the concept of transient pulse permeability (Brace et al., 1968). The apparatus is hand-held and needs to be employed carefully to maintain a consistent seal 337 338 between the nozzle of the permeameter and rock surface throughout the measurements (lasting up to a few tens of minutes). It may be used to determine the permeability of rocks between approximately 10-339 ¹² to 10⁻¹⁶ m² (Farquharson et al., 2015; Kendrick et al., 2016; Lamur et al., 2017). In this study, three 340 341 transects were measured across the central shear zone and all measurements were performed twice to 342 ensure precision of the method (as determined in Lamur et al., 2017).

343

344 3. Observations and results





345 The 1994-95 spine structure at Mount Unzen is exposed in several large segmented blocks (Fig. 346 1c-d; Fig. 2a-b). A thorough structural description of the main spine structure and subsidiary block (e.g., 347 CSZ) can be found in Smith et al. (2001); here we highlight the main features. The lava spine is split into a few very large, primary blocks, ~20-30 m wide and high (Fig. 1c-d, 2a-b), broken roughly 348 perpendicular to extrusion direction: westward and inclined (see Fig. 2b). The CSZ block seen in Figure 349 350 1c shows a >8-m wide variably deformed core (I) lying adjacent to a 2-m wide intensely sheared zone 351 (II), bordered to the north by a dextral fault and coupled to a large indurated breccia (III), uplifted from 352 the surrounding dome emplaced. The lower and southern edges were not exposed. The upper edge of the spine was not accessible, but we noted large, incoherent brecciated blocks. The rear of this outcrop 353 as well as the main *in situ* spine structure exhibit irregular, metre-scale polygonal joints, although these 354 355 are not developed in the face of the outcrop studied here (Fig. 1c). Additional fragments of the spine 356 occur in a few subsidiary blocks (e.g. Fig. 1d), located a few tens of meters to the east of the main spine 357 (Fig. 2a). These blocks, which were emplaced prior to the main spine, expose several sections through 358 the spine, and reveal the evolving architecture of the shear zone in the shallow magmatic conduit. One 359 such block, shown in Figure 1d, exhibits a ~1-m wide shear zone, bordered to the left by a set of oblique 360 tensile fractures, reaching 2-5 m in length and spaced at \sim 3 to \sim 10 cm intervals, and to the right by an 361 indurated breccia. This prominent block was not sampled or further studied to preserve its integrity.

362

363 3.1 The marginal shear zone

364 3.1.1. Structural and microtextural observations

365 Our primary field location for this study was a 4.7-m wide block of the spine, exposing the 366 northern marginal shear zone consisting of gouge, sheared lava and the spine core (Fig. 2c-d). The outcrop displayed mild surface weathering, in the form of a thin (micron-size) veneer of unknown 367 368 precipitate on the rock surface (which was inclined at an angle of ca. 40° towards the West). This thin 369 veneer did not visually obstruct any primary magmatic textures and structures, and the shear texture 370 was clearly visible, yet we it would prevent accurate field permeability constraints. Four distinct degrees 371 of shear were visually defined through textural examination and changes in surface roughness across 372 this section of the conduit (Fig. 2c-d): a fault gouge zone (sample A) bordering a high-shear zone 373 (samples B, C, D), a moderate-shear zone (samples E, F) and low-shear spine core (samples G, H) in 374 decreasing order of surface roughness and visually observable fracture density variations; quantitation 375 of fracture density was not attempted as we deemed the thin veneer may have prevented meaningful 376 accuracy. This shear-based division is consistent with a complementary investigation of the 377 mineralogical characteristics of this shear zone (Wallace et al., 2019). The contacts between shear zones trend approximately E-W in the outcrop (Fig. 2c,d), and so roughly parallel to the spine emplacement 378 379 direction to the ENE, despite the detachment of this spine block from the main intact spine body to the 380 west. Eight samples were systematically collected across this shear zone for further analysis (labelled 381 A-G in Fig. 2c,d): eight for 2D microstructural analysis (PPL, UV light and BSE imagery; Fig. 3), four for tomographic imaging (Fig. 4) and five for porosity and permeability determination (Fig. 5-6). [Note 382 383 that multiple cores were obtained from the five blocks sampled for laboratory measurements.]

The spine core, termed low shear herein (~1.5 m wide; Fig. 2c, d), exhibited a smooth surface and the phenocrysts showed no preferred orientation at the macroscopic scale. In samples G and H collected from the low shear zone (Fig. 3), phenocrysts of plagioclase, amphibole, biotite (plus minor quartz) are typically euhedral, largely intact and up to ~5mm in length (Fig. 3); groundmass microlites also show no preferred orientation in BSE images. The porous structure is characterised by a diktytaxitic texture, composed of some large, irregular, vesicles with 'ragged' edges, appearing intrinsically related





to the presence of surrounding phenocrysts (single white arrows on UV light images in Fig. 3). Small
fractures are often seen to originate from these large vesicles, penetrating pervasively through both
phenocrysts and the groundmass (double white arrows in Fig. 3). The groundmass contains abundant
small vesicles, showing a high degree of connectivity as revealed by tomography (Fig. 4g-h).

394 The moderate-shear zone is approximately 2 m wide (Fig. 2c, d). In this zone, we observed an 395 increased fracturing of phenocrysts and changes in the distribution of porosity. Scrutinising the sample 396 E under microscopy, we observe that the phenocrysts, which rarely exceed 2 mm in size in this zone, 397 are commonly micro-fractured (Fig. 3). The vesicles are occasionally large and connected (Fig 3, 4e-f), 398 and while the vesicular texture remains diktytaxitic (as in the low shear spine core), the vesicles in 399 sample E appears increasingly aligned and localised around phenocrysts as the magnitude of shear 400 increases towards the fault; similarly, the microlites show increasing degrees of alignment (revealed by 401 undulose extinction angles; see Wallace et al., 2019). Thin bands (<200 µm width) of reduced porosity 402 are observed to localise in the groundmass (see facing double arrows in UV light images in Fig. 3), 403 which are notably absent in the low shear zone; these are (sub-)parallel to the shear plane. The 404 tomographic reconstructions show irregular vesicles, which are surrounded by fractures and invaded by 405 rock fragments (Fig. 4e). These vesicles enhance the connectivity of the porous network (Fig. 4f).

406 The high shear zone is approximately 1 m wide (Fig. 2c, d) and marks the beginning of micro-407 and meso-scopic shear bands, at a scale of the order of a few millimetres, near-parallel with the direction of shear; these increase in abundance and scale nearer the fault, especially within the final 0.1-0.2 m 408 409 (see features denoted in Fig. 2c-d as well as enlarged in the inset). The bands, which form a pervasive 410 foliation (S), consist of elongate, white porphyritic plagioclase lenses, fractured and crenulated. The C-411 S fabrics are parallel in this area. These porphyritic bands are flanked by reddish-brown groundmass as 412 well as thin, elongate biotite phenocrysts (see sample B "fresh surface" in Fig. 3). The plagioclase and 413 biotite commonly exhibit a mineral fish texture. Under the microscope, we observe that the biotite show 414 undulose extinction from crystal plastic deformation (see Wallace et al., 2019, for a detailled crystal 415 plasticity study). Intense banding (observed as faint lineations of reduced porosity under UV light in 416 the moderate shear zone; Fig. 3) is observed adjacent to, and running parallel with, the fault-gouge 417 contact. The bands are up to 1 mm wide and display variations in porosity under UV light (Fig. 3), as 418 also revealed by tomography (Fig. 4c-d). The dense bands are traversed by hairline fractures a few 419 hundred microns in length and contain a few isolated millimetre-size vesicles, generally adjacent to 420 large phenocryst fragments (samples B and C in Fig. 3). More porous bands display disordered and 421 fragmental textures (sample B), with abundant, irregular large pores and cracks, and pulverised 422 phenocrysts (PPL and UV light in Fig. 3); macroscopically, the most porous bands often appear like 423 ragged tensile fractures. The transition between dense and porous bands is abrupt, occurring over a few 424 tens of microns (BSE images of samples B and C in Fig. 3). Microlites and microphenocrysts are aligned 425 with the banding, and thus with shear and extrusion direction (Fig. 3). The high shear region of the 426 spine is further crosscut by multiple sub-parallel curvilinear extensional bands (i.e., weakly defined 427 fractures), up to ~ 1 m in length, and trending $\sim 57^{\circ}$ from the primary C-S fabrics in a Riedel-like fashion 428 (Fig. 2c, d); some of these bands extend into the moderate shear zone but only faintly. These bands, 429 spaced by 3-6 cm (~4.5 cm in average), show opening of ca. 1-2 mm in places. [Note that the blue traces 430 in Figure 2 denote the general attitude, not the spacing, of the bands]. The Riedel fractures appear to be associated with a set of faint, conjugate fractures (R'), although their observation is not ubiquitous 431 432 across the high-shear zone.

433 The fault zone hosts up to ca. 0.2-m thick gouge material (Fig. 2c,d). The contact between the 434 gouge and the high shear zone is generally sharp, and often planar, although we observed small 435 embayments, especially along C-S fabrics in the neighbouring high shear zone (Fig. 2d). [Note that the





436 extent of the gouge is not exposed equally across the outcrop as material was likely lost during 437 separation of this block from the main spine upon eruption; so the surface does not reflect the contact 438 geometry. This material loss also led to obliteration of vestiges of a pseudotachylyte, suggested by local partial melting textures presented by Wallace et al. (2019)]. The gouge is typified by well-consolidated, 439 440 fine-grained cataclasite with some larger rounded clasts up to ~15 mm in diameter (sample A; Fig. 2c 441 inset). The gouge is matrix supported and displays a strong foliation parallel to spine extrusion direction. 442 Conjugate fractures form a dominant feature contributing to the porosity of the gouge. Microscopically, 443 the rock is pervasively fragmented (sample A in Fig. 3); the few phenocrysts that remain relatively 444 intact often display signs of deformation. The fragments in the gouge are generally densely compacted 445 and the porosity is uniformly distributed, with little banding or preferred orientation of fragments at the microscopic scale, although connected pores occasionally exhibit a degree of alignment at small scale 446 447 (Fig. 3) and at large scale as observed via x-ray tomography (Fig. 4a-b).

448

449 3.1.2 Connected porosity across the marginal shear zone

450 The porosity of the rocks, determined via pycnometry, indicates variations between 8 % and 27 451 % across the shear zone and in the fault gouge; Figure 5a displays the average of multiple measurements 452 from the different cores prepared from each sample. The measurements indicate that the high shear zone 453 generally holds slightly lower porosities than surrounding areas. Within the high-shear zone (sample B) 454 we measured important variations in porosity ranging between 8 % and 15 % due to flow bands (e.g., 455 in sample B); yet, the coarseness of sample measured prevent from accurately quantifying the highly 456 variable degrees of porosity visually observable in hand specimen.

457 When loading the samples (cored parallel with to spine extrusion direction) in the hydrostatic 458 pressure vessel, we observed a nonlinear decrease in porosity of up to 4 % by increasing the effective 459 pressure to 100 MPa (Fig. 5b). The data shows a similar dependence of porosity on effective pressure 460 for the coherent samples from the low, moderate and (densest part of) high shear areas, with a slightly 461 larger reduction in porosity with effective pressure in the initially most porous, high shear bands and 462 granular gouge sample (Fig. 5b).

463

464 *3.1.3 Permeability across the marginal shear zone*

The permeability of the rocks collected across the spine segment reveals a ~1-m wide region of low permeability in the high shear zone, compared with the moderate shear zone, the low shear spine core and fault gouge (Figs. 5, 6). There appear to be abrupt variations in permeability (decrease and increase) in sheared rocks directly adjacent to the fault gouge, due to the alternation between dense and porous shear bands.

470 The data show considerable differences in the permeability parallel and perpendicular to the 471 plane of shear (Fig. 3c,d) across the shear zone (Fig 6a,b). In the high shear zone permeability was 472 found to be higher in the plane of shear (*i.e.*, parallel with extrusion direction) than perpendicular to it, 473 whereas in the moderate and low shear zones, as well as in the gouge, permeability was essentially 474 isotropic. Anisotropy is cast here as a ratio between the permeability parallel and perpendicular to the 475 shear plane (Fig. 6c). The anisotropy is most pronounced in the high shear zones, where, in one instance, 476 the permeability ratio increases dramatically from three to over seven times larger parallel than 477 perpendicular to the shear plane with increasing confining pressure in a hydrostatic pressure vessel (Fig. 6c). In other samples, the anisotropy increase with pressure is less or even negligible, indicating the 478





479 heterogenous nature of the high shear zone. This sensitivity to confinement is due to the presence of the 480 distinct dense and porous bands in the sheared lava (Fig. 5b, 6); in the cores parallel to the shear plane, 481 fluid can flow through porous bands from top to bottom of the sample, whereas perpendicular to shear, fluids must pass through both and dense and porous bands to traverse the sample. Fluid flow in the 482 483 denser areas will be dominated by channelling through narrow fractures (sub-horizontal in BSE images 484 in samples B and C in Fig. 3), which are more susceptible to closure by increasing effective pressure 485 than equant pores (e.g., Kendrick et al., 2021). Although this process occurs during confinement in both 486 orientations, it only impacts permeability perpendicular to shear direction, and so contributes to enhanced anisotropy of permeability in banded shear fabrics under confinement (Kendrick et al., 2021). 487

- 488 3.2 Central shear zone
- 489 3.2.1 Structural observations

490 The second feature of interest is the cavity exposed in the central shear zone block (Fig. 1c and 491 2a). This section of the spine has been described in detail by Smith et al. (2001); here, we review key 492 aspects observed in the field as no samples were collected to conserve the exposure of this world-class 493 feature. We only examined the rocks forming this structure and performed non-destructive, *in-situ* 494 testing.

495 The central shear zone (CSZ) is located near the centre of the spine core (Fig. 1c). Its primary 496 feature is the presence of a porous cavity, which curves and pinches out (upward) from the end of a 497 dominant, 9-cm wide fracture, extending approximately 3 m in length (determined from the visible 498 extent of the exposure). Unlike the aforementioned marginal shear zone, which displays an increased 499 degree of shear towards the spine margin, the central shear zone exhibits an increase in shear towards 500 the centre of the spine. From left to right (i.e., northward) on Figure 7, we note an increase in aligned, bent and broken phenocrysts as well as aligned shear bands (ostensibly parallel with the dominant 501 502 fracture), fractures and surface roughness, which terminates upon intersecting the end cavity; beyond 503 which point, the rocks show no clear evidence of shear, including shear bands, elongate pores or aligned 504 crystals. This is evident in the field photograph (Fig. 7) as steeply inclined porous bands which ends 505 against the southern (i.e., right) side of the cavity; on the southern side the sheared lava exhibit a higher 506 porosity than the surrounding undeformed rocks (although this could not be quantified in the field). 507 Approximately 1 m above the pinched-out tip of the main cavity, we observe the presence of a 508 secondary porous cavity (Fig. 1c inset), approximately 60 cm long, and elongated parallel to the fracture 509 that connects to the main cavity.

510

511 3.2.2 Permeability across the central shear zone

The permeability of the rocks in the central shear zone was measured along three transects in two field campaigns (in November 2013 and May 2016) to negate potential influence from variable degrees of water saturation of the rocks at different times of year. Our field measurements are consistent with one another. The permeability varies very little in the undeformed areas of the outcrop (i.e., on the right-hand side of the fracture in Fig. 7) for all transects, with an abrupt increase in permeability up to three orders of magnitude in the 9cm wide central cavity, and elevated permeability in the ~40 cm wide proximal sheared area to the left of the fracture.

519

520 4. Interpretation





521 The contrasting permeability, porosity and (micro)structural changes observed across the 522 marginal and central shear zones reveal the impact of shear and distinct modes of magma deformation 523 during shallow conduit ascent. Here we interpret each of these key features for the development of 524 volcanism at lava domes.

525 Marginal shear zone

526 The marginal shear zone is characterised by a 3-m wide zone in which strain caused changes in 527 the porous structure, via crushing of the pore walls as well as distortion and failure of the crystalline 528 phase; these promoted an increased reduction in pore volume and permeability towards the fault, 529 especially in the high shear zone. Smith et al. (2001) invoked the effects of gravitational forces during 530 post-emplacement flow of the lobes as a mechanism for the development of 'ragged' pores and 531 porous/dense flow banding in dome lavas at Unzen volcano. Yet, such diktytaxitic structure have been 532 observed in small surficial dome blocks at Santiaguito volcano (Guatemala), which have not suffered from gravitational effects associated with flow along the flanks (Rhodes et al., 2018). This diktytaxitic 533 534 texture has been observed in the experimental products of lavas compacted under uniaxial (Ashwell et al., 2015) and triaxial (Kushnir et al., 2017) conditions. Similarly, they can be reproduced (to a high 535 536 degree of similarity) through shear-enhanced compaction of porous rocks under high effective pressures (Heap et al., 2015a; Heap et al., 2015b). The commonality between these experiments is that they were 537 538 carried out in the ductile field, through which material may sustain substantial compaction without the propensity for developing localised strain (Rutter, 1986) - a regime that results in a permeability 539 540 reduction through shear (Ashwell et al., 2015; Kushnir et al., 2017; Heap et al., 2015a; Heap et al., 2015b). In this regime, magma deformation may result in crystal plastic distortion and failure (Kendrick 541 542 et al., 2016), as witnessed at Unzen (Wallace et al., 2019). Thus, we interpret the bulk of the marginal 543 shear zone as the result of ductile deformation, which resulted in distributed, pervasive shear over a 544 width of 3 m. Within this part of the conduit, the high shear zone displayed the highest degree of shear-545 enhanced compaction.

However, ductility alone is insufficient to describe the marginal shear zone. For instance, the high-shear 546 547 area exhibits a foliation (S plane) and fractures (C plane) parallel to the shear plane, which is then 548 crosscut (parallel but undulating) by a marginal fault hosting gouge formed by comminution and 549 cataclasis, containing conjugate fractures. The composite C-S fabric in the high shear zone is 550 increasingly penetrative towards the fault core (at the gouge contact), and its parallel C and S planes 551 indicates that the shear zone accommodated significant strain. This is supported by observation that 552 curvilinear Riedel fractures have developed and overprinted the C-S fabric at an angle of 57° (cf. 553 Ramsay, 1980). Such an angle is consistent with a lava body undergoing rupture following sustained 554 ductile deformation (e.g., Lavallée et al., 2013); it is also consistent with the progressive thickening of 555 a shear zone formed via simple shear with a small component (<10 %) of pure shear (assuming pure 556 and simple shear are planar; Fossen and Cavalcante, 2017); this minor pure shear component is further 557 supported by the presence of weakly defined conjugate fractures crosscutting the Riedel fractures. Both 558 the gouge and the fractures through the high shear zone were constrained to have locally higher 559 permeability and porosity than the bulk of the shear zones: features characteristic of dilational 560 deformation resulting from macroscopically brittle failure (Heap et al., 2015a; Heap et al., 2015b; 561 Laumonier et al., 2011). Riedel fractures generated in experimentally deformed magma have been 562 described as important pathways to redistribute fluids across shear zones (Laumonier et al., 2011), and 563 we anticipate the impact would be similar at Unzen; the Riedel fracturs in the marginal shear zone only 564 reached ~1m in length, but the marginal shear zone in other blocks (Fig. 1d) contain oblique Riedel fractures that reach 2-5 m in length (Fig. 1d) which would have formed efficient fluid flow pathways. 565 566 Thus, we interpret the marginal shear zones to reflect the evolution of magma shearing across the ductile 567 to brittle transition during shallowing of the ascending spine, which impacted fluid flow during eruption.





568 Central shear zone

The central shear zone detailed in this study has a very different character. Macroscopic 569 570 observations of numerous cracks suggest that it is dominantly dilational, as supported by the drastic 571 increase in permeability towards the fault and cavity. Despite having opened by ~9 cm, the main fracture tip is blunted as it terminates in a curvilinear cavity, and seemingly disappears before reappearing as a 572 573 secondary cavity 1 m above (Fig. 1c inset). This is akin to areas of reduced density that develop ahead 574 of a crack tips during material failure in the lab (e.g., Célarié et al., 2003) and indicates immature shear 575 that was insufficient to enable the continuous propagation of a fault across the whole spine. This, in 576 conjuncture with the observation that shear becomes more pronounced towards the centre of the spine, 577 suggests that the areas undergoing shear may have locally shifted towards the conduit core; yet, 578 displacement was not extensive. The reason for this shift is difficult to assert, but we posit that the 579 shallow calving of blocks from the spine front, progressive inward cooling and/ or the higher porosity 580 of the magmas in the conduit core (compared to a denser, compacted and strained conduit margin) may have shifted the locus of deformation towards the conduit core at the end of the eruption. 581

582 The shear zones studied here indicate that the dominant deformation regime of magma may 583 evolve spatially and temporally during ascent in volcanic conduits, which would modify the magma's 584 permeability and its ability to localise and channel outgassing during the effusion of lava domes.

585

586 5. Discussion

587 Permeability in volcanic environments

588 The power of volcanic eruption models relies on an understanding of the coupling between 589 magma and volatiles in volcanic conduits (Sparks, 1997), yet a description of dynamic permeability of 590 deforming magma eludes us. The studies of eruptive products have provided first order constraints on 591 the relationship between permeability and porosity (Fig. 8; Klug and Cashman, 1996; Mueller et al., 592 2005; Farquharson et al., 2015) for various types of volcanic rocks (e.g. explosive clasts vs effusive 593 lavas), including the presence of heterogeneous structures (Farquharson et al., 2016c; Kolzenburg et al., 594 2012; Lamur et al., 2017; Kendrick et al., 2021), and these constraints have been invoked in diverse 595 models to assess how magma permeability may evolve leading to eruption (Burgisser et al., 2019; 596 Edmonds et al., 2003). However, the deformability of magma imposes constant changes to the porous 597 permeable network and to date, only a few studies have measured or assessed the transience of 598 permeability and porosity during magma deformation (Okumura et al., 2010, 2012; Kendrick et al., 599 2013; Ashwell et al., 2015; Kennedy et al., 2016), especially in operando (Kushnir et al., 2017; 600 Wadsworth et al., 2017; Wadsworth et al., 2021). Considering the range of pressure conditions (e.g., 601 pore pressure gradient, local deviatoric stress) and magma properties, none of these studies has yet 602 succeeded in fully reconstructing the evolution of porosity and permeability of magma shearing during 603 ascent in volcanic conduits.

604 The rocks sampled across the shear zone and in the fault gouge at Mount Unzen vary in porosity 605 between 8 % and 27 %; this range is slightly narrower than the porosity range (4-48 %) covered by 606 blocks shed by pyroclastic density currents originating from the domes during the 5-year eruption (see Fig. 8; Kueppers et al., 2005; Coats et al., 2018; Kendrick et al., 2021; Scheu et al., 2007; Mueller et 607 608 al., 2005). The narrower range exhibited by the spine shear zones may reflect the occurrence of fewer 609 porosity-modifying mechanisms (e.g. post-fragmentation vesiculation) in the highly viscous spine lava 610 compared to those which occurred throughout the entire course of the eruption, which are represented by the blocks at the foot of the volcano. We see the largest contrast when we compare the permeability 611





range of the lavas which erupted through the spine at the end of the eruption ($\sim 10^{-15}$ to $\sim 10^{-14}$ m², at the 612 lowest effective pressure) with that obtained from rocks recovered by drilling through the eruptive 613 conduit at a depth of ~1.5 km (~10⁻¹⁷ to ~10⁻¹⁹ m²) in the framework of the Unzen Scientific Drilling 614 615 Project, drill hole 4 (USDP-4) (Watanabe et al., 2008). The latter rocks, originating from magma stalling at depth, reflect greater time under compactant conditions and porosity infill and reduction from 616 617 secondary mineral precipitation (Yilmaz et al., 2021). The large difference in permeability between the 618 two datasets alludes to the highly variable spatial and temporal variation of magma permeability within 619 even a single volcanic system.

620 Previous investigations of permeability in shallow volcanic conduits have highlighted the 621 existence of dilational shear zones, whereby the conduit margin is bound by a permeable 'damage halo'; 622 this has been proposed through both field (Saubin et al., 2019; Pallister et al., 2013a; Gaunt et al., 2014; 623 Wallace et al., 2019) and laboratory (Lavallée et al., 2013; Laumonier et al., 2011) studies. These 624 constraints indicate a high-permeability zone, with a strong component of anisotropy, with fluid flow 625 preferentially developed in the direction of extrusion due to shear fabrics (Wright et al., 2006; Gaunt et 626 al., 2014; Wallace et al., 2019). Connectivity is enhanced by fractures, which would contribute to the 627 development of anisotropy and preferential channelling of fluids along the conduit margin, promoting 628 concentric or ring-like gas emissions, as for instance exemplified at Santiaguito, Guatemala (Lavallée 629 et al., 2013). Here, at the conduit centre at Unzen we observed a localised dilational shear zone up to 630 three orders of magnitude more permeable than the surrounding magma. This zone spans a relatively 631 narrow section of the conduit and appears to be a late, immature feature that is possibly related to shear during the final stages of ascent and/ or structural readjustment during failure and calving of portions 632 of the spine to the ENE. Instead, the primary (and volumetrically most significant) marginal shear zone 633 634 studied at Unzen is mostly compactional and exhibits a lower permeability than the surrounding magma, 635 particularly in the plane perpendicular to shear direction. It appears to have formed at depth, before 636 being overprinted by shallower faulting. Seismic analysis indicated that seismogenic faulting was episodic and shallow, likely originating in the upper 500 m of the conduit (Umakoshi et al., 2008; Lamb 637 638 et al., 2015); the pulsatory magma shearing above this depth would have resulted in switches between 639 compactional and dilatant shear, causing locally higher permeability fractures through the sheared 640 magma, and a permeable marginal fault gouge by cataclasis (Fig. 9). Such intermittent seismic stressing 641 may also serve to weaken surrounding country rocks and modify permeable pathways (Schaefer et al., 642 2020).

643

644 Ductile-brittle transition in ascending magma

645 The presence and overprinting of compactional and dilational shearing modes in close 646 proximity in a given magmatic extrusion demands appraisal. The ductile-brittle transition of materials 647 has long been studied and is generally better understood for rocks than lavas as more low-temperature 648 tests have been carried out (Paterson and Wong, 2005; Rutter, 1986; Heap et al., 2015a). Reconstruction 649 of yield caps (or curves), based on the shear stress required for rupture or flow of materials at different 650 effective mean stress, have shown that porous rocks undergo a transition from macroscopically brittle to ductile deformation modes with increasing effective pressure (Fig. 9b); this transition sets in at lower 651 652 effective pressure (i.e., either at shallower depths or with higher pore pressures) if the material is more 653 porous (Heap et al., 2015a; Coats et al., 2018). However, magma is viscoelastic, thus depending on the 654 timescale of observations magma may behave as a solid; in essence, as a rock. Magmas abide to the 655 glass transition so that at long observation timescales or under slow deformation, they flow; but at short 656 timescales or if strain rate is high, they may rupture (Dingwell, 1996). The strain rate to meet this





657 transition decreases if melt viscosity increases due to cooling, crystallisation, degassing, and/ or vesiculation (Wadsworth et al., 2018; Dingwell and Webb, 1989, 1990; Cordonnier et al., 2012; 658 Cordonnier et al., 2009; Coats et al., 2018; Lavallée et al., 2013; Lavallée et al., 2008). The glass 659 transition of silicate melts, which controls the deformation mechanisms of magmas (viscous or brittle), 660 thus impacts their deformation modes, brittle or ductile (be it viscous flow or cataclastic flow); 661 662 applicability of the concept of yield caps to volcanic rocks and magma, as shown in Figure 9b, have 663 been reviewed by Lavallée and Kendrick (2020). In a scenario where magma ascends, deforms and 664 outgasses during an eruption, such as during spine extrusion at Unzen, magma may undergo a transition from a macroscopically ductile to brittle deformation mode due to a reduction in effective pressure 665 (from ascent or due to pore pressure increase; Heap et al., 2017b), densification (Heap et al., 2015a; 666 Coats et al., 2018), viscosity increase (cf. Dingwell and Webb 1990) or if the strain rate locally increases 667 (Coats et al., 2018; Lavallée et al., 2013; Lavallée et al., 2008). 668

669 Nakada and Motomura (1999) proposed that faulting of this spine formed due to a lower 670 effusion rate that resulted in more complete degassing and crystallisation that increased the magma 671 viscosity. We advance that fluctuations in pore pressure (Farquharson et al., 2016a) and local strain 672 rates (Coats et al., 2018; Lavallée et al., 2013; Wadsworth et al., 2019) may be especially important in 673 triggering embrittlement of otherwise ductile magma. In the ductile regime, strain is accommodated 674 over prolonged duration without necessarily leading to any substantial stress drop (Coats et al., 2018). 675 Thus, under such conditions, we do not expect to detect any, or much, seismicity that would characterise 676 magma rupture near the conduit margin (e.g., Neuberg et al., 2006; Thomas and Neuberg, 2012; Kendrick et al., 2014b). As a result, we anticipate that magma shearing below the point of rupture (ca. 677 0.5 km at Unzen; Umakoshi et al., 2008) would have compacted and partially shut the permeability of 678 679 the conduit margin, with the shear zone creating an impermeable barrier preventing gas from escaping 680 to the surrounding country rock and promoting outgassing through the more permeable conduit core, at 681 least up to the point of rupture (cf. Collinson and Neuberg, 2012). Upon further ascent, changes in the 682 stress fields and physical properties of the magmas during pulsatory ascent would have favoured 683 transition to a macroscopically brittle response to shear (Lavallée and Kendrick, 2020), triggering 684 seismic rupture (Umakoshi et al., 2008; Lamb et al., 2015) and initiation of predominantly fault-685 controlled, stick-slip dynamics in the final stint of magma ascent and spine extrusion (Hornby et al., 686 2015). In brief periods of high discharge rate, shear may have localised along the primary seismogenic 687 fault, simultaneously creating a Riedel fracture, but in periods with lower discharge rates, shear would 688 have been distributed over a wide area and the fault would become inactive (stick phase), shifting the 689 Riedel fracture to shallower depth; upon renewed discharge rate increase, shear would narrow again, 690 and faulting would generate another Riedel fracture, and so on (Fig. 9a). Indeed, using seismic events 691 as a proxy for the ductile-brittle transition it was possible to identify its migration through time as the 692 inclined spine loaded and compacted its lower shear zone as it grew, dilating the upper fault zone (Lamb 693 et al., 2015). This is further indicated by the localisation of fumaroles along the upper spine margin 694 (also observed during our latest field campaign in 2016), showing that the fault zone around the inclined 695 spine controlled fluid circulation in the upper conduit (Lamb et al., 2015; Yamasato, 1998). Finally, a 696 late lateral shift in dilational shearing, from the conduit margin to the conduit core, suggest that the location of shear may migrate during magma ascent in conduits as a result of changes in local stresses 697 698 (e.g., upon extrusion and/ or blocks calving), likely resulting from a combination of pore pressure 699 fluctuations, strain rate reduction and progressive inward cooling which would have favoured 700 deformation in the core of the spine. Thus, the rheology of magma and the dominant shearing mode 701 may evolve during ascent, which in turn dynamically modifies the permeability distribution across the 702 conduit through time (Fig. 9a).





704 Rheological assessment of magma switching from ductile to brittle deformation

705 The above rheological description is primarily based on the unavoidable decompression of erupting 706 magma (which degases, crystallises and viscously stiffens), yet previous observations at Unzen suggest 707 that the conditions for magmatic flow may have fluctuated (Umakoshi et al., 2008; Lamb et al., 2015), thus contributing to rheological shifts. Here, we invoke findings from the literature to assess the 708 709 conditions leading to rupture. The discharge rates associated with spine extrusion in 1994-95 varied, 710 although Yamashina et al. (1999) constrained a relatively constant spine protrusion rate of 0.8 m d⁻¹ 711 over a week-long period in early November 1994. Scrutinising within this period, however, seismicity 712 indicated a pulsatory magma ascent in the conduit at shorter timescales (Umakoshi et al. 2008; Lamb 713 et al. 2015). In particular, waveform correlation of the seismic record performed by Lamb et al. (2015) 714 revealed rhythmic seismicity punctuated by two primary clusters that were attributed to recurring 715 rupture associated with stick-slip cycles. They identified 668 repetitive events over the course of the 36 716 days examined: 487 from cluster 1 and 181 from cluster 2. Progressive shallowing of cluster 1 source 717 location was argued to result from progressive compaction of the lower shear zone (underneath the 718 inclined magma column) as eruption slowly waned; in contrast, cluster 2, which was accompanied by 719 low-frequency coda associated with fluid resonance, showed deepening of source location due to 720 dilation on the overside of the inclined conduit. Considering the events in cluster 1, we define the 721 recurrence rate of fault slip at 13.5 events per day; so each 'stick' interval for viscous flow would have 722 lasted on average 106 minutes. Hornby et al. (2015) statistically analysed the slip duration of seismic 723 events in clusters 1 and 2, defining a mode and mean of 0.1 s. In order to pursue a quantitative analysis 724 of stick-slip behaviour, we must first turn our attention to our knowledge of Unzen magma flow and 725 failure conditions.

726 Coats et al. (2018) studied the rheology of Unzen's porous lavas to define a failure criterion. 727 Considering the estimated eruptive temperature of ca. 870-900 °C (Holtz et al., 2005; Venezky and 728 Rutherford, 1999) and measured glass transition temperature (at 10 °C min⁻¹) of 790 °C (Wallace et al., 729 2019), Coats et al. (2018) empirically defined that Unzen magma would break if experiencing strain 730 rates exceeding $\sim 10^{-3}$ s⁻¹; otherwise, magma would undergo ductile flow. But these determination were done at atmospheric pressure, so the melt was considered dry; Kusakabe et al. (1999) determined the 731 732 concentration of magmatic water dissolved in the groundmass glass of eruptive products at 0.1-0.5 wt. 733 %; however, the concentration of dissolved water at the point of rupture, at 500 m depth or ~10 MPa 734 pressure considering a nominal rock density of ~2,000 kg m⁻³ (Scheu et al., 2006), would have been ~1 735 wt. % (Liu et al., 2005). Such a higher concentration would lower the viscosity of the interstitial melt 736 one order of magnitude; as the strain rate limit shares an inverse relationship with viscosity (e.g., 737 Dingwell and Webb, 1989), we advance that the presence of dissolved water in the melt would have 738 shifted the strain rate limit by approximately one order of magnitude. If we omit any upscaling of the 739 above failure conditions for simplification and assume that deformation was localised in the ~ 1 m-wide 740 high shear area of the spine, rupture would have occurred when the ascent rate exceeded 1 mm.s⁻¹. As 741 such high deformation rate episodes are inferred to have triggered fault slip events lasting on average 742 0.1 s (Hornby et al. 2015), each slip event may have resulted in a mere ≥ 0.1 mm of displacement. With 743 13.5 events per day, this would culminate in $\gtrsim 1.35$ mm of magma ascent ascribed to faulting activity, 744 signifying that deformation associated with the ~0.8 m daily ascent was predominantly ductile and 745 aseismic.

746 We can then turn our attention to geometrical constraints from our structural analysis to frame magma 747 ascent conditions that satisfy the above failure criterion. The Riedel fractures that are observed at regular 748 intervals of ~4.5 cm in the high shear zones have been shown to be important stress and strain rate 749 distribution markers in multiphase materials containing a weak phase, such as melt and bubbles (Finch





750 et al., 2020), and can thus be used to constrain rates. Considering the ephemeral nature of Riedel fracture development (Finch et al., 2020), here we assume that their formation may be encouraged during brief 751 periods of high strain rate, and they thus portray the clockwork ticking of seismogenic slip events during 752 753 magma ascent. Bearing in mind an average spacing of 4.5 cm and an angle of 57° with respect to the 754 main C-S fabric, we estimate the offset of the loci of rupture events at 5.4 cm. Recalling the 0.1 mm of 755 displacement ascribed to faulting events (detailed in the previous paragraph), this suggests that ductile 756 deformation was responsible for 5.3 cm of magma ascent during inter-seismic periods (i.e., inter-757 seismicity deformation, ISD; Fig 9a). Again, considering shear over 1 m area and inter-seismic periods 758 of 106 minutes, we estimate that ductile deformation would have proceeded at an average rate of 8×10^{-10} 759 6 s⁻¹; a value well within the ductile regime as experimentally constrained by Coats et al. (2018). The 760 above rates (of magma flow in the ductile regime and of faulting) may be conservative estimates, 761 especially if we consider the rheological consequences of dissolved water at depth. Even if the threshold 762 strain rate for seismogenic faulting were an order of magnitude higher, at 10^{-2} s⁻¹, this would only require 763 13.5 mm of magma ascent in each brittle faulting event and that inter-seismic periods of ductile 764 deformation at a rate of $\sim 8 \times 10^{-5}$ s⁻¹ would have dominated spine extrusion.

In concert the physical and structural description bolstered by the rheological analysis argue for 765 766 changes in magma rheology during decompression and pulsatory ascent. We propose that throughout 767 its journey to the Earth's surface, magma may undergo several cycles of expansion (from vesiculation 768 and dilation) and collapse (from outgassing and compaction) due to variable permeability and pore 769 pressure, which may promote switches in shearing regimes that trigger further changes in the 770 permeability structure of shallow conduits. For instance, the vesicles of low permeability magma may 771 accumulate fluid, thus reducing the effective pressure and promoting brittle, dilatant rupture; rupture 772 would in turn allow magma outgassing and a reduction in effective pressure, promoting compaction 773 and lowering of permeability; and the cycle may recur. The picture portrayed here highlights the need 774 to understand the coupling between magma and fluid flow dynamics and, importantly, pressure 775 fluctuations (Michaut et al., 2013) in volcanic conduits with increased spatial and temporal complexities 776 in order to resolve the transient state of magma and reconcile gas emission data and volcanic eruption 777 style (Edmonds and Herd, 2007).

778

779 6. Conclusions

780 The present detailed study of the Mount Unzen spine reveals the competing occurrence of 781 compactional and dilational shear regimes during magma ascent in volcanic conduits. At depth, in areas 782 subjected to high effective pressure, shearing may induce pore compaction, thereby lowering the 783 permeability of the system and inhibiting lateral outgassing to the country rock. At shallower depth, 784 where the effective pressure may be low, shearing may favour localised dilation that enhances 785 permeability. Both shear regimes result in the development of permeability anisotropy, with permeability generally being highest parallel or sub-parallel to the direction of extrusion, and lowest 786 perpendicular to the shear plane. The observation of shearing mode overprints suggests that fluctuations 787 788 in effective pressure and strain rates, during stick-slip cycles, may result in magma switching between compactant and dilational shearing regimes, thus dynamically reshaping fluid circulation at a range of 789 790 scales, and in turn controlling outgassing efficiency during magma ascent and eruption.

791

792 Acknowledgements





We are thankful to Guðjón Eggertsson for help with the maintenance of the permeameter. This project
was financially supported by a European Research Council (ERC) Starting Grant on Strain Localisation
in Magma (SLiM, No. 306488) and an award from the DAIWA Anglo-Japanese Foundation (grant No.
11000/11740). YL and JEK acknowledge support from the Leverhulme Trust (ECF-2016-325 and RF2019-526\4, respectively). HT was supported by a University Research Fellowship from the Royal

798 Society.

799

800 References

- 801 Acocella, V.: Hazard mitigation of unstable volcanic edifices, EOS, 91, 2, 2010.
- Alidibirov, M., and Dingwell, D. B.: Magma fragmentation by rapid decompression, Nature, 380, 146 148, 1996.
- Ashwell, P. A., Kendrick, J. E., Lavallée, Y., Kennedy, B. M., Hess, K. U., von Aulock, F. W., Wadsworth,
 F. B., Vasseur, J., and Dingwell, D. B.: Permeability of compacting porous lavas, Journal of Geophysical
- 806 Research-Solid Earth, 120, 1605-1622, 10.1002/2014jb011519, 2015.
- Baker, D. R., Brun, F., O'Shaughnessy, C., Mancini, L., Fife, J. L., and Rivers, M.: A four-dimensional X ray tomographic microscopy study of bubble growth in basaltic foam, Nature Communications, 3,
 10.1038/ncomms2134, 2012.
- Ball, J. L., Stauffer, P. H., Calder, E. S., and Valentine, G. A.: The hydrothermal alteration of cooling lava
 domes, Bulletin of Volcanology, 77, 10.1007/s00445-015-0986-z, 2015.
- 812 Blower, J. D.: Factors controlling permeability-porosity relationships in magma, Bulletin of 813 Volcanology, 63, 497-504, 2001.
- Blundy, J., Cashman, K., and Humphreys, M.: Magma heating by decompression-driven crystallization
 beneath andesite volcances, Nature, 443, 76-80, 10.1038/nature05100, 2006.
- Brace, W. F., Walsh, J. B., and Frangos, W. T.: Permeability of granite under high pressure, Journal of
 Geophysical Research, 73, 2225-&, 10.1029/JB073i006p02225, 1968.
- Browning, J., Tuffen, H., James, M. R., Owen, J., Castro, J. M., Halliwell, S., and Wehbe, K.: Postfragmentation vesiculation timescales in hydrous rhyolitic bombs from Chaitén volcano, Journal of
- 820 South American Earth Sciences, 104, 102807, <u>https://doi.org/10.1016/j.jsames.2020.102807</u>, 2020.
- Burgisser, A., Chevalier, L., Gardner, J. E., and Castro, J. M.: The percolation threshold and permeability
 evolution of ascending magmas, Earth and Planetary Science Letters, 470, 37-47,
 10.1016/j.epsl.2017.04.023, 2017.
- Burgisser, A., Bechon, T., Chevalier, L., Collombet, M., Arbaret, L., and Forien, M.: Conduit processes
 during the February 11, 2010 Vulcanian eruption of Soufriere Hills, Montserrat, Journal of Volcanology
 and Geothermal Research, 373, 23-35, 10.1016/j.jvolgeores.2019.01.020, 2019.
- Caricchi, L., Burlini, L., Ulmer, P., Gerya, T., Vassalli, M., and Papale, P.: Non-Newtonian rheology of
 crystal-bearing magmas and implications for magma ascent dynamics, Earth and Planetary Science
 Letters, 264, 402-419, 2007.
- Castro, J. M., and Gardner, J. E.: Did magma ascent rate control the explosive-effusive transition at the
 Inyo volcanic chain, California?, Geology, 36, 279-282, 10.1130/g24453a.1, 2008.
- Castro, J. M., Cordonnier, B., Tuffen, H., Tobin, M. J., Puskar, L., Martin, M. C., and Bechtel, H. A.: The
 role of melt-fracture degassing in defusing explosive rhyolite eruptions at volcan Chaiten, Earth and
 Planetary Science Letters, 333, 63-69, 10.1016/j.epsl.2012.04.024, 2012.
- Célarié, F., Prades, S., Bonamy, D., Ferrero, L., Bouchaud, E., Guillot, C., and Marliere, C.: Glass breaks
- like metal, but at the nanometer scale, Physical Review Letters, 90, 10.1103/PhysRevLett.90.075504,
 2003.
- 838 Chouet, B. A.: Long-period volcano seismicity: Its source and use in eruption forecasting, Nature, 380,
- 839 309-316, 1996.





- Coats, R., Kendrick, J. E., Wallace, P. A., Miwa, T., Hornby, A. J., Ashworth, J. D., Matsushima, T., and
- 841 Lavallée, Y.: Failure criteria for porous dome rocks and lavas: a study of Mt. Unzen, Japan, Solid Earth,
- 842 9, 1299-1328, 10.5194/se-9-1299-2018, 2018.
- Collinson, A. S. D., and Neuberg, J. W.: Gas storage, transport and pressure changes in an evolving permeable volcanic edifice, Journal of Volcanology and Geothermal Research, 243, 1-13,
- 845 10.1016/j.jvolgeores.2012.06.027, 2012.
- 846 Colombier, M., Wadsworth, F. B., Gurioli, L., Scheu, B., Kueppers, U., Di Muro, A., and Dingwell, D. B.:
- The evolution of pore connectivity in volcanic rocks, Earth and Planetary Science Letters, 462, 99-109,
 10.1016/j.epsl.2017.01.011, 2017.
- 849 Cordonnier, B., Hess, K. U., Lavallée, Y., and Dingwell, D. B.: Rheological properties of dome lavas: Case
- 850 study of Unzen volcano, Earth and Planetary Science Letters, 279, 263-272, 851 10.1016/j.epsl.2009.01.014, 2009.
- Cordonnier, B., Caricchi, L., Pistone, M., Castro, J., Hess, K. U., Gottschaller, S., Manga, M., Dingwell,
 D. B., and Burlini, L.: The viscous-brittle transition of crystal-bearing silicic melt: Direct observation of
- 854 magma rupture and healing, Geology, 40, 611-614, 10.1130/g3914.1, 2012.
- Degruyter, W., Bachmann, O., Burgisser, A., and Manga, M.: The effects of outgassing on the transition
 between effusive and explosive silicic eruptions, Earth and Planetary Science Letters, 349, 161-170,
 10.1016/j.epsl.2012.06.056, 2012.
- Dingwell, D. B., and Webb, S. L.: Structural relaxation in silicate melts and non-Newtonian melt
 rheology in geologic processes, Physics and Chemistry of Minerals, 16, 508-516, 1989.
- Bingwell, D. B., and Webb, S. L.: Relaxation in silicate melts, European Journal of Mineralogy, 2, 427449, 1990.
- B62 Dingwell, D. B.: Volcanic dilemma: flow or blow?, Science, 273, 1054-1055, 1996.
- 863 Dingwell, D. B., Lavallée, Y., Hess, K. U., Flaws, A., Martí, J., Nichols, A. R. L., Gilg, H. A., and Schillinger,
- B.: Eruptive shearing of tube pumice: pure and simple, Solid Earth, 7, 1383-1393, 10.5194/se-7-13832016, 2016.
- Edmonds, M., Oppenheimer, C., Pyle, D. M., Herd, R. A., and Thompson, G.: SO2 emissions from
 Soufriere Hills Volcano and their relationship to conduit permeability, hydrothermal interaction and
 degassing regime, Journal of Volcanology and Geothermal Research, 124, 23-43, 10.1016/s03770273(03)00041-6, 2003.
- Edmonds, M., and Herd, R. A.: A volcanic degassing event at the explosive-effusive transition,
 Geophysical Research Letters, 34, 10.1029/2007gl031379, 2007.
- Eggertsson, G. H., Lavallée, Y., Kendrick, J. E., and Markússon, S. H.: Improving fluid flow in geothermal
 reservoirs by thermal and mechanical stimulation: The case of Krafla volcano, Iceland, Journal of
- 874 Volcanology and Geothermal Research, in press, 1-14, 2018.
- Eichelberger, J. C., Carrigan, C. R., Westrich, H. R., and Price, R. H.: Non-explosive silicic volcanism,
 Nature, 323, 598-602, 10.1038/323598a0, 1986.
- Farquharson, J., Heap, M. J., Varley, N. R., Baud, P., and Reuschle, T.: Permeability and porosity
 relationships of edifice-forming andesites: A combined field and laboratory study, Journal of
 Volcanology and Geothermal Research, 297, 52-68, 10.1016/j.jvolgeores.2015.03.016, 2015.
- Farquharson, J., Heap, M. J., Baud, P., Reuschle, T., and Varley, N. R.: Pore pressure embrittlement in
 a volcanic edifice, Bulletin of Volcanology, 78, 10.1007/s00445-015-0997-9, 2016a.
- Farquharson, J. I., Heap, M. J., and Baud, P.: Strain-induced permeability increase in volcanic rock,
 Geophysical Research Letters, 43, 11603-11610, 10.1002/2016gl071540, 2016b.
- Farquharson, J. I., Heap, M. J., Lavallée, Y., Varley, N. R., and Baud, P.: Evidence for the development
- of permeability anisotropy in lava domes and volcanic conduits, Journal of Volcanology and Geothermal Research, 323, 163-185, 10.1016/j.jvolgeores.2016.05.007, 2016c.
- 887 Farquharson, J. I., Wadsworth, F. B., Heap, M. J., and Baud, P.: Time-dependent permeability evolution
- 888 in compacting volcanic fracture systems and implications for gas overpressure, Journal of Volcanology
- and Geothermal Research, 339, 81-97, 10.1016/j.jvolgeores.2017.04.025, 2017.





- Finch, M. A., Bons, P. D., Steinbach, F., Griera, A., Llorens, M.-G., Gomez-Rivas, E., Ran, H., and de
 Riese, T.: The ephemeral development of C' shear bands: A numerical modelling approach, Journal of
 Chrustward Conference 120, 104001, 10, 10401, 2020.
- 892 Structural Geology, 139, 104091, 10.1016/j.jsg.2020.104091, 2020.
- Fossen, H., and Cavalcante, G. C. G.: Shear zones A review, Earth-Science Reviews, 171, 434-455,
 0.1016/j.earscirev.2017.05.002, 2017.
- Gaunt, H. E., Sammonds, P. R., Meredith, P. G., Smith, R., and Pallister, J. S.: Pathways for degassing
 during the lava dome eruption of Mount St. Helens 2004-2008, Geology, 42, 947-950,
 10.1130/g35940.1, 2014.
- Gaunt, H. E., Sammonds, P. R., Meredith, P. G., and Chadderton, A.: Effect of temperature on the
 permeability of lava dome rocks from the 2004-2008 eruption of Mount St. Helens, Bulletin of
 Volcanology, 78, 10.1007/s00445-016-1024-5, 2016.
- 901 Gonnermann, H. M., and Manga, M.: The fluid mechanics inside a volcano, Annual Review of Fluid 902 Mechanics, 39, 321-356, 2007.
- Goto, A.: A new model for volcanic earthquake at Unzen Volcano: Melt rupture model, Geophysical
 Research Letters, 26, 2541-2544, 1999.
- Goto, A., Fukui, K., Hiraga, T., Nishida, Y., Ishibashi, H., Matsushima, T., Miyamoto, T., and Sasaki, O.:
 Rigid migration of Unzen lava rather than flow, Journal of Volcanology and Geothermal Research, 407,
- 907 10.1016/j.jvolgeores.2020.107073, 2020.
- Hale, A. J., and Wadge, G.: The transition from endogenous to exogenous growth of lava domes with
 the development of shear bands, Journal of Volcanology and Geothermal Research, 171, 237-257,
 2008.
- Heap, M. J., Farquharson, J. I., Baud, P., Lavallée, Y., and Reuschle, T.: Fracture and compaction of
 andesite in a volcanic edifice, Bulletin of Volcanology, 77, 10.1007/s00445-015-0938-7, 2015a.
- Heap, M. J., Kennedy, B. M., Pernin, N., Jacquemard, L., Baud, P., Farguharson, J. I., Scheu, B., Lavallee,
- 914 Y., Gilg, H. A., Letham-Brake, M., Mayer, K., Jolly, A. D., Reuschle, T., and Dingwell, D. B.: Mechanical 915 behaviour and failure modes in the Whakaari (White Island volcano) hydrothermal system, New of 916 Zealand, Journal Volcanology and Geothermal Research, 295, 26-42, 917 10.1016/j.jvolgeores.2015.02.012, 2015b.
- Heap, M. J., and Kennedy, B. M.: Exploring the scale-dependent permeability of fractured andesite,
 Earth and Planetary Science Letters, 447, 139-150, 10.1016/j.epsl.2016.05.004, 2016.
- Heap, M. J., Kennedy, B. M., Farquharson, J. I., Ashworth, J., Mayer, K., LEtham-Brake, M., Reuschlé,
 T., Gilg, H. A., Scheu, B., Lavallée, Y., Siratovich, P. A., Cole, J. W., Jolly, A. D., Baud, P., and Dingwell, D.
- B.: A multidisciplinary approach to quantify the permeability of the Whakaari/ White Island volcanic
- hydrothermal system (Taupo Volcanic Zone, New Zealand), Journal of Volcanology and Geothermal
- 924 Research, 10.1016/j.jvolgeores.2016.12.004, 2017a.
- Heap, M. J., Violay, M., Wadsworth, F. B., and Vasseur, J.: From rock to magma and back again: The
 evolution of temperature and deformation mechanism in conduit margin zones, Earth and Planetary
 Science Letters, 463, 92-100, 10.1016/j.epsl.2017.01.021, 2017b.
- 928 Heap, M. J., Troll, V. R., Kushnir, A. R. L., Gilg, H. A., Collinson, A. S. D., Deegan, F. M., Darmawan, H.,
- 929 Seraphine, N., Neuberg, J., and Walter, T. R.: Hydrothermal alteration of andesitic lava domes can lead
- to explosive volcanic behaviour, Nature Communications, 10, 5063, 10.1038/s41467-019-13102-8,
 2019.
- Heap, M. J., and Violay, M. E. S.: The mechanical behaviour and failure modes of volcanic rocks: a
 review, Bulletin of Volcanology, 83, 10.1007/s00445-021-01447-2, 2021.
- Heiken, G., Wohletz, K., and Eichelberger, J.: Fracture fillings and intrusive pyroclasts, Inyo domes,
 California, Journal of Geophysical Research-Solid Earth and Planets, 93, 4335-4350,
 10.1029/JB093iB05p04335, 1988.
- 937 Holland, A. S. P., Watson, I. M., Phillips, J. C., Caricchi, L., and Dalton, M. P.: Degassing processes during
- 938 lava dome growth: Insights from Santiaguito lava dome, Guatemala, Journal of Volcanology and
- 939 Geothermal Research, 202, 153-166, 10.1016/j.jvolgeores.2011.02.004, 2011.





Holtz, F., Sato, H., Lewis, J., Behrens, H., and Nakada, S.: Experimental petrology of the 1991-1995
 Unzen dacite, Japan. Part I: Phase relations, phase composition and pre-eruptive conditions, Journal

- 942 of Petrology, 46, 319-337, 10.1093/petrology/egh077, 2005.
- 943 Hornby, A. J., Kendrick, J. E., Lamb, O. D., Hirose, T., De Angelis, S., von Aulock, F. W., Umakoshi, K.,
- Miwa, T., Henton De Angelis, S., Wadsworth, F. B., Hess, K.-U., Dingwell, D. B., and Lavallée, Y.: Spine
 growth and seismogenic faulting at Mt. Unzen, Japan, Journal of Geophysical Research: Solid Earth,
 120, 2169-9356, 10.1002/2014JB011660, 2015.
- Hornby, A. J., Lavallée, Y., Kendrick, J. E., De Angelis, S., Lamur, A., Rietbrock, A., and Chigna, G.: Brittleductile deformation and tensile rupture of dome lava during inflation at Santiaguito, Guatemala,
- Journal of Geophysical Research, in press, 10.1029/2018JB017253, 2019.
- 950 Jaupart, C., and Allègre, C. J.: Gas content, eruption rate and instabilities or eruption regime in silicic
- volcanoes, Earth and Planetary Science Letters, 102, 413-429, 10.1016/0012-821x(91)90032-d, 1991.
 Kendrick, J. E., Lavallée, Y., Ferk, A., Perugini, D., Leonhardt, R., and Dingwell, D. B.: Extreme frictional
- processes in the volcanic conduit of Mount St. Helens (USA) during the 2004-2008 eruption, Journal
 of Structural Geology, 38, 61-76, 10.1016/j.jsg.2011.10.003, 2012.
- Kendrick, J. E., Lavallée, Y., Hess, K. U., Heap, M. J., Gaunt, H. E., Meredith, P. G., and Dingwell, D. B.:
 Tracking the permeable porous network during strain-dependent magmatic flow, Journal of
 Volcanology and Geothermal Research, 260, 117-126, 10.1016/j.jvolgeores.2013.05.012, 2013.
- Kendrick, J. E., Lavallée, Y., Hess, K. U., De Angelis, S., Ferk, A., Gaunt, H. E., Meredith, P. G., Dingwell,
 D. B., and Leonhardt, R.: Seismogenic frictional melting in the magmatic column, Solid Earth, 5, 199-
- 960 208, 10.5194/se-5-199-2014, 2014a.

Kendrick, J. E., Lavallée, Y., Hirose, T., Di Toro, G., Hornby, A. J., De Angelis, S., and Dingwell, D. B.:
Volcanic drumbeat seismicity caused by stick-slip motion and magmatic frictional melting, Nature
Geoscience, 7, 438-442, 10.1038/ngeo2146, 2014b.

Kendrick, J. E., Lavallée, Y., Varley, N. R., Wadsworth, F. B., Lamb, O. D., and Vasseur, J.: Blowing off
steam: Tuffisite formation as a regulator for lava dome eruptions, Frontiers in Earth Science, 4,
10.3389/feart.2016.00041, 2016.

- Kendrick, J. E., Schaefer, L. N., Schauroth, J., Bell, A. F., Lamb, O. D., Lamur, A., Miwa, T., Coats, R.,
 Lavallée, Y., and Kennedy, B. M.: Physical and mechanical rock properties of a heterogeneous volcano:
 the case of Mount Unzen, Japan, Solid Earth, 12, 633-664, 10.5194/se-12-1-2021, 2021.
- Kennedy, B. M., Wadsworth, F. B., Vasseur, J., Schipper, C. I., Jellinek, A. M., von Aulock, F. W., Hess,
- K. U., Russell, J. K., Lavallee, Y., Nichols, A. R. L., and Dingwell, D. B.: Surface tension driven processes
 densify and retain permeability in magma and lava, Earth and Planetary Science Letters, 433, 116-124,
 10.1016/j.epsl.2015.10.031, 2016.
- Kennedy, L. A., and Russell, J. K.: Cataclastic production of volcanic ash at Mount Saint Helens, Physics
 and Chemistry of the Earth, 45-46, 40-49, 10.1016/j.pce.2011.07.052, 2012.
- Klug, C., and Cashman, K. V.: Permeability development in vesiculating magmas: Implications for
 fragmentation, Bulletin of Volcanology, 58, 87-100, 10.1007/s004450050128, 1996.
- Kolzenburg, S., Heap, M. J., Lavallée, Y., Russell, J. K., Meredith, P. G., and Dingwell, D. B.: Strength and
 permeability recovery of tuffisite-bearing andesite, Solid Earth, 3, 191-198, 10.5194/se-3-191-2012,
- 980 2012.
- Kueppers, U., Scheu, B., Spieler, O., and Dingwell, D. B.: Field-based density measurements as tool to
 identify preeruption dome structure: set-up and first results from Unzen volcano, Japan, Journal of
 Volcanology and Geothermal Research, 141, 65-75, 2005.
- Kusakabe, M., Sato, H., Nakada, S., and Kitamura, T.: Water contents and hydrogen isotopic ratios of
 rocks and minerals from the 1991 eruption of Unzen volcano, Japan, Journal of Volcanology and
 Geothermal Research, 89, 231-242, 10.1016/s0377-0273(98)00134-6, 1999.
- 987 Kushnir, A. R. L., Martel, C., Bourdier, J. L., Heap, M. J., Reuschle, T., Erdmann, S., Komorowski, J. C.,
- 988 and Cholik, N.: Probing permeability and microstructure: Unravelling the role of a low-permeability
- 989 dome on the explosivity of Merapi (Indonesia), Journal of Volcanology and Geothermal Research, 316,
- 990 56-71, 10.1016/j.jvolgeores.2016.02.012, 2016.





- Kushnir, A. R. L., Martel, C., Champallier, R., and Wadsworth, F. B.: Permeability Evolution in Variably
 Glassy Basaltic Andesites Measured Under Magmatic Conditions, Geophysical Research Letters, 44,
- . 993 10262-10271, 10.1002/2017gl074042, 2017.
- Lamb, O. D., De Angelis, S., Umakoshi, K., Hornby, A. J., Kendrick, J. E., and Lavallée, Y.: Repetitive
 fracturing during spine extrusion at Unzen volcano, Japan, Solid Earth, 6, 1277-1293, 10.5194/se-61277-2015, 2015.
- Barner, A., Kendrick, J. E., Eggertsson, G. H., Wall, R. J., Ashworth, J. D., and Lavallée, Y.: The
 permeability of fractured rocks in pressurised volcanic and geothermal systems, Scientific Reports,
 2017.
- Lamur, A., Kendrick, J. E., Wadsworth, F. B., and Lavallée, Y.: Fracture healing and strength recovery in magmatic liquids, Geology, 47, 195-198, 10.1130/g45512.1, 2019.
- Laumonier, M., Arbaret, L., Burgisser, A., and Champallier, R.: Porosity redistribution enhanced by strain localization in crystal-rich magmas, Geology, 39, 715-718, 10.1130/g31803.1, 2011.
- Lavallée, Y., Hess, K.-U., Cordonnier, B., and Dingwell, D. B.: Non-Newtonian rheological law for highly
 crystalline dome lavas, Geology, 35, 843-846, 10.1130/g23594a.1, 2007.
- Lavallée, Y., Meredith, P. G., Dingwell, D. B., Hess, K. U., Wassermann, J., Cordonnier, B., Gerik, A., and
 Kruhl, J. H.: Seismogenic lavas and explosive eruption forecasting, Nature, 453, 507-510,
 10.1038/nature06980, 2008.
- 1009 Lavallée, Y., Varley, N. R., Alatorre-Ibargueengoitia, M. A., Hess, K. U., Kueppers, U., Mueller, S.,
- 1010 Richard, D., Scheu, B., Spieler, O., and Dingwell, D. B.: Magmatic architecture of dome-building
- 1011 eruptions at Volcan de Colima, Mexico, Bulletin of Volcanology, 74, 249-260, 10.1007/s00445-0111012 0518-4, 2012.
- Lavallée, Y., Benson, P. M., Heap, M. J., Hess, K.-U., Flaws, A., Schillinger, B., Meredith, P. G., and Dingwell, D. B.: Reconstructing magma failure and the degassing network of dome-building eruptions,
- 1015 Geology, 41, 515-518, 10.1130/g33948.1, 2013.
- Lavallée, Y., Dingwell, D. B., Johnson, J. B., Cimarelli, C., Hornby, A. J., Kendrick, J. E., von Aulock, F. W.,
 Kennedy, B. M., Andrews, B. J., Wadsworth, F. B., Rhodes, E., and Chigna, G.: Thermal vesiculation
 during volcanic eruptions, Nature, 528, 544-547, 10.1038/nature16153, 2015.
- Lavallée, Y., and Kendrick, J. E.: A review of the physical and mechanical properties of volcanic rocks
 and magmas in the brittle and ductile regimes, in: Forecasting and planning for volcanic hazards, risks,
 and disasters. Vol. 2, 2nd Edition ed., edited by: Papale, P., Elsevier, 2020.
- Lavallée, Y., and Kendrick, J. E.: Strain localisation in magmas, in: Magmas, Melts, ILquids and Glasses:
 Experimental Insights edited by: Neuville, D. R., Henderson, G. S., and Dingwell, D. B., Reviews in
 Minerlogy and Geochemistry, Mineralogical Society of America, 2021.
- Lejeune, A. M., and Richet, P.: Rheology of Crystal-Bearing Silicate Melts an Experimental-Study at
 High Viscosities, Journal of Geophysical Research-Solid Earth, 100, 4215-4229, 1995.
- Lejeune, A. M., Bottinga, Y., Trull, T. W., and Richet, P.: Rheology of bubble-bearing magmas, Earthand Planetary Science Letters, 166, 71-84, 1999.
- Liu, Y., Zhang, Y. X., and Behrens, H.: Solubility of H2O in rhyolitic melts at low pressures and a new
 empirical model for mixed H2O-CO2 solubility in rhyolitic melts, Journal of Volcanology and
 Geothermal Research, 143, 219-235, 10.1016/j.jvolgeores.2004.09.019, 2005.
- Loaiza, S., Fortin, J., Schubnel, A., Gueguen, Y., Vinciguerra, S., and Moreira, M.: Mechanical behavior
 and localized failure modes in a porous basalt from the Azores, Geophysical Research Letters, 39,
 10.1029/2012gl053218, 2012.
- Mader, H. M., Llewellin, E. W., and Mueller, S. P.: The rheology of two-phase magmas: A review and
 analysis, Journal of Volcanology and Geothermal Research, 257, 135-158,
 10.1016/j.jvolgeores.2013.02.014, 2013.
- 1038 Matoza, R. S., and Chouet, B. A.: Subevents of long-period seismicity: Implications for hydrothermal 1039 dynamics during the 2004-2008 eruption of Mount St. Helens, Journal of Geophysical Research-Solid
- 1040 Earth, 115, 10.1029/2010jb007839, 2010.
- 1041 Melnik, O., and Sparks, R. S. J.: Nonlinear dynamics of lava dome extrusion, Nature, 402, 37-41, 1999.





1042 Michaut, C., Ricard, Y., Bercovici, D., and Sparks, R. S. J.: Eruption cyclicity at silicic volcanoes 1043 potentially caused by magmatic gas waves, Nature Geoscience, 6, 856-860, 10.1038/ngeo1928, 2013. Mueller, S., Melnik, O., Spieler, O., Scheu, B., and Dingwell, D. B.: Permeability and degassing of dome 1044 1045 lavas undergoing rapid decompression: An experimental determination, Bulletin of Volcanology, 67, 1046 526-538. 2005. 1047 Mueller, S., Scheu, B., Spieler, O., and Dingwell, D. B.: Permeability control on magma fragmentation, 1048 Geology, 36, 399-402, 10.1130/g24605a.1, 2008. 1049 Nakada, S., Miyake, Y., Sato, H., Oshima, O., and Fujinawa, A.: Endogenous growth of dacite dome at Geology, 1050 Unzen volcano (Japan), 1993-1994, 23, 157-160, 10.1130/0091-1051 7613(1995)023<0157:egodda>2.3.co;2, 1995. Nakada, S., and Motomura, Y.: Petrology of the 1991-1995 eruption at Unzen: effusion pulsation and 1052 1053 groundmass crystallization, Journal of Volcanology and Geothermal Research, 89, 173-196, 1054 10.1016/s0377-0273(98)00131-0, 1999. Nakada, S., Shimizu, H., and Ohta, K.: Overview of the 1990-1995 eruption at Unzen Volcano, Journal 1055 1056 of Volcanology and Geothermal Research, 89, 1-22, 10.1016/s0377-0273(98)00118-8, 1999. 1057 Navon, O., Chekhmir, A., and Lyakhovsky, V.: Bubble growth in highly viscous melts: theory, 1058 experiments, and autoexplosivity of dome lavas, Earth and Planetary Science Letters, 160, 763-776, 1059 10.1016/s0012-821x(98)00126-5, 1998. 1060 Neuberg, J. W., Tuffen, H., Collier, L., Green, D., Powell, T., and Dingwell, D.: The trigger mechanism of 1061 low-frequency earthquakes on Montserrat, Journal of Volcanology and Geothermal Research, 153, 37-1062 50, 2006. 1063 Newhall, C. G., and Melson, W. G.: Explosive activity associated with the growth of volcanic domes, 1064 Journal of Volcanology and Geothermal Research, 17, 111-131, 10.1016/0377-0273(83)90064-1, 1983. 1065 Ohba, T., Hirabayashi, J.-I., Nogami, K., Kusakabe, M., and Yoshida, M.: Magma degassing process 1066 during the eruption of Mt. Unzen, Japan in 1991 to 1995: Modeling with the chemical composition of 1067 volcanic gas, Journal of Volcanology and Geothermal Research, 175, 120-132, 1068 10.1016/j.jvolgeores.2008.03.040, 2008. 1069 Okumura, S., Nakamura, M., and Tsuchiyama, A.: Shear-induced bubble coalescence in rhyolitic melts 1070 with low vesicularity, Geophysical Research Letters, 33, 10.1029/2006gl027347, 2006. 1071 Okumura, S., Nakamura, M., Tsuchiyama, A., Nakano, T., and Uesugi, K.: Evolution of bubble microstructure in sheared rhyolite: Formation of a channel-like bubble network, Journal of 1072 1073 Geophysical Research-Solid Earth, 113, 10.1029/2007jb005362, 2008. 1074 Okumura, S., Nakamura, M., Takeuchi, S., Tsuchiyama, A., Nakano, T., and Uesugi, K.: Magma deformation may induce non-explosive volcanism via degassing through bubble networks, Earth and 1075 1076 Planetary Science Letters, 281, 267-274, 10.1016/j.epsl.2009.02.036, 2009. 1077 Okumura, S., Nakamura, M., Nakano, T., Uesugi, K., and Tsuchiyama, A.: Shear deformation 1078 experiments on vesicular rhyolite: Implications for brittle fracturing, degassing, and compaction of 1079 magmas in volcanic conduits, Journal of Geophysical Research-Solid Earth, 115, 1080 10.1029/2009jb006904, 2010. 1081 Okumura, S., Nakamura, M., Nakano, T., Uesugi, K., and Tsuchiyama, A.: Experimental constraints on 1082 permeable gas transport in crystalline silicic magmas, Contributions to Mineralogy and Petrology, 164, 1083 493-504, 10.1007/s00410-012-0750-8, 2012. 1084 Okumura, S., Nakamura, M., Uesugi, K., Nakano, T., and Fujioka, T.: Coupled effect of magma degassing 1085 and rheology on silicic volcanism, Earth and Planetary Science Letters, 362, 163-170, 1086 10.1016/j.epsl.2012.11.056, 2013. 1087 Okumura, S., and Sasaki, O.: Permeability reduction of fractured rhyolite in volcanic conduits and its 1088 control on eruption cyclicity, Geology, 42, 843-846, 10.1130/g35855.1, 2014. 1089 Pallister, J. S., Cashman, K. V., Hagstrum, J. T., Beeler, N. M., Moran, S. C., and Denlinger, R. P.: Faulting 1090 within the Mount St. Helens conduit and implications for volcanic earthquakes, Geological Society of

1091 America Bulletin, 125, 359-376, 10.1130/b30716.1, 2013a.





- Pallister, J. S., Diefenback, A. K., Burton, W. C., Muñoz, J., Griswold, J. P., Lara, L. E., Lowernster, J. B.,
 and Valenzuela, C. E.: The Chaitén rhyolite lava dome: Eruption sequence, lava dome volumes, rapid
- 1094 effusion rates and source of the rhyolite magma, Andean Geology, 40, 277-294, 2013b.
- 1095 Paterson, M. S., and Wong, T.-F.: Experimental Rock Deformation- The Brittle Field., Science-1096 Technology, 347p, 2005.
- Pistone, M., Caricchi, L., Ulmer, P., Burlini, L., Ardia, P., Reusser, E., Marone, F., and Arbaret, L.:
 Deformation experiments of bubble- and crystal-bearing magmas: Rheological and microstructural
 analysis, Journal of Geophysical Research-Solid Earth, 117, 10.1029/2011jb008986, 2012.
- 1100 Platz, T., Cronin, S. J., Procter, J. N., Neal, V. E., and Foley, S. F.: Non-explosive, dome-forming eruptions
- 1101 at Mt. Taranaki, New Zealand, Geomorphology, 136, 15-30, 10.1016/j.geomorph.2011.06.016, 2012.
- 1102 Radon, J.: On the determination of functions from their integral values along certain manifolds, leee
- 1103 Transactions on Medical Imaging, 5, 170-176, 10.1109/tmi.1986.4307775, 1986.
- 1104 Ramsay, J. G.: Shear zone geometry: A review, Journal of Structural Geology, 2, 83-99, 10.1016/01911105 8141(80)90038-3, 1980.
- Rhodes, E., Kennedy, B. M., Lavallée, Y., Hornby, A., Edwards, M., and Chigna, G.: Textural Insights Into
 the Evolving Lava Dome Cycles at Santiaguito Lava Dome, Guatemala, Frontiers in Earth Science, 6,
 10.3389/feart.2018.00030, 2018.
- 1109 Rohnacher, A., Rietbrock, A., Gottschämmer, E., Carter, W., Lavallée, Y., De Angelis, S., Kendrick, J. E.,
- and Chigna, G.: Source mechanism of seismic explosion signals at Santiaguito volcano, Guatemala:
 New insights from seismic analysis and numerical modeling, Frontiers in Earth Science, 8, 740,
- 1112 10.3389/feart.2020.603441 2021.
- 1113 Rust, A. C., and Manga, M.: Bubble shapes and Orientations in low Re simple shear flow, Journal of 1114 Colloid and Interface Science, 249, 476-480, 10.1006/jcis.2002.8292, 2002.
- Rust, A. C., Manga, M., and Cashman, K. V.: Determining flow type, shear rate and shear stress in
 magmas from bubble shapes and orientations, Journal of Volcanology and Geothermal Research, 122,
 1117 112, 2003.
- Rust, A. C., and Cashman, K. V.: Permeability of vesicular silicic magma: inertial and hysteresis effects,
 Earth and Planetary Science Letters, 228, 93-107, 2004.
- Rust, A. C., and Cashman, K. V.: Permeability controls on expansion and size distributions of pyroclasts,
 Journal of Geophysical Research-Solid Earth, 116, 17, B11202
- 1122 10.1029/2011jb008494, 2011.
- 1123 Rutter, E. H.: On the nomenclature of mode of failure transitions in rocks, Tectonophysics, 122, 381-1124 387, 10.1016/0040-1951(86)90153-8, 1986.
- 1125 Sahagian, D.: Volcanology Magma fragmentation in eruptions, Nature, 402, 589-+, 1999.
- 1126 Sahetapy-Engel, S. T., and Harris, A. J. L.: Thermal structure and heat loss at the summit crater of an 1127 active lava dome, Bulletin of Volcanology, 71, 15-28, 10.1007/s00445-008-0204-3, 2009.
- 1128 Sato, H., Suto, S., Ui, T., Fujii, T., Yamamoto, T., Takarada, S., and Sakaguchi, K.: Flowage of the 1991
- Unzen lava; discussions to Goto et al., 2020 'Rigid migration of Unzen lava rather than flow' J. Volcanol.
 Geotherm. Res., 110, 107073, Journal of Volcanology and Geothermal Research, 107343,
 https://doi.org/10.1016/j.jvolgeores.2021.107343, 2021.
- Saubin, E., Kennedy, B., Tuffen, H., Villeneuve, M. C., Davidson, J., and Burchardt, S.: Comparative field 1132 1133 study of shallow rhyolite intrusions in Iceland: Emplacement mechanisms and impact on country 1134 rocks, Journal of Volcanology 388, and Geothermal Research. 106691. https://doi.org/10.1016/j.jvolgeores.2019.106691, 2019. 1135
- 1136 Schaefer, L. N., Kennedy, B. M., Kendrick, J. E., Lavallée, Y., and Miwa, T.: Laboratory Measurements
- of Damage Evolution in Dynamic Volcanic Environments: From Slow to Rapid Strain Events, 54th U.S.
 Rock Mechanics/Geomechanics Symposium, 2020,
- 1139 Scheu, B., Spieler, O., and Dingwell, D. B.: Dynamics of explosive volcanism at Unzen volcano: an
- 1140 experimental contribution, Bulletin of Volcanology, 69, 175-187, 2006.





- 1141Scheu, B., Kueppers, U., Mueller, S., Spieler, O., and Dingwell, D. B.: Experimental volcanology on1142eruptive products of Unzen, Journal of Volcanology and Geothermal Research, 175, 110-119,
- 1143 10.1016/j.jvolgeores.2008.03.023, 2007.
- Shields, J. K., Mader, H. M., Pistone, M., Caricchi, L., Floess, D., and Putlitz, B.: Strain-induced
- outgassing of three-phase magmas during simple shear, Journal of Geophysical Research-Solid Earth,
 1146 119, 6936-6957, 10.1002/2014jb011111, 2014.
- 1147 Smith, J. V., Miyake, Y., and Oikawa, T.: Interpretation of porosity in dacite lava domes as ductile-
- brittle failure textures, Journal of Volcanology and Geothermal Research, 112, 25-35, 10.1016/s03770273(01)00232-3, 2001.
- Smith, J. V.: Structural analysis of flow-related textures in lavas, Earth-Science Reviews, 57, 279-297,
 Pii s0012-8252(01)00081-2
- 1152 10.1016/s0012-8252(01)00081-2, 2002.
- Sparks, R. S. J.: Causes and consequences of pressurisation in lava dome eruptions, Earth and Planetary
 Science Letters, 150, 177-189, 1997.
- Sparks, R. S. J., Murphy, M. D., Lejeune, A. M., Watts, R. B., Barclay, J., and Young, S. R.: Control on the
 emplacement of the andesite lava dome of the Soufriere Hills volcano, Montserrat by degassinginduced crystallization, Terra Nova, 12, 14-20, 2000.
- Sparks, R. S. J.: Dynamics of magma degassing, in: Volcanic Degassing, edited by: Oppenheimer, C.,
 Pyle, D. M., and Barclay, J., Geological Society Special Publication, 5-22, 2003.
- Stasiuk, M. V., Barclay, J., Carroll, M. R., Jaupart, C., Ratte, J. C., Sparks, R. S. J., and Tait, S. R.: Degassing
 during magma ascent in the Mule Creek vent (USA), Bulletin of Volcanology, 58, 117-130, 1996.
- Stix, J., Layne, G. D., and Williams, S. N.: Mechanisms of degassing at Nevado del Ruiz volcano,
 Colombia, Journal of the Geological Society, 160, 507-521, 2003.
- 1164 Tait, S., Jaupart, C., and Vergniolle, S.: Pressure, gas content and eruption periodicity of a shallow,
- 1165
 crystallizing magma chamber, Earth and Planetary Science Letters, 92, 107-123, 10.1016/0012

 1166
 821x(89)90025-3, 1989.
- 1167Thomas, M. E., and Neuberg, J.: What makes a volcano tick--A first explanation of deep multiple1168seismic sources in ascending magma, Geology, 40, 351-354, 10.113/G32868.1, 2012.
- 1169 Tuffen, H., Dingwell, D. B., and Pinkerton, H.: Repeated fracture and healing of silicic magma generate1170 flow banding and earthquakes?, Geology, 31, 1089-1092, 2003.
- 1171 Tuffen, H., and Dingwell, D. B.: Fault textures in volcanic conduits: evidence for seismic trigger 1172 mechanisms during silicic eruptions, Bulletin of Volcanology, 67, 370-387, 2005.
- 1173 Umakoshi, K., Takamura, N., Shinzato, N., Uchida, K., Matsuwo, N., and Shimizu, H.: Seismicity
 1174 associated with the 1991-1995 dome growth at Unzen Volcano, Japan, Journal of Volcanology and
 1175 Geothermal Research, 175, 91-99, 10.1016/j.jvolgeores.2008.03.030, 2008.
- 1176 Varley, N. R., and Taran, Y.: Degassing processes of popocatepetl and Volcan de Colima, Mexico, in:
- 1177 Volcanic Degassing, edited by: Oppenheimer, C. P. D. M. B. J., Geological Society Special Publication,1178 263-280, 2003.
- 1179 Vasseur, J., Wadsworth, F. B., Lavallée, Y., Hess, K.-U., and Dingwell, D. B.: Volcanic sintering:
 1180 Timescales of viscous densification and strength recovery, Geophysical Research Letters, 40, 56581181 5664, 10.1002/2013gl058105, 2013.
- Venezky, D. Y., and Rutherford, M. J.: Petrology and Fe-Ti oxide reequilibration of the 1991 Mount
 Unzen mixed magma, Journal of Volcanology and Geothermal Research, 89, 213-230, 10.1016/s0377-
- 1184 0273(98)00133-4, 1999.
- Wadsworth, F. B., Vasseur, J., von Aulock, F. W., Hess, K.-U., Scheu, B., Lavallée, Y., and Dingwell, D.
 B.: Nonisothermal viscous sintering of volcanic ash, Journal of Geophysical Research-Solid Earth, 119,
 8792-8804, 10.1002/2014jb011453, 2014.
- 1188 Wadsworth, F. B., Vasseur, J., Scheu, B., Kendrick, J. E., Lavallee, Y., and Dingwell, D. B.: Universal
- scaling of fluid permeability during volcanic welding and sediment diagenesis, Geology, 44, 219-222,
- 1190 10.1130/g37559.1, 2016a.





Wadsworth, F. B., Vasseur, J., Scheu, B., Kendrick, J. E., Lavallée, Y., and Dingwell, D. B.: Universal
scaling of fluid permeability during volcanic welding and sediment diagenesis, Geology, 44, 219-222,
10.1130/g37559.1, 2016b.

Wadsworth, F. B., Vasseur, J., Llewellin, E. W., Dobson, K. J., Colombier, M., von Aulock, F. W., Fife, J.
L., Wiesmaier, S., K.-U., H., Scheu, B., Lavallée, Y., and Dingwell, D. B.: Topological inversions in

- 1196 coalescing granular media control fluid-flow regimes, PHYSICAL REVIEW E, 96, 033113, 2017.
- 1197 Wadsworth, F. B., Witcher, T., Vossen, C. E. J., Hess, K.-U., Unwin, H. E., Scheu, B., Castro, J. M., and
- 1198 Dingwell, D. B.: Combined effusive-explosive silicic volcanism straddles the multiphase viscous-to-
- brittle transition, Nature Communications, 9, 10.1038/s41467-018-07187-w, 2018.
- Wadsworth, F. B., Witcher, T., Vasseur, J., Dingwell, D. B., and Scheu, B.: When Does Magma Break?,
 in: Volcanic Unrest: From Science to Society, edited by: Gottsmann, J., Neuberg, J., and Scheu, B.,
- 1202 Advances in Volcanology, 171-184, 2019.
- 1203 Wadsworth, F. B., Vasseur, J., Llewellin, E. W., Brown, R. J., Tuffen, H., Gardner, J. E., Kendrick, J. E.,
- Lavallée, Y., Dobson, K. J., Heap, M. J., Dingwell, D. B., Hess, K.-U., Schauroth, J., von Aulock, F. W.,
 Kushnir, A. R. L., and Marone, F.: A model for permeability evolution during volcanic welding, Journal
 of Volcanology and Geothermal Research. 409. 107118.
- 1206 of Volcanology and Geothermal Research, 40 1207 https://doi.org/10.1016/j.jvolgeores.2020.107118, 2021.
- Wallace, P. A., Kendrick, J. E., Ashworth, J. D., Miwa, T., Coats, R., De Angelis, S. H., Mariani, E., Utley,
 J. E. P., Biggin, A., Kendrick, R., Nakada, S., Matsushima, T., and Lavallée, Y.: Petrological architecture
 of a magmatic shear zone: A multidisciplinary investigation of strain localisation during magma ascent
- 1211 at Unzen Volcano, Japan, Journal of Petrology, 60, 791-826, 10.1093/petrology/egz016, 2019.
- Watanabe, T., Shimizu, Y., Noguchi, S., and Nakada, S.: Permeability measurements on rock samples
 from Unzen scientific drilling project drill hole 4 (USDP-4), Journal of Volcanology and Geothermal
 Research, 175, 82-90, 10.1016/j.jvolgeores.2008.03.021, 2008.
- 1215 Westrich, H. R., and Eichelberger, J. C.: Gas transport and bubble collapse in rhyolitic magma an 1216 experimental approach, Bulletin of Volcanology, 56, 447-458, 10.1007/bf00302826, 1994.
- 1217 Woods, A. W., and Koyaguchi, T.: Transitions between explosive and effusive eruptions of silicic 1218 magmas, Nature, 370, 641-644, 1994.
- 1219 Wright, H. M. N., Roberts, J. J., and Cashman, K. V.: Permeability of anisotropic tube pumice: Model 1220 calculations and measurements, Geophysical Research Letters, 33, L17316
- 1221 10.1029/2006gl027224, 2006.
- Wright, H. M. N., and Weinberg, R. F.: Strain localization in vesicular magma: Implications for rheology
 and fragmentation, Geology, 37, 1023-1026, 10.1130/g30199a.1, 2009.
- Yamasato, H.: Nature of infrasonic pulse accompanying low frequency earthquake at Unzen volcano,
 Japan, Bulletin of the volcanological society of Japan, 43, 1-13, 10.18940/kazan.43.1_1, 1998.
- Yamashina, K., Matsushima, T., and Ohmi, S.: Volcanic deformation at Unzen, Japan, visualized by a
 time-differential stereoscopy, Journal of Volcanology and Geothermal Research, 89, 73-80,
 10.1016/s0377-0273(98)00124-3, 1999.
- Yilmaz, T. I., Wadsworth, F. B., Gilg, H. A., Hess, K. U., Kendrick, J. E., Wallace, P. A., Lavallée, Y., Utley,
 J. E. P., Vasseur, J., Nakada, S., and Dingwell, D. B.: Rapid alteration of fractured volcanic conduits
 beneath Mt Unzen, Bulletin of Volcanology, 83, 34, 2021.
- 1232 Yoshimura, S., and Nakamura, M.: Fracture healing in a magma: An experimental approach and
- implications for volcanic seismicity and degassing, Journal of Geophysical Research-Solid Earth, 115,
 10.1029/2009jb000834, 2010.
- 1235 Zhang, Y. X.: H2O in rhyolitic glasses and melts: Measurement, speciation, solubility, and diffusion,
 1236 Reviews of Geophysics, 37, 493-516, 1999.
- 1237

1238





1239 Figure Caption

1240 Figure 1. a) © Google Earth image showing the location of Unzen volcano on the island of Kyushu, 1241 Japan. b) Photograph of Unzen volcano, looking northwest, viewed from near Onokoba in the suburbs 1242 of Shimabara city. c)Photo of the relict 1994-95 spine at Unzen volcano (looking westward), showing 1243 (I) the central shear zone (i.e., the cavitation structures detailed in Smith *et al.*, 2001, further expanded in the inset); (II) the marginal shear zone, bordered by a fault (dark orange-brown colour), and (III) a 1244 1245 large block of sintered breccia of earlier domes, which has become welded to the fault material and extruded with the spine. Adapted from Hornby et al. (2015). d) Photograph of a fragment of the spine 1246 1247 showing the primary internal structure of the shear zone, bordered by a set of closely spaced, inclined fractures to the left and indurated breccia to the right. 1248

1249

1250 Figure 2. Location of the lava spine blocks and characteristics of the marginal shear zone. a) An aerial 1251 view of Unzen lava dome summit showing the remnants of the 1994-95 lava spine, including the main spine, the central shear zone (CSZ) block and the marginal shear zone (MSZ) block; image taken from 1252 1253 © Google Earth. b) Photograph of the main spine inclined towards the east. c) 3D construction of the 1254 marginal shear zone block (created using the photogrammetry 3DF Zephyr by 3Dflow). The outcrop 1255 is annotated to show the location of samples (A-H) as well as the 4 main regions (gouge as well as 1256 high-, moderate- and low-shear zones) and key features, including the fault contact (red dashed 1257 curve), shear zone transitions (yellow dashed curves), extension of tensile fracture (C; green lines) and Riedel fractures (blue curves). The inset shows detail of the fault plane, dividing the gouge and 1258 high-shear zone. Directional arrows X, Y and Z show the orientation of sample coring relative to the 1259 1260 shear plane. d) View of the MSZ block parallel to the shear plane and perpendicular to the shear 1261 plane. Insets show surface textures across the shear zone.

1262

1263 Figure 3. Composite figure of the microtextural characteristics across the marginal shear zone 1264 consisting of photograph of fresh surface textures, plane polarised light (PPL) photomicrographs, 1265 ultraviolet (UV) light photomicrographs and backscattered electron (BSE) images of the groundmass. 1266 Images of the fresh surface were taken following cutting the sample perpendicular to shear. 1267 Phenocryst observed include plagioclase (P), amphibole (A), biotite (B) and quartz (Q). Green boxes 1268 on PPL photomicrographs show the location of the UV light images, which highlight the pore 1269 structures across the MSZ. On UV light images, two white arrows pointing away from each other 1270 show the location of fractures within the groundmass (samples G and H), single arrows point to large pores adjacent to large phenocryst (samples G and H), and two arrows pointing towards each other 1271 1272 show compaction bands (their spacing represents the width of each band; samples B and C).

1273

1274	Figure 4. Tomographic reconstructions of four samples across the shear zones: a-b) A, c-d) C, e-f) E,
1275	g-h) H.; The upper row shows density-based images of tomographic reconstructions, whereas the
1276	lower row highlights the porous network in blue and the solid fraction is transparent. The
1277	reconstruction shows that the porous fraction becomes increasingly localised towards the fault plane
1278	(i.e., from right to left).

1279

1280 Figure 5. a) Porosity and permeability (parallel and perpendicular to shear plane) profile across the

1281 shear zone, showing the compactant (ductile) nature of the high shear zone, overprint by localised,





1282 dilational (brittle) fractures. Measurements on the gouge sample are plotted at a distance of 0 m. b) 1283 Porosity reduction as a function of effective pressure, derived from the volume of water expelled 1284 during loading in effective pressure of samples cored parallel to shear. Note that the initial porosity 1285 value (at Peff \simeq 5 MPa) is that of the sample initial porosity (before loading); the exact quantity of 1286 volume expelled between 0.1 and 5MPa cannot be accurately determined due to the method used, 1287 hence we simply show the porosity reduction from this point onward. 1288 1289 Figure 6. Permeability of the marginal shear zone as a function of effective pressure and direction to shear: measurements conducted a) parallel and b) perpendicular to the shear plane. The data shows a 1290 1291 reduction in permeability with effective pressure; yet the permeability profile across the shear zone remains, irrespective of the pressure conditions tested. The data shows contrasting permeabilities as a 1292 1293 function of direction, which create c) permeability anisotropy, cast here as the ratio between the 1294 permeability parallel and perpendicular to the shar plane. The anisotropy is most pronounced in the high shear zone and generally increases as samples were loaded to higher effective pressure due to 1295 1296 fracture closure. Note that the x-axis was truncated and the scale was expanded for the near-fault high-shear zone for which we conducted more measurements due to the structurally complex nature 1297 1298 of this area of the spine. Measurements on the gouge sample are plotted at a distance of 0 m. 1299 Figure 7. a) Photograph showing measurement locations for the field-based permeability 1300 1301 measurements, for the upper (orange) and lower (green) transects. b) Permeability data for the upper 1302 (orange) and lower (green) transects, plotted against distance. The data shows a drastic increase in 1303 permeability of ~3 orders of magnitude. 1304

Figure 8. Permeability-porosity relationship for Unzen dome lavas and similar effusive lavas. Blue
and red circles represent data from this study, made parallel and perpendicular to the plane of shear,
respectively. Grey circles show porosity data for Unzen from Mueller *et al.* (2005) and Kendrick *et al.*(2021), and open circles show permeability measurement on USDP drill cores from Watanabe *et al.*(2008). Other symbols show data for effusive products at similar dome eruptions.

1310

Figure 9. a) Conceptual model showing rheological shifts and evolution of permeability (seen as fluid 1311 1312 flow vectors) during pulsatory magma ascent and stick-slip faulting. The sketches illustrate the 1313 evolution of the extent of active shear zones (in orange), inactive areas (dark reds), active faults (blue) 1314 and inactive faults (grey), during magma discharge fluctuations. The dominant rheology in each area 1315 is numbered (1-4) and is linked to the deformation mechanism map for magma (shown in b). The 1316 sketches (a) show that shear narrows toward the eruption point as magmas is subjected to lower 1317 effective pressure (as shown in b). Compaction of the outer margin of the shear zones (dark red-1318 brown) would generate a zone of lower permeability (which may act as a local fluid flow barrier) As 1319 discharge rates increase, the width of the shear zone also narrows, and promote a switch to brittle 1320 failure at shallow depth (~500 m), causing the propagation of a primary fault plane and an adjacent 1321 Riedel fracture (which channels fluid flow; blue arrow). Upon discharge rate reduction, the shear zone 1322 would widen again and the fault would become inactive (stick phase), shifting the Riedel fracture to 1323 shallower depth. Upon renewed discharge rate increase, shear would narrow again, and faulting would 1324 generate another Riedel fracture. Thus, the distance between Riedel fractures may be used to resolve 1325 the magma ascent associated with inter-seismicity deformation (ISD). b) Deformation mechanism





- 1326 map for magma (adapted from Lavallée and Kendrick, 2020). Yield caps are displayed by blue and
- 1327 green lines representing brittle rupture (dilatant shear) and ductile cataclastic flow (compactant shear),
- 1328 respectively, and showing an increase in strength as a function of strain rate ($\dot{\varepsilon}$). At low strain rates or
- 1329 at high effective mean stress, magma flow viscously. The numbers refer to scenarios as displayed for
- 1330 different parts of the magmatic column in panel (a).





1331



Figure 1.





1332



Figure 2.







Figure 3.







Figure 4.







Figure 5.







Figure 6.











1338





Figure 8.





1339 a





Figure 9.