Geodynamics 101 Revision - Review Boris Kaus

This is the authors' response to the review of Boris Kaus. Our responses are indicated in green, with the original review presented in black.

This is a lengthy, but well written and useful manuscript that highlights many aspects of geodynamic modelling. I have a number of minor suggestions to further strengthen the manuscript in the attached pdf.

Review of van Zelst et al. 101 Geodynamic modelling: How to design, carry out, and interpret numerical studies
Boris Kaus, Mainz 1.6.2021

This is a well-written paper that is likely to be very useful for people that are new to geodynamic modelling. Putting all of that together in one place is challenging and I congratulate the authors for doing a good job with it. Unsurprisingly it is rather long and feels more like a book (which partly explains why it took me so long to get the review back to you - apologies!). This certainly deserved to be published in SE. Yet, I do have a number of comments which I believe would be good to address first. This can mostly be done with some rewriting, so I don’t think it should take you a lot of time (yet, it may help to point new geodynamicists in the right direction).

1) Modelling tools
You discuss quite a bit about the open source modelling packages like ASPECT, Underworld or LaMEM. Yet, missing from this discussion are alternative approaches that are based on smaller (sometimes one-page) scripts. Those are not full-blown modelling packages but rather simpler scripts that solve a particular problem very well (and fast). The classical example of that in geodynamics is the MILAMIN code (http://milamin.org) which remains one of the fastest codes to solve the incompressible Stokes equations on 2D unstructured meshes, and may be very helpful for those interested to solve for example, problems with viscous inclusions or interacting crystals. Other examples are the M2Di scripts which are concise matlab and julia codes to solve viscoelastic problems, including regularized plasticity, available from https://bitbucket.org/laess/m2di/src/master/. The most recent development in this direction is the ParallelStencil julia package (https://github.com/omlins/ParallelStencil.jl#stencilcomputations-with-math-close-notation) which comes with many geodynamic examples.

Geodynamicists that are interested to work on technical developments as well, may find such scripts much easier to understand than the big software packages that can do it all. In fact, with the ParallelStencil julia package it is possible to write a very compact code that scales to very large parallel GPU-based supercomputers with almost no effort (provided you use a pseudotransient iterative solver approach). Other efforts (under
development) allow calling PETSc and its staggered grid interface from Julia. In my opinion such approaches may become increasingly important in the future as it allows PhD students to go from writing an experimental solver to a fully blown (parallel) production code in a rather straightforward manner. This will help to address new multiphysics problems, such as the coupling between reactions & deformation. Given the informative nature of your current paper it would be good if you can discuss these topics as well in your manuscript (and give some of the links above).

We added a paragraph at the end of section 3.8 (l 801-814) about these recent approaches in building multi-physics applications. In this paragraph, we also addressed the comment mentioned later on doing multiphysics beyond the coupling of codes through boundary conditions.

2) Parameter sensitivity/controlling parameters
Typical forward models used in many geodynamic applications indeed have a large amount of parameters (as you discuss around lines 116 and lines 845). Usually, such sensitivity studies are done ‘by hand’ by modifying input parameters, making a model run and checking the difference with respect to a reference model. Yet, part of this can be done automatically by computing scaling exponents which directly show which of the parameters control the velocity at a certain point (as discussed in Reuber et al. 2018 Tectonophysics and used in Reuber et al. 2018, Front. Earth Science). In case adjoint methods are used to compute the gradients this is even computationally extremely efficient, and gives you the sensitivity to all model parameters at the same time. This would go a long way in determining which of those are of first order importance and which are not. This method only gives the sensitivity of the model results for a particular timestep/geometry but for many of the cases we looked at so far this sensitivity did not change drastically during a model simulation making this a quite powerful techniques (provided adjoints are available).
I suppose that the reason that this is not yet more widely applied is that it is not yet implemented in many of the codes currently in use, but I can well imagine that this may change soon. It would be good to highlight this as it is a very useful and automatic way to map and reduce the model parameter space (section 5.2.2) and will help to reduce the number of required simulations and thus the CO2 emissions of a study (line 1455).

This is indeed an exciting direction for future modelling. To include this into the manuscript, we added a sentence to this approach in the introduction (section 1.1):

“An alternative approach is to incorporate automatic parameter scaling routines or use adjoint methods to test parameter sensitivities in models (Reuber et al., 2018a,c), which could considerably reduce the amount of models required.”

To emphasise how adjoint methods will make scanning the parameter space more efficient - particularly in the case of general modelling studies where large parameter sweeps are common, we added a further explanation of adjoint methods at the end of paragraph 5.2.2:
“The mapping of a parameter space is often done through manual variation of a single model parameter and comparison of the resulting model predictions. However, recent developments allow for scaling laws between the model solution and the model parameters to be computed automatically through adjoint methods. Besides solving inverse problems (e.g. Ismail-Zadeh et al., 2003; Ghelichkhan and Bunge, 2016; Colli et al., 2018), adjoint methods can efficiently compute the scaling exponent for all model parameters with one linear solve (for a specific model time step) (Reuber et al., 2018c). These scaling exponents (that are based on the derivative of the solution parameter to model parameter) indicate which parameters control the model solution and which have a lesser effect (e.g. Reuber et al., 2018a; Crawford et al., 2018). Knowledge of the relative importance of each model parameter can help decrease the parameter space that is to be investigated (see also Section 7.3).”

3) Section 2.1.1: Mass conservation
Later in the manuscript (and around line 179/180 & lines 203-205) you discuss bulk compressibility, poison ratio etc., but you don’t show how that should be added to the mass conservation equation in equations (same in line 452). It would clarify matters if you can add this.

We incorporated this comment in various places in the manuscript:

Section 2.2.5: We now include an explicit equation for the density in our section on the equation of state (the simplest case of \( \rho = \rho_0 (1 - \alpha \Delta T) \)). We realise that there are many more complex equations of state, but there are many different ways for how, for example, the compressibility can be included in the equation of state, and many of them can not even be expressed as an equation anymore, for example a look-up table (we also added a sentence to highlight such efforts). In addition, choosing a different equation of state does not change the equations (2)-(7) beyond them being compressible or incompressible or using a specific approximation, which is something we already discuss. The only part that changes is the values of \( \rho \) (and how it depends on the solution variables). Because of this, we have decided to not add any other specific examples for equations of state.

Equation 2: We now also include an explicit reference to the equation of state section from the section about mass conservation.

After eq 4.: We also added the definition of Poisson’s ratio - defining it in terms of Lame’s first parameter and shear modulus for a homogeneous isotropic linear elastic material.

4) Solution methods
You discuss different solution methods in section 3.4, but what is completely missing is a discussion about multigrid preconditioners. Users of any 3D geodynamic code will run into having to use multigrid at some point or the other and will wonder why it sometimes does not converge (and sometimes does). As people new to geodynamic modelling are the target audience of this paper, it would be good if you can add a paragraph to
discuss this (and why, for many lithosphere dynamics problems, it is important to have a coarse grid that still “feels” the main viscosity structure of the model and is thus not too coarse).

What I believe is also important to discuss are pseudotransient solvers (there is much recent work by Raess et al.) as they result in compact solvers that scale particularly well on GPU’s systems (see link to the ParallelStencil julia package for sample codes).

We clarified and added that the solver methods described can be augmented with multigrid for optimized performance, also in response to a comment from Reviewer #1 (section 3.4). We have also added the following information on pseudo-transient solvers (lines 711-714):

“Recently, iterative pseudo-transient solvers have been used to solve coupled sets of equations (Räss et al., 2019; Reuben et al. 2020). These methods introduce a physics-based transient-term (a time derivative) to a steady-state equation, in order to iterate towards the steady-state solution. The matrix-free, finite difference pseudo-transient schemes of Räss et al. (2019) are well-suited for GPU (Graphical Processing Unit) accelerated systems (Räss et al. 2020).”

5) Particle-in-cell methods
I agree that for lithosphere dynamics codes, the particle-in-cell method is the most popular one (and I don’t really understand why you say it is difficult to parallelize). Perhaps you can explain here why that as, which is in my opinion because it is the easiest method to take things like phase transitions, history variables like strain as well as large deformations in a simple manner. Many alternative approaches have been suggested over the years, and many of those are good if dealing with a limited amount of phases (e.g. bubbles interacting with crystals can be well approximated with level sets). Yet, somehow none of these other methods withstood the test of time, perhaps because they are not general enough.

We removed the sentence that PIC methods are difficult to parallelize and significantly improved our section on tracking materials (section 3.7) (l 759-785; also in response to a comment from Reviewer #1).

6) Multiphysics
Multiphysics is indeed an important avenue of future research (perhaps even one of the most important ones in geodynamics). Yet achieving this by “code coupling” of different, unrelated, codes really only works if there is only a loose coupling between the physics (section 3.8). This is perhaps the case when coupling models of surface processes and lithosphere deformation. Other problems, on the other hand, require a much stronger coupling on the solver level which implies that new solvers should be developed. An example are the two phase flow equations that describe magma migration, which roughly consists in a Stokes-like and a Darcy-like problem. Getting the solution in an efficient manner cannot be done by taking a Stokes code and a separate Darcy code and coupling that using batch scripts. Instead, the coupled set of equations needs to be
solved in a tightly integrated manner. Our knowledge on the individual systems is still useful, as we can use multigrid preconditioners that work well for Stokes as part of this, for example. With this in mind, PETSc developed the multi-physics framework (see papers by Jed Brown) as well as the recently developed DMStag interface which allows you to add a Darcy-like code to an existing Stokes solver in a straightforward manner. I think it is important that you clarify this here, as there have been too many fruitless attempts in the past already to do loose coupling of different codes (which sounds intuitively easier but has very strong limitations for more tightly coupled problems).

That is a good point. We have added a paragraph at the end of the section, saying:

“However, some multi-physics problems are so closely intertwined that solving the coupled system of equations requires different numerical methods than solving the problems individually (for example, coupled magma/mantle dynamics). In these cases, coupling requires the development of new code that tightly integrates the different physical processes.”

This fits well with comment 1 on APIs, because using them could overcome these limitations on coupling codes that are incompatible at a fundamental level (i.e. defining a coupled system of equations).

7) Creep rheologies & modelling manuscripts
You spend quite some space describing how to write a modelling paper, which is certainly quite useful for people that are new to the community and/or to the paperwriting business. Yet, after having taught a class for over 10 years in Mainz in which students are supposed to reproduce published geodynamic modelling papers with a different code, I think the two most commonly made mistakes are not described here.

a) The first one is that in many papers, the model parameters are incorrectly listed in the tables, perhaps because the units of say dislocation creep laws are non-intuitive (the prefactor has units of the form of MPa^(-n) s^(n-1) where n is the powerlaw exponent; if n>1, transferring this to Pa^n s^-1 takes a bit of work). In addition, creep law experiments are usually given in terms of principal stresses, whereas geodynamic models need to have this in a tensor format, which involves correction factors (see the textbook of Gerya for a nice derivation). All these issues add up to make it quite difficult to fully reproduce the experiments unless the input scripts are attached. In most cases these are obvious typos, as employing the stated parameters results in wide spread drip tectonics, rather than stable subduction (as many generations of students in mainz have experienced), so it seems quite clear that it is actually correctly implemented in the code. Yet, whether the correction factor is taken into account or not is often less clear as it has a smaller effect. To me this points to a deeper underlying problem: there is currently no central database that collects all experimental creep rheology data in a format that is directly implementable in different geodynamic simulations. If we would have that it would eliminate quite a few of those mistakes.
b) A second issue is that often not all parameters are listed and some info is missing (nor just in terms of the material parameters employed but also with respect to the numerical convergence criteria etc.). I suppose that the reason for this is that there are often so many model parameters that it is quite easy to overlook some. A potential way around that would be to develop scripts that automatically generates tables with model parameters from input scripts which would eliminate another source of mistakes (which is that it seems that often some parameters are forgotten to be described in detail). Of course publishing the input script used to perform the simulations, as you point out, helps as well, but in that case it is still not fully guaranteed that the creep law used in one code are implemented in the same manner in a different code (while using the same correction factors).

To address both comments a,b, we added a paragraph in section 8.1 on how information should be given in the methods section (l 1416 - 1430).

I think it would be good if you can discuss these topics in the manuscript. To minimize the chance of mistakes in the future, developing standardized databases and automatized ways to create the parameter tables could be very helpful.

These are great ideas, but we think they should be a community effort and are therefore unfortunately beyond the scope of this manuscript. Summarised potential ideas for the community:

1. Central database that collects all experimental creep rheology data in a format that is directly implementable in geodynamic models.
2. Database with standardised model inputs/parameters tables: create automatized ways to print input parameters as tables ready for publications, to avoid editing mistakes.

Additional comments

line-wise (mostly typos, with some longer comments)

l. 55: One of the reference textbooks on analytical solutions is Turcotte and Schubert, which would be good to list here. Similarly, the textbook of Neil Ribe (theoretical mantle dynamics) has lots of useful info as well.

We added these additional references.

l. 170: highly viscous fluid

We corrected this typo.

l. 259: please define deviatoric stress in equations, and not just in words.

We include the equations in Section 2.2.1, and we now refer to that section.
I. 263: you did not explain what D/Dt means (or give equations)

We have added the definition (including the equation).

I. 269: The shear heating term should only involve the non-recoverable deviatoric strain rate components (that is, non-elastic) and not the full strain rate. I realize that you define strain rate only later in the paper and that for incompressible viscous rheologies this is equivalent. For compressible viscoelastic materials it is however not the same, so it seems important to point that out here.

We clarified the text.

I. 268: thin phase transitions? Not sure what you mean by that.

We changed it to “For phase transitions that occur over a narrow pressure/temperature range”

I. 349: It would be good if you can give the mathematical definition of how to go from total to deviatoric strain rate (which indeed simplifies to what you write here in the incompressible case).

We added an explanation after eq. 9.

I. 368: fH2O looks as if it could be several symbols; perhaps better to write as f_{H_2O}

Fixed equation 10.

I. 368; please indicate the units of all variables. In fact that is an issue throughout the manuscript and without specifying the units of parameters you employed, modelling results will not be reproducible and replicable…

Our equations are formulated independently of the units; i.e. the pressure could be given in Pa, or bar or any other unit, as long as the material parameters are adjusted accordingly. Consequently, we have decided to not indicate any units here (and throughout the manuscript). To emphasise the importance of units in normal studies of geodynamic modelling, we added sentences on how to properly include parameter tables (including units). Also see our response to comment 7b above and our response to the next comment.

I. 368/369: Can you add the units of all parameters?

We specifically do not include the units of the parameters in the manuscript, for multiple reasons:

- This would needlessly clutter the (already long) manuscript and reduce the readability of the manuscript in our opinion
- The units of some parameters (specifically constants even if we would stick to SI units) can differ in different formulations of the constitutive equations, making it an inherently incomplete and potentially confusing exercise to include the units.
- As mentioned above, as long as no values for specific parameters are given, the variables can be in any units as long as they are consistent with each other within the equation, as they are formulated independent of the units.

However, we agree that the consistent use of units is important for writing geodynamic modeling manuscripts and have added a sentence about this in Section 8.1 (also see our response to comment 7b).

I. 402: Plasticity can also be used for pressure-independent yield criteria. Examples are Griffith’s criteria (tensile failure) or implementing an ultimate yield stress in geodynamic models. It’s therefore better to remove ‘pressure-dependent’

We removed ‘pressure-dependent’.

I. 410: equation 14 is actually not identical to equation 13, unless cohesion C has a different meaning in eq. 14 compared to eq. 13. We actually had a discussion about this before in Solid Earth, so please have a look at https://se.copernicus.org/articles/11/1333/2020/ to see an illustration of the difference. The reason that cosine and sine terms appear in eq. 13 instead of tan(\(\mu_f\)) is that this is the yield stress function.

We agree that equation 13 and 14 indeed are not the same. We significantly rewrote section 2.2.4 to better introduce the various yield criteria.

I. 416/417: If lithostatic pressure is used to evaluate the yield criteria, the shear band orientation is always 45 degree and there is no difference between compression and extension. I realize that many large-scale convection codes use that assumption, and assuming lithostatic pressure is fine within the mantle. Yet, within the stronger lithosphere there can be quite strong deviations between lithostatic and dynamic P (up to factor 2 under compression). It’s probably good to point that out to the readers.

We added this further explanation:

“If lithostatic pressure is used in the yield criterion, shear bands are always orientated 45° under both extension and compression. This assumption is allowed for the mantle, where the total pressure is close to the lithostatic pressure. However, pressure can deviate strongly from the lithostatic pressure in the lithosphere, which can have major effects on the results.”


We added this reference.
I. 558: A big difference between FD and FE methods is that in FD, you solve the partial differential equation in a pointwise manner, whereas FE approximate the equations on average per element. It would be useful to add a comment on that here.

We added:

“This also controls how the partial differential equations are solved on the grid, with finite difference methods solving the equations pointwise and finite element methods averaging the equations per element.”

I. 562: Perhaps add a small remark on why finite elements should use a mixed formulation with higher order for velocity compared to pressure to get reliable results in the (near)-incompressible limit.

We added this as "Stable elements are typically characterised by m>n." and made a link with figure 10.

I. 584: 10^{24} Pas

We fixed this.

I. 740: the method of manufactured solutions also works for nonlinear problems which is perhaps good to mention here.

We changed the sentence to reflect this point.

I. 745: It is perhaps interesting to mention that community benchmarks is a process that typically takes several years…

We added a sentence to highlight this.

I. 963: geomio is spelled “geomIO”

We fixed this.

I. 1026: Implicit timestepping, in which advection becomes part of the nonlinear solution step, is an even better method to deal with the drunken sailor effect (was described in Popov & Sobolev 2008, even if only very briefly).

The Popov & Sobolev (2008) paper does not make a clear statement about implicit time stepping avoiding the drunken sailor effect, so we refrained from adding this reference and remedy.
Fig. 10: Maybe it is good to mention that the reason the drunken sailor effect occurs is that the typical density difference between rocks and air is much much larger than the typical density difference within the Earth. Moving the Earth’s surface by one meter therefore causes a much larger stress perturbation than moving the Moho by a meter and that is why the models have a tendency to become unstable.

We have clarified this in the main text (l. 1168-1172) as well as in the caption of Figure 10. We now say:

“The stark density contrasts (approx. 1.2~$\text{km}^{-3}$ versus 2830~$\text{km}^{-3}$) lead to much larger stress perturbations from topographic changes compared to similar topography variations at a typical density contrast inside the Earth (e.g. the density jump at the continental crust-mantle boundary is $\sim$280~$\text{km}^{-3}$; \cite{Martinec1994}).”

Fig. 10b: Even when you can ‘fix’ the wrong pressure field in this case by smoothing there are other, more heterogeneous, setups where such smoothing does not fully remove the artefacts. It is thus clearly better to employ LBB stable elements (like Q_2P_-1) throughout, The other undesired side-effect of unstable elements such as Q_1P_0 is that they require more iterations for a higher resolutions, if combined with iterative solvers. Stable elements fix that.

We added:

"Using an LBB stable element, like $Q_2 P_1$ or $Q_2 P_1$, avoids this problem (Donea & Huerta, 2003)."

to the figure caption, and a more extensive explanation to the main text. See our response to the reviewer’s comment about line 1047 (below).

Fig. 10c: I am not sure why you claim that the higher-resolution model is better here. It seems that both models are performed with non-regularized plasticity so both are actually numerically non-convergent. On one hand this manuscript argues about the importance of reproducible and numerically trustworthy results. On the other hand you show a key example in geodynamics where this is not the case. In my opinion there is no way around using a form of regularized plasticity (together with sufficient resolution such that the plasticity length scale is captured) and it would thus be better to use an example of that within the current manuscript. There are some recent papers by Duret et al. showing that this can be done, for example by using viscoplastic regularization, so it seems more appropriate to show examples of such computations in this figure.

In the main text (l 1154-1164), we state that common plasticity implementations are non-convergent with resolution, as illustrated in the figure. The figure however indeed suggests that the higher-resolution solution is better, but this was only meant to indicate that the fault
angle at higher resolution is closer to the theoretical angle. To avoid ambiguity, we have removed the check mark symbol from the higher resolution figure and updated the caption of Fig 10 to:

"The angle and thickness of the shear bands is dependent on the mesh resolution. Regularised plasticity implementations and sufficient resolution are required to achieve convergence with resolution (e.g. Duretz et al., 2020)."

We also added the Duretz et al. 2018 and 2019 and DeBorst et al. 2020 references to the main text.

I. 1047: Or better avoid unstable elements altogether and use stable ones (which also scale better on parallel computers using multigrid solvers).

We have included the following at line 1194 to point this out: "Stable elements, which fulfil the Ladyzhenskaya- Babuska-Brezzi compatibility condition (LBB or inf-sup condition), do not exhibit pressure artefacts (Donea and Huerta, 2003). Moreover, the required number of outer iterations does not increase significantly with mesh resolution compared to the Q1×P0 element (Thieulot and Bangerth, 2021)."

I. 1184: probabilistic

We fixed the typo.

I. 1311: In the geosciences it is common practice to acknowledge the work that reviewers put in going through the manuscript and making suggestions. Not everyone in computational geodynamics follows this (unwritten) rule, however, which I know doesn’t go well with many of my colleagues (including myself). So perhaps it is good to spell this out here.

We agree it is good practice to thank the reviewers for their work and have therefore added:

"Acknowledging the often substantial contributions of reviewers is a common courtesy."

**Online supplement**

Byerlee’s law: Byerlee’s law was originally derived for small-scale laboratory samples and is in itself quite amazing in that it shows that the maximum stress that rocks can withstand is nearly independent on the rock composition. What is even better is that it is nearly perfectly consistent with in-situ stress measurements in drill holes around the world (the classical reference would be Townend and Zoback, 2000). This shows that we can safely upscale Byerlee’s law from small samples to the upper crust. For other geodynamically relevant parameters, such as the effective viscosity of rocks such upscaling does not appear to be that easy. It would be good to point that our here (or in
We added the following sentence to the glossary entry for Byerlee’s law:

“Originally derived from small-scale laboratory samples, it has been shown that Byerlee's law can be safely upscaled to crustal conditions (Townend & Zoback 2000); something which is not straightforward for most geodynamically relevant parameters.”