

Contribution of gravity gliding in salt-bearing rift basins – A new experimental setup for simulating salt tectonics under the influence of sub-salt extension and tilting

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Abstract. Basin-scale salt flow and the evolution of salt structures in rift basin is mainly driven by sub- and supra-salt faulting and sedimentary loading. Crustal extension is often accompanied and followed by thermal subsidence leading to tilting of the graben flanks, which might induce an additional basinward directed driver for salt tectonics. We designed a new experimental analogue apparatus capable of integrating the processes of sub-salt graben extension and tilting of the flanks, such that the overlapping effects on the deformation of a viscous substratum and the brittle overburden can be simulated. The presented experimental study was performed to demonstrate the main functionality of the experimental procedure and setup demonstrating the main differences in structural evolution between conditions of pure extension, pure tilting and extension combined with tilting. Digital image correlation of top view stereoscopic images was applied to reveal the 3D displacement and strain patterns. Results of these experiments suggest that in salt basins affected by sub-salt extension and flank inclination, the salt flow and downward movement of overburden affects the entire flanks of the basin. Supra-salt extension occurring close to the graben centre is overprinted by the downward movement, i.e. amount of extension is reduced or extensional faults zones are shortened. At the basin margins, thin-skinned extensional faults developed, ~~which resemble fault zones observed on basin flanks offset from the central graben zone.~~ as a result of gravity gliding. A comparison with natural examples reveals that such fault zones can also be observed at the margins of many salt-bearing rift basins indicating that gravity gliding played a role in these basins.

15 1 Introduction

Salt layers in sedimentary basins play a key role in the structural and sedimentary evolution. As the bulk mechanical behaviour of rock salt on a geological time scale is that of a viscous fluid, it is able to flow in response to an internal pressure gradient and external shear forces (Vendeville and Jackson, 1992; Jackson and Hudec, 2017). Salt flow is mainly driven by two factors (Jackson and Hudec, 2017): (1) differential loading due to lateral differences in thickness and density of the suprasalt overburden (e.g. deltaic progradation, minibasin subsidence etc.) can trigger widespread squeezing of the salt and vertical collapse and lateral expansion of the overburden ('gravity spreading'; Schultz-Ela (2001)), and (2) any vertical displacement (e.g. tilting

of the salt layer) that can cause gravitational downward salt flow and, under some circumstances, gravitational collapse of overburden strata ('gravity gliding'; Schultz-Ela (2001)).

Driving processes of salt flow are well explained for salt-bearing passive margin basins, e.g. Lower Congo Basin, Western Mediterranean and the Santos Basin (Fort et al., 2004; Quirk et al., 2012; Jackson et al., 2015; Mianaekere and Adam, 2020), where both driving factors mostly act in the same seaward direction (Fig. 1a) (e.g. Brun and Fort, 2011; Peel, 2014). Scaled laboratory experiments focusing on passive margin salt tectonics reveal that seaward directed tilting of the basin floor and/or sediment progradation typically cause the formation of a downdip compressional and an updip extensional domain (Fig. 1a) (e.g. Jackson and Cramez, 1989; Mauduit et al., 1997; Brun and Mauduit, 2009; Brun and Fort, 2011; Ge et al., 2019b). The initiation of gravity gliding and spreading depends on the angle of the basin slope and of the top surface (e.g. Vendeville, 2005; Rowan et al., 2012), i.e. the gravitational potential has to be sufficient to overcome the compressional, frictional strength of the overburden strata at the foot of the slope (Rowan et al., 2004).

In salt-bearing rift basins (SBRB), these driving forces mostly act oppositely (Fig. 1b). On one side, the sedimentary regime in SBRB is commonly characterized by aggradation with maximum sediment thickness in the basin centre and less sedimentation or erosion at the basin margins. In many SBRB, such as the Glückstadt Graben (Northern Germany), the Central Graben (North Sea) and the Polish Basin (Central Poland) (Penge et al., 1999; Stewart, 2007; Krzywiec, 2004), salt is almost completely evacuated from the basin centre towards diapirs on the basin flanks, whereas the graben centre is filled up with several kilometres of syn-kinematic sediments (Krzywiec, 2004; Maystrenko et al., 2005; Warsitzka et al., 2017). This evacuation was likely driven by upward directed gravity spreading due to sedimentary loading (Hudec and Jackson, 2007; Warsitzka et al., 2018). Conversely, the graben flanks are often inclined towards a basin centre mainly due to thermal subsidence and flexural loading (e.g. Buchanan et al., 1996) causing basinward directed stresses. Downward gravity gliding is often inferred from the occurrence of a domain of thin-skinned extensional structures at the upslope basin margin (e.g. Penge et al., 1999; Geil, 1991; Stewart and Coward, 1995; Best, 1996; Withjack and Callaway, 2000; Thieme and Rockenbauch, 2001; Mohr et al., 2005; Jackson and Larsen, 2009; Tvedt et al., 2013; Vackiner et al., 2013; Labaume and Teixell, 2020). For instance, ~~in the northern Central Graben (central North Sea) and the Sole Pit Basin (southwestern North Sea)~~ in many rift basins in the central and southern North Sea, such as the Central Graben, the southern Viking Graben and the Sole Pit Basin (e.g. Hughes and Davison, 1993; Stewart and Coward, 1995; Jackson and Larsen, 2009), thin-skinned rollovers and separated blocks ('rafts') are observed above the tilted outer platforms of the rift basins (Hodgson et al., 1992; Buchanan et al., 1996; Penge et al., 1999; Stewart, 2014). It is suggested that these structures formed by a combination of gravitationally and tectonically driven extension ('rift-raft tectonics') (Stewart and Clark, 1999; Penge et al., 1999) and, in some cases, coincided with thin-skinned shortening in the basin centre (e.g. Hughes and Davison, 1993; Coward and Stewart, 1995; Stewart and Coward, 1995). ~~Thus~~ However, it is unclear ~~to-date~~ how gravity gliding can take place during sediment aggradation in the basin centre, because thick overburden sediments may act as resisting buttress against gravity gliding. ~~The oppositely acting processes of gravity spreading and gliding in SBRB provoke the question, which geological configurations have to be fulfilled to initiate gravity gliding in SBRB, i.e. which minimum topographic gradient and basin slope is required~~ The oppositely acting processes of gravity spreading and gliding in SBRB prompt the question of which geological configurations have to be fulfilled to initiate gravity gliding in the

case of SBRB, i.e. which minimum topographic gradient and basin slope is required? Furthermore, it is unclear, how gravity gliding affects basin-wide salt flow and the evolution of salt structures in SBRB and which driving process prevail at what period of the basin evolution.

60 Analogue models dedicated to the opposing interplay between gliding and spreading in SBRB show that salt flows downward if the depocentre above the downthrown block is underfilled and returns to upward directed flow as soon as syn-kinematic sediments are accumulated in the depocentre (e.g. Vendeville et al., 1995; Koyi et al., 1993; Ge et al., 1997; Dooley et al., 2005; Burliga et al., 2012; Warsitzka et al., 2015; Lymer et al., 2018; Roma et al., 2018; Dooley and Hudec, 2020). Basin-scale models reproduced thin-skinned extensional structures occurring on the flanks of the graben resulting from the decoupling effect of the salt layer (Nalpas and Brun, 1993; Koyi et al., 1993; Dooley et al., 2005; Ferrer et al., 2014; Moragas et al., 2017). Other experiments demonstrated how progradation of sedimentary wedges across extensional systems modify sedimentary patterns in the overburden by introducing an additional component of gravity spreading (Loncke et al., 2010; Rojo et al., 2020). ~~However, the influence of basin-wide tilting of the subsalt basement and, hence, effects of gravity gliding~~ However, the influence of basin-scale tilting of the graben flanks and, hence, effects of gravity gliding, on the evolution of supra-salt sedimentary structures and salt flow patterns have not been investigated yet.

70 The key motivation of this ~~project work~~ is to evaluate the influence of regional-scale gravity gliding on salt flow processes and the ~~initiation of minibasins in~~ tectono-sedimentary evolution of SBRB. We designed a new experimental apparatus for analogue modelling to integrate the simulation of crustal scale extension and tilting of the graben flanks. Here, the basic concept of the apparatus and results of preliminary experiments ~~is~~ are presented. Based on the techniques of digital image correlation, we compare displacement and strain patterns between different scenarios: (1) only extension, (2) only tilting of the flanks, (3) combined extension and tilting and (4) extension, tilting and syn-kinematic sedimentation. In order to demonstrate the effect of ~~the presence of~~ a viscous detachment, these scenarios were performed with and without a viscous substratum. The experimental results illustratively demonstrate the functionality of the apparatus and reveal first main differences between rift basins with tilted flanks and those surrounded by flat flanks.

80 2 Geological prototype and experimental setup

The experimental setup and the procedure are inspired by a generalized natural salt-bearing rift basin (Fig. 1c; Tab. A1 in Appendix A). Boundary conditions and basin configurations, such as ~~width of the basin, maximal amount of tilting, extension rate, layer thicknesses and timing of extension~~ age and thickness of the salt layer, timing of rifting, thickness of the pre-kinematic overburden, maximal amount of tilting and fault displacement as well as basin width, were derived from a survey of various SBRB worldwide (Tab. A1). In most basins, the initial phase of rifting occurred prior to deposition of the major salt layer and in some case continued during deposition. The early post-salt evolution (~ 10 to 20 Ma after salt deposition) is usually characterized by multiple phases of rifting each followed by post-rift thermal subsidence. Thermal relaxation and sedimentary loading usually cause subsidence in a wider region compared to that of the stretched extensional graben (Watts et al., 1982). We assume that these subsidence processes are the main drivers for the basinward tilting of the graben flanks. Such tilting can be as much

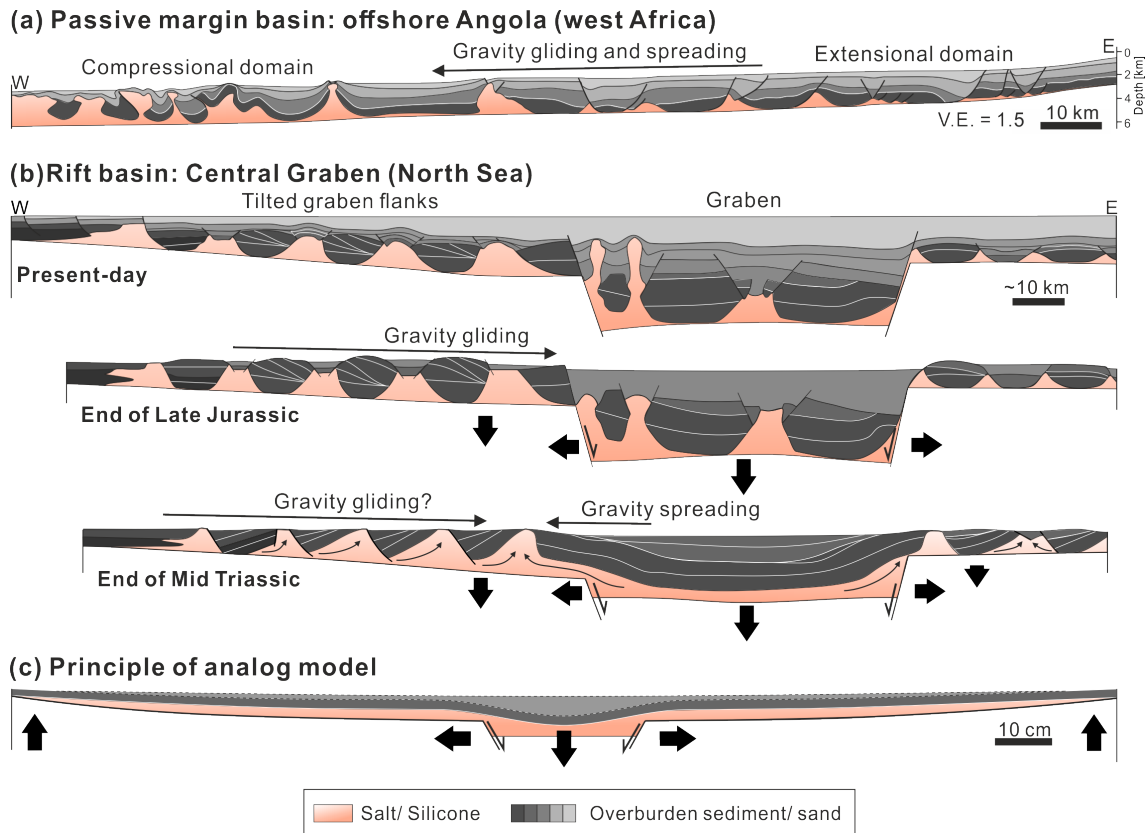


Figure 1. (a) Geological cross section of the Angolan passive margin, where basinward directed gravity gliding and spreading lead to the formation of an upslope extensional and a downslope compressional domain (Fort et al., 2004). (b) ~~Geological cross section of the Central Graben (North Sea) (van Winden et al., 2018). Diapirs and other salt-related structures develop due to sub- and supra-salt extension.~~ (c) ~~Restored cross section of (b) to the top Upper Triassic showing inclined flanks of the rift basin. Note that redistribution of the salt was included in the restoration procedure by van Winden et al. (2018). However, restored thickness of Triassic layers and estimates of the original salt layer thickness (Ten Veen et al., 2012) provide reasonable implications for the geometry of the sub-salt base.~~ (d) Schematic cross section showing structures typically observed in SBRB, such as thin-skinned extensional "rafts" and rollovers on the graben flanks and as well as diapirs and minibasins in the centre. (e) Concept of the experimental setup applied in the presented study. Graben extension is simulated by a vertically moving central block and two laterally moving flanks. Tilting of the flank (due to thermal subsidence) is modelled by an upward bending of the outer margins. Schematic evolution of the structure and the salt flow patterns in the Central Graben (North Sea) (modified from Stewart and Clark, 1999). It is suggested that gravity gliding above the tiled western flank of the graben coincided with upward salt flow from the graben centre towards the rift shoulders (Penge et al., 1999). (c) Conceptual sketch of the experimental setup applied in this study. Graben extension is simulated by a vertically moving central block and two laterally moving flanks. Tilting of the flanks (due to thermal subsidence) is modelled by upward bent plates. α – Angle of the salt base, V.E. - vertical exaggeration.

90 as 10° (4.5° in average; ~~over a length of up to 100 km (60 km in average;~~ Tab. A1). The supposed maximal thickness of the original salt layer in the basin centre ranges between ~~800 600~~ and 4000 m (~~1800 1600~~ m in average), but pinches out towards the basin margins. The pre-kinematic layer (~ 400 m in average; Tab. A1) is defined as the post-salt overburden deposited prior to the first post-salt extensional phase. In some basins, ~~the thickness of this layer is probably 0 (~ 450 m in average)~~ there was probably no pre-kinematic overburden, because syn-salt rifting continued after end of salt deposition. The lithology of the pre-kinematic overburden can be fine- or coarse-grained clastics or carbonates. Therefore, we assume an averaged lithological composition when relating physical properties between model and nature (see below).
95

3 Method

3.1 Experimental setup

The experimental setup combines two main approaches of previous analogue modelling studies: 1. the simulation of upper
100 crustal extension (e.g. Koyi et al., 1993; Dooley et al., 2005; Ferrer et al., 2014; Moragas et al., 2017; Roma et al., 2018) and 2. the effect of basin-scale tilting of the basin floor as it was applied to model salt tectonics at passive margins (e.g. Mauduit et al., 1997; Fort et al., 2004; Dooley et al., 2017; Ge et al., 2019b). ~~The experimental apparatus consists of a ~ 30 cm wide graben structure in the centre, which can be moved vertically~~ The apparatus consists of a ~ 30 cm wide, rigid graben block, which is bounded by two 60° faults and can be moved vertically by a step motor (Fig. 2a). The graben flanks are made of ~ 100 cm long
105 (x-axis) and 60 cm wide (y-axis) bendable plates (~~~ 100 cm~~), which are pushed outward by the down-going central block to simulate tectonic extension. The dimensions of the apparatus are chosen to simulate the wide range of variably sized rift basins with a maximal extent of ~ 220 km (Tab. A1) ~~Independent from~~ Besides the lateral movement, the outer edges of the flanks can be uplifted by independently controlled step motors so that the plates are bent and ~~inclined~~ tilted. Thermal subsidence in nature leads to widespread, saucer-shaped subsidence of the basin. However, the net effect of thermal subsidence is differential
110 vertical movement between the graben centre and the flanks, which is here mimicked by an uplift of the flanks. The flanking plates are made of 4 mm thick steel and vertically supported at the footwall tips of the basal faults. This ensures that neither downpull of the graben block nor loading of the analogue materials lead to flexural bending of the plates in the basin centre. ~~The depth (y-axis) of the plates is roughly 60 cm.~~ In the presented experiments, the silicone basin extends over an area of roughly 160×50 cm and is confined by sand walls (Fig. 2b). The width of the silicone basin (y-axis) is sufficient to reduce
115 boundary effects ~~at the confining lateral sand walls (Fig. 2b).~~ The silicone layer is thickest in the basin centre ($h_v = \sim 1.5$ cm) and pinches out towards the basin margins (Fig. 2c).

The experimental surface is recorded by 2 stereoscopic CMOS sensor cameras ('Imager M-lite 12M camera', 12-bit monochrome, 12 MPx resolution) in top view to analyse the evolution of the surface displacements and strains (Adam et al., 2005). Images are acquired at intervals of 300 s and processed by the StrainMaster[®] (LaVision GmbH) 3D digital image correlation (DIC)
120 software (Adam et al., 2005; LaVision, 2018). The 3D-DIC analysis is an optical, non-intrusive method to determine shapes, displacement and deformation in ~~high~~ sub-millimeter resolution.

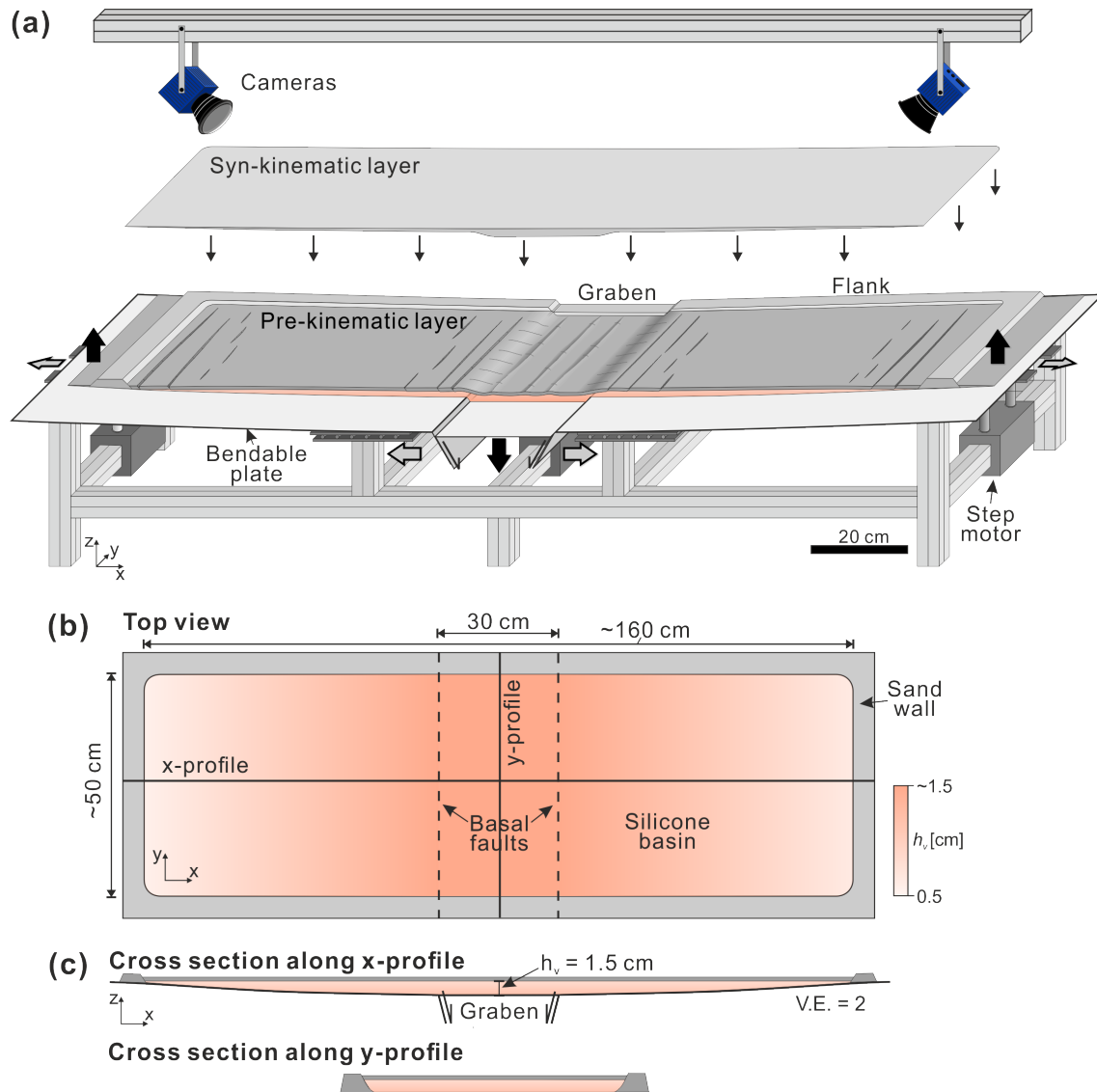


Figure 2. (a) Sketch of the experimental setup. The apparatus consists of a central, vertically moving graben structure bounded by two bendable metal plates. The latter can be moved laterally and their outer edges can be uplifted. All parts can be controlled separately. The experimental surface is monitored by 2 stereoscopic cameras from above and images are utilized for 3D digital image correlation (Adam et al., 2005). Black arrows denote active movement of the basal parts driven by motors (i.e. down pull of the central graben and uplift at the outer edges of the flanks), and grey arrows display passive movement (i.e. lateral push of the flanking plates). (b) Top view sketch of the initial silicone basin displaying the thickness distribution of the silicone layer. The basin is confined by sand walls, which are not shown in (a). Top view of the silicone basin, which is confined by sand walls. (c) Cross section views of the initial silicone layer, which is covered by a 3 mm thick pre-kinematic sand layer (not shown in (b)). V.E. - vertical exaggeration.

3.2 Experimental materials

The viscous behaviour of salt (e.g. Urai et al., 2008) is simulated by a silicone putty (Polydimethylsiloxane, PDMS; Rudolf et al. (2016)) ~~sold under the name Korasilon G 20 OH (Kurt Obermeier GmbH & Co. KG), which is a KORASILON M-type~~
125 ~~fluid sold by the company Kurt Obermeier GmbH & Co. KG.~~ The viscosity of the silicone was measured in the analogue laboratory of the Helmholtz Centre Potsdam GFZ German Research Centre for Geosciences using a cone-plate rheometer (RC20.1-CPS-P1, RheoTec). The dynamic viscosity is $2.2\text{--}2.8 \times 10^4$ Pas, which is slightly temperature-dependent (decrease of 1.4–1.9 % per 1 K increase), and the density is 970 kg m^{-3} at 25°C . It has a Newtonian behaviour at low strain rates ($<10^{-2}$) and a strain rate softening behaviour above relatively high strain rates ($>10^{-2} \text{ s}^{-1}$) (Rudolf et al., 2016).

130 Frictional-plastic behaviour of overburden sediments is modelled by granular mixtures of quartz sand (bulk density = 1620 kg m^{-3}) and silicate cenospheres (bulk density = 430 kg m^{-3}) ~~with various mixing ratios.~~ In order to adapt the density ratio between the ~~ductile~~ viscous silicone and the overburden sand to match the natural density ratio, various mixtures of both granulates were used (~~see Appendix C for details~~). Similar to upper crustal rocks, such granulates deform according to the Mohr–Coulomb failure criterion (Equ. 1; (e.g. Lohrmann et al., 2003; Panien et al., 2006)). Frictional properties (coefficient of
135 internal friction μ , cohesion C ; Tab. B1) and densities of the mixtures were determined in ring shear tests in the laboratory of the Helmholtz Centre Potsdam GFZ German Research Centre for Geosciences. Details of the measurements and the material properties can be found in Warsitzka et al. (2021).

3.3 Scaling

For the analogue model to be representative of the natural prototype, models have to be geometrically, dynamically, kinemati-
140 cally and rheologically similar to their natural prototype (Hubbert, 1937; Weijermars et al., 1993) ~~Ramberg, 1981; Weijermars and Schmeling, 1986; Davy and Cobbold, 1991; Bonini et al., 2012.~~ In the presented models, we simulated the deformation processes of tectonically driven extension and gravity driven gliding and spreading and identified the following characteristic parameters describing the boundary conditions of these processes (cf. Fig. 3): layer thickness (h_b , ~~h_a~~ h_x), densities (ρ_b , ~~ρ_a~~ ρ_v), cohesion and coefficient of internal friction of the brittle overburden (C , μ), viscosity of the ~~ductile~~ substratum (η), extension
145 rate (~~tt_e~~ v_e) and tilting rate (~~tt_t~~ v_t). Extension and tilting rates depend on the geometrical and kinematical scaling and determine the amount of extension and tilting during experimental time.

In order to relate these parameters between nature and model, we applied standard scaling procedures as in other studies about geological systems consisting of viscous substratum and brittle cover (Hubbert, 1937; Ramberg, 1981; Weijermars and Schmeling, 1986; Davy and Cobbold, 1991; Weijermars et al., 1993; Bonini et al., 2012). The detailed scaling procedure is
150 described in Appendix B.

The geometrical scaling ratio is set to 10^{-5} meaning that 1 cm in the model represents 1 km in nature (Tab. B1 in Appendix B). Principles of dynamical scaling require that driving and resisting forces acting during deformation have to be related properly (Weijermars and Schmeling, 1986; Pollard et al., 2005). Drivers for salt flow are produced by gravitational buoyancy forces due to loading of sediments on top of the salt layer, and by tectonically driven lateral pressure forces (Jackson and

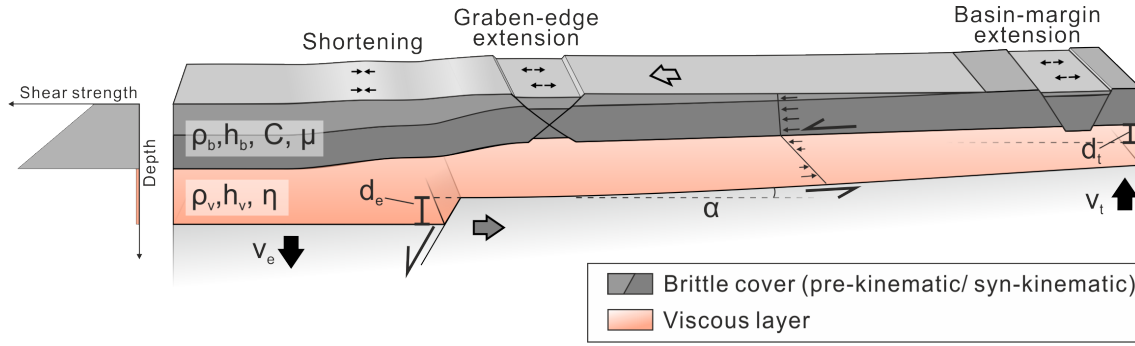


Figure 3. Physical parameters describing the deformation processes in the experiments. Lateral displacement due to basal extension and downward gliding of the overburden lead to shearing at the base and the top of viscous layer. Vertical displacement of the graben block and the flanks cause gravitational instability in the viscous layer. Syn-kinematic sedimentation, i.e. filling the subsided depocentres, imposes differential loading on top of the viscous layer. Governing parameters describing these processes are: α – tilt angle of the basin flank, d_e – vertical displacement of the graben centre, d_t – vertical displacement of the graben flanks, h_b – thickness of the brittle overburden, h_v – thickness of the ductile viscous substratum, η – dynamic viscosity of the ductile material, μ – Coefficient of internal friction, C – cohesion, v_e – (vertical) extension rate, v_t – uplift rate v_t – tilting rate

155 Hudec, 2017). Resisting forces result from the frictional strength of the brittle overburden (S_F) and viscous stresses in the substratum (S_V) (Jackson and Hudec, 2017). The shear strength of brittle sediments and experimental granular materials obeys the Mohr-Coulomb criterion (Byerlee, 1978; Dahlen, 1990; Jaeger et al., 2007), where the mean stress σ_m and the maximum shear stress τ_{max} are related by means of the angle of internal friction ϕ and the cohesion C :

$$\tau_{max} = \sigma_m \sin\phi + C \cos\phi \quad (1)$$

160 This is translated to an expression by the maximum and the minimum principal stresses σ_1 and σ_3 to define the frictional strength S_F (see Appendix B). The force F_F required for extensional failure ($\sigma_1 =$ lithostatic pressure) is then derived by integrating S_F over the thickness h_b (Schueller and Davy, 2008):

$$F_F = \frac{1}{2} \rho_b g h_b^2 (1 - \lambda) - \frac{1}{2} \rho_b g h_b^2 (1 - \lambda) \frac{1 - \sin\phi}{1 + \sin\phi} + 2Ch_b \sqrt{\frac{1 - \sin\phi}{1 + \sin\phi}} \quad (2)$$

165 The large-scale viscous behaviour of the salt layer is commonly characterized by the stress-strain rate relationship of rock salt (halite), which usually constitutes the main volumetric proportion of a mobilized evaporitic succession (Jackson and Hudec, 2017). For modelling purposes, the salt layer is here approximated as temperature-independent, mechanically uniform viscous material with a linear (Newtonian) stress-strain rate behaviour (Weijermars and Schmeling, 1986). This condition might be true for wet salt ($> \sim 0.05$ wt% water content) with relatively small grain sizes (< 0.5 mm) (Carter and Hansen, 1983; Schläder

et al., 2008; Urai et al., 2008). Under this assumption, the differential (viscous) stress is calculated by (Weijermars et al., 1993; 170 Bonini et al., 2012; Turcotte and Schubert, 2014):

$$S_V = \eta \dot{\gamma} \quad (3)$$

Here, η is the dynamic viscosity and $\dot{\gamma}$ the shear strain rate in x-direction. The viscous force is then given by integrating Eq. (3) over h_v (see Appendix B) and approximating $\dot{\gamma}$ by the ratio of the shear velocity at the base of the viscous layer v_x and the thickness of the viscous layer h_v :

$$175 \quad F_V = \eta \dot{\gamma} h_v = \eta v_x \quad (4)$$

For correct dynamical scaling, the ratio between brittle and viscous force has to be similar in model and nature (e.g. Sornette et al., 1993; Schueller and Davy, 2008). This ratio is here called 'Brittle-viscous-coupling' (BVC):

$$BVC = \frac{F_F}{F_V} \quad (5)$$

Accurate scaling is difficult to achieve, because natural values of viscosity ($10^{16} - 10^{19}$ Pas; Van Keken et al. (1993); 180 Mukherjee et al. (2010)) and strain rates ($10^{-16} - 10^{-14}$ s $^{-1}$; Jackson and Hudec (2017)) can be highly variable. We inserted a wide range of possible natural and experimental values (Tab. B1) in Eq. (4) and Eq. (2) to calculate the BVC, which shows that the BVC is relatively similar between model and nature (Fig. B1a).

Equation (2) reveals that the thickness of the overburden h_b mainly governs the coupling between brittle and viscous layer due to its quadratic input. Frictional forces are 1–2 orders of magnitude higher than viscous forces, at least at the beginning of 185 the experiment when strain in the viscous layer is distributed widely. During progressive strain localization in later stage of the experiment, higher strain rates and, hence, higher viscous forces arise. Therefore, the BVC can be a measure for strain localization (Sornette et al., 1993; Schueller and Davy, 2008) in future advanced experiments when strain evolution is quantified. The ratio between brittle and viscous material properties also sets the scaling for extension, tilting and sedimentation rates (see Appendix B) (Weijermars et al., 1993; Allen and Beaumont, 2012). Applying average values of viscosity, thickness and density 190 (Tab. B1), the velocity scaling ratio is $\sim 10^4$. Natural extension rate can vary between 0.1 and 10 mm a $^{-1}$ so that experimental extension rate should be in the range of ~ 0.1 to 10 mm h $^{-1}$. We choose an intermediate rate of 1 mm h $^{-1}$ for practical reasons so that a simulation duration of several hours is achieved.

The main drivers in the scenario modelled here are related to tectonic and buoyancy forces. The tectonic force is approximated by the integrated strength of the brittle and the viscous layer and, therefore, the sum of F_F (Eq. (2)) and F_V (Eq. (4)) 195 (e.g. Schueller and Davy, 2008). The buoyancy force F_B results from isostatic sinking of the overburden into the viscous

substratum and, therefore, from the weight of the overburden related to the weight of the equivalent volume of the viscous substratum. The buoyancy force can be calculated by (see Appendix B) (e.g. Schueller and Davy, 2008):

$$F_B = \frac{1}{2} g \rho_b h_b^2 \left(\frac{\rho_b}{\rho_v} - 1 \right) \quad (6)$$

200 F_B is mainly influenced by the density ratio between overburden (ρ_b) and salt (ρ_v). Salt density is constant with depth, whereas the density of other sediments increases with depth due to compaction (see Appendix C). During early stage of the basin history, the overburden density ρ_b is smaller than the salt density ρ_v so that F_B is negative. It changes to positive values when ρ_b exceeds ρ_v .

For scaling purposes, we adapted the so-called 'Argand number' (Ar) from analogue models of lithospheric-scale brittle-viscous systems (e.g. Schueller and Davy, 2008), which is the ratio between the buoyancy and the tectonic force:

$$205 \quad Ar = \frac{F_B}{F_F + F_V} \quad (7)$$

Figure B1b shows that the curves of Ar are similar between model and nature when inserting a range of potential values (Tab. B1) into Eq. (7). Ar changes from negative to positive values meaning the general behaviour of the influence of buoyancy force is adequately approximated in the models.

3.4 Experimental procedure

210 Since all parts of the apparatus can be controlled separately, different modes of basin evolution can be tested, ~~such as pure extension without tilting, pure uplift of the flanks without extension and combined motion of graben and flanks. In the experiment with combined modes (Exp. 3, 6 and 7), extension and tilting was applied simultaneously. This procedure such as pure basal extension, pure tilting of the flanks and simultaneous basal extension and tilting (Tab. 1).~~ The latter scenario refers to the geological condition that in most rift basins salt was deposited during or after an initial rifting phase (Tab. A1).
 215 Hence, it can be assumed that early post-salt history is characterized by thermal sag basin subsidence accompanied by additional post-salt rifting.

In a first series of experiments (~~E1, T1, ET1~~Eb, Tb, ETb), no silicone layer was included in order to present benchmark experiments for the displacement ~~and strain~~ patterns of the basal parts and to better illustrate the effects of the viscous layer in subsequent experiments. In experiments including a silicone layer (~~E2, T2, ET2, ETS~~Ev, Tv, ETv, ETSv), the silicone layer
 220 was covered by ~~a pre-kinematic an even, ~3 mm thick~~ layer of granular material (bulk density = $\sim 850 \text{ kg m}^{-3}$) prior to onset of deformation ('pre-kinematic layer'). In experiments in which basal extension was applied, the central graben block (Fig. 2) was pulled down with a rate of $v_e = 1 \text{ mm h}^{-1}$ to a max. vertical displacement of $d_e = 20 \text{ mm}$ (Tab. 1). This downward pull pushed the graben flanks outward up to max. horizontal extensional displacement of 12 mm per side. In experiments with tilting, both flanks were uplifted with a rate of $v_t = 1 \text{ mm h}^{-1}$ to a max. vertical displacement of $d_t = 20 \text{ mm}$ at the outer edges of the flanks
 225 (Tab. 1). This caused a parabolic bending of the flanking plates resulting in a max. tilting angle of $\alpha = 5^\circ$ at the margins of the

Table 1. Experiments and key parameters. In the experiment's name, E stands for extension, T for tilting, S for syn-kinematic sedimentation, b for (only) brittle and v for viscous. d_e – total amount of vertical displacement of the graben centre, d_t – total amount of vertical uplift of the graben flanks, h_b the thickness of the pre-kinematic sand layer, h_a – ~~thickness of the silicone~~ h_{vmax} – max. thickness of the silicone in the basin centre, T – ~~total duration of the experiment~~ T_{syn} – duration of the syn-deformation phase (extension and/or tilting), T_{post} – duration of the post-deformation phase, $u_e v_e$ – (vertical) extension rate, $u_t v_t$ – uplift rate of the flanks. See Fig. 3 for illustration of the parameters.

Experiment no.	Experiment name	h_b [mm]	h_v/h_{vmax} [mm]	$u_e v_e$ [mm h ⁻¹]	$u_t v_t$ [mm h ⁻¹]	d_e [mm]	d_t [mm]	T_{syn} [h]	T_{post} [h]	Comment
1	E1Eb	0.33	0	1	0	20	0	20	0	No silicone
2	E2Ev	0.33	1.515	1	0	20	0	20	12	
3	T1Tb	0.33	0	0	1	0	20	20	0	No silicone
4	T2Tv	0.33	1.515	0	1	0	20	20	3	
5	ET1ETb	0.33	0	1	1	20	20	20	0	No silicone
6	ET2ETv	0.33	1.515	1	1	20	20	20	12	
7	ETS1ETSv	0.33	1.515	1	1	20	20	22	13	Including syn-kinematic sedimentation

silicone basin and a tilting of $\sim 1.5^\circ$ on average across the entire flank. In experiments with a viscous layer, the syn-extensional or syn-tilting phase was followed by a phase in which we continued to monitor the surface deformation resulting from ongoing flow of the silicone (T_{post} ; Tab. 1). Note that T_{post} in exp. Tv is shorter than in the other experiments due to the lack of data in this experiment.

230 Except for experiment ETS1ETSv, no syn-kinematic sand accumulation sedimentation was applied to demonstrate the deformation patterns resulting only from the specific mode. Only ETS1 ETSv involves all relevant processes, which should be investigated with this apparatus (extension, tilting, syn-kinematic sedimentation). In intervals of 5 h, syn-kinematic sand is sieved on the model surface, whereas largest thickness is applied in the graben centre and decreasing thickness on the flanks. syn-kinematic sand was sieved on the model surface until the graben centre was completely filled up. Furthermore, 235 syn-kinematic sand with a uniform thickness of ~ 1 mm was added on the flanks. The total volume of the syn-kinematic layer was the same in each sedimentation interval, whereas its density was gradually increased according to the procedure described in Appendix C. This is supposed This sieving procedure is aimed to mimic aggradational sedimentation pattern as observed in most continental basins. Due to interruptions during sieving, the experimental duration of exp. ETSv is longer than in the other experiments (Tab. 1). After finishing experiment ETSv and soaking the sand layers with water, cross sections were cut 240 and photographed in order to examine internal, syn-kinematic structures.

4 Experimental results

In the following, we compare experiments, which were conducted without a viscous silicone layer (Eb, Tb, ETb) with those containing a viscous substratum (Ev, Tv, ETv). Then, we show differences in displacement and strain patterns between the

245 brittle-viscous experiment with only extension (Ev) and the experiment with extension and tilting (ETv). Finally, the experiment including syn-kinematic sedimentation simultaneous to extension and tilting (ETSv) is presented.

4.1 With vs. without viscous substratum

A first series of experiments (E1, T1, ET1) was conducted without a viscous silicone layer to illustrate the movement of the basal parts and to compare them to displacement patterns developed in experiments containing a viscous substratum (E2, T2, ET2).

250 4.1.1 Basal extension

When only ~~lateral~~basal extension was applied as boundary condition (~~experiment E1 in Eb (without viscous layer)~~), both flanks were pushed apart laterally by the downward moving central graben block (Fig. 4a and c). Except for straight fault zones above the basal normal faults, no further deformation structures could be observed in Eb. In experiment ~~E2 (including viscous layer), this outward directed movement was overprinted by deformation of the cover layer and the viscous substratum close to the central graben~~ Ev, basal extension caused deformation of the viscous substratum and the cover layer close to the graben, overprinting the outward directed movement of the graben flanks (Fig. 4b and d). The cover layer moved inward (max. 70 mm in 20 h) indicating an influx of silicone into the graben. Figure ~~4e and de and f~~ display the evolution of x-displacement d_x the cumulative x-displacement d_{xcum} extracted from a profile through the centre of the box. It shows that the inward movement in ~~E2~~Ev began during the late stages of the syn-extensional phase (after ~ 1.5 cm vertical displacement of the graben) and continued after extension stopped (post-extensional phase). Figure ~~4e and fg and h~~ display the averaged cumulative x-displacement $\overline{d_{xcum}}$ separately for the left (blue) and the right (green) half of the box. In both experiments, E1Eb and E2Ev, $\overline{d_{xcum}}$ increased linearly with progressive lateral extension. However, towards the end of the extensional phase of exp. E2, $\overline{d_{xcum}}$ was increased $\overline{d_{xcum}}$ increased stronger in Ev. This reflects the effect of the movement of the cover towards the graben structure and implies that this movement was, which was faster than the lateral extension.

265 4.1.2 Flank tilting

In experiment ~~T1Tb~~ (without silicone) in which only tilting of the flanks was applied, minor x-displacement ~~of the flanks can be observed (Fig. 5a and e)~~ can be observed at the margins of experimental surface resulting from rotational movement of the flanks (Fig. 5a and c). In ~~the experiment T2~~experiment Tv (including a viscous layer), the tilting of the flanks caused a widespread movement of the cover ~~layer~~ towards the centre of the box (max. 3 mm in 20 h) indicating a downward redistribution of the silicone (Fig. 5b and d). The evolutionary plot of the cumulative x-displacement d_{xcum} (Fig. 5f) shows that this downward movement began after a relatively small amount of tilting ($\sim 1^\circ$ after 5 h ~ 5 mm flank uplift). The downward movement only continued after tilting stopped, for a short extent (~ 2 h; continued after the tilting has been stopped (Fig. 5f and h).

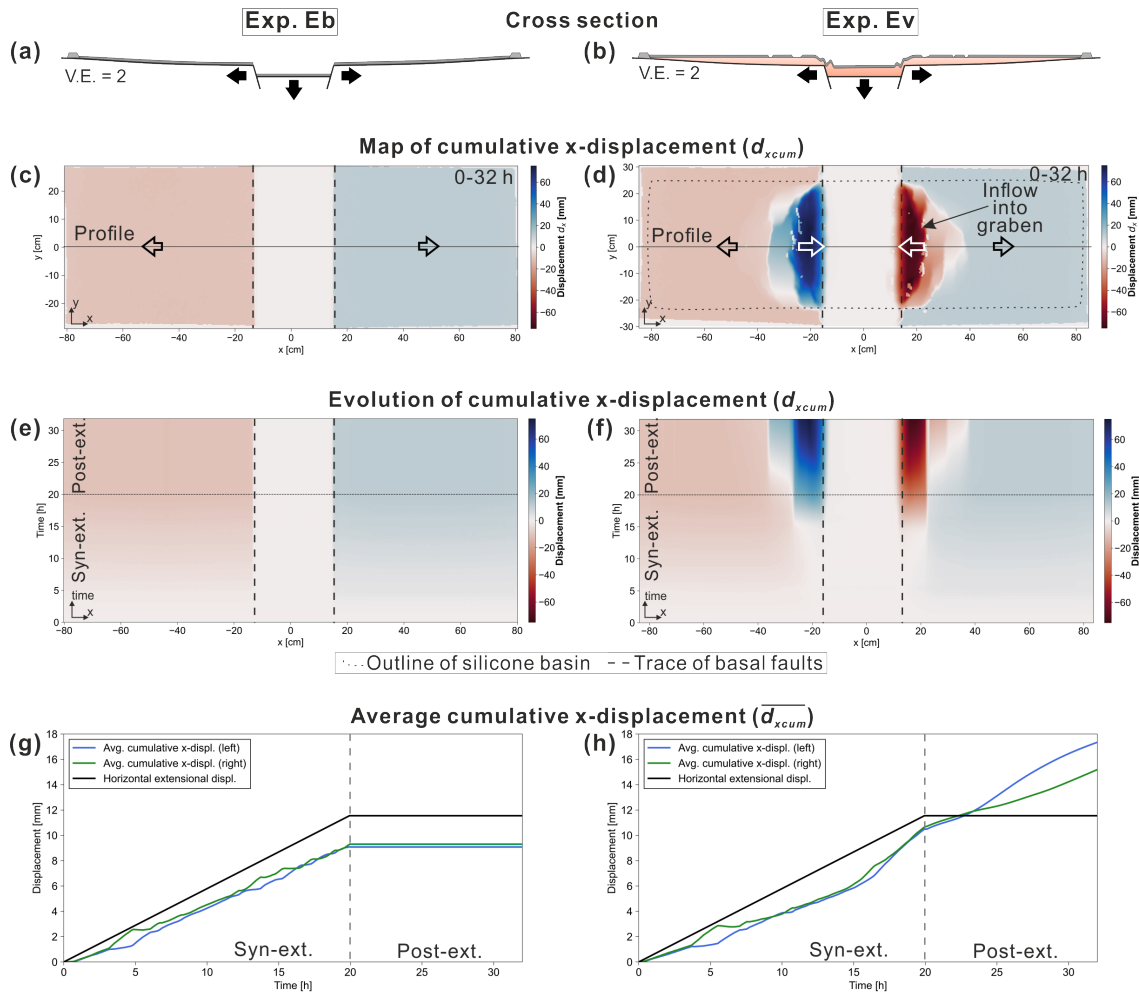


Figure 4. Results of the DIC of experiment **E1Eb** (without viscous substratum) and **E2Ev** (with viscous substratum) in which only **lateralbasal** extension with a rate of 1 mm h^{-1} was applied. The total vertical displacement of the central block d_e was 2 cm. **(a/b)** **Schematic cross sections through the centre of the experiments showing the final structures.** **(c/d)** **Top view of the cumulative x-displacement d_{xcum} at the end of the experiment (32 h).** Blue colours reflect rightward and red colours leftward movement. **(e/f)** **Incremental x-displacement (d_x) Evolutionary plots of d_{xcum} for experiment Eb and Ev, respectively,** extracted from a central profile (shown in (c) and (d)) and plotted against experimental time. **(g/h)** **Averaged cumulative x-displacement $\overline{d_{xcum}}$ for the left (blue) and the right (green) half of the box plotted against experimental time.** Note that the black line displays the lateral extension (max. $\sim 12 \text{ mm per side}$), which is less than the vertical movement of the central block (max. 20 mm).

4.1.3 Extension and tilting

In experiments **ET1ETb** and **ET2ETv**, graben extension and flank tilting were applied **as boundary conditions simultaneously** (Fig. 6a and b). **In ETb, x-displacement of the experimental surface is dominated by the lateral movement of the graben flanks**

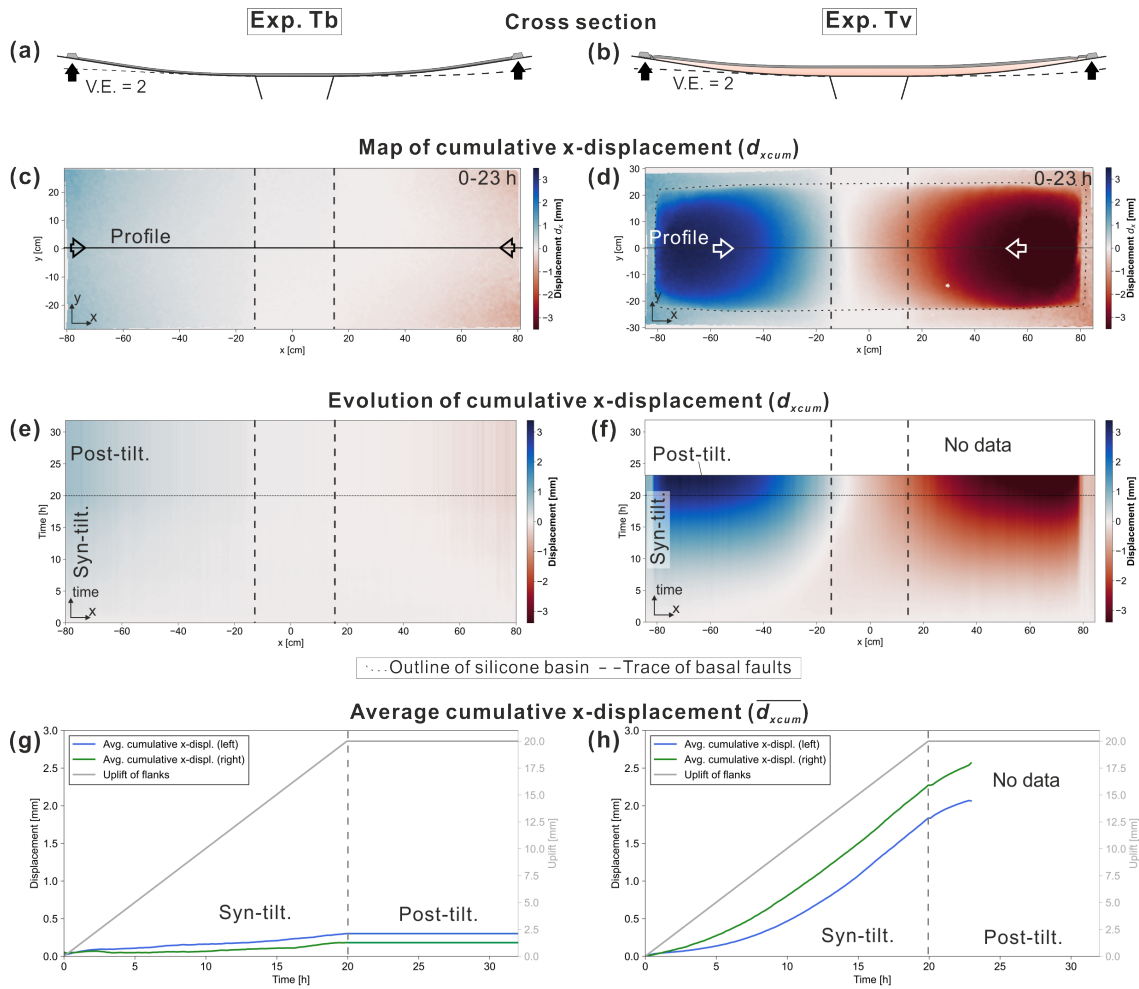


Figure 5. Results of the DIC of experiment **T1Tb** (without viscous substratum) and **T2Tv** (with viscous substratum) in which tilting was applied with a rate of 1 mm h^{-1} . The total vertical displacement of the outer edges of the flanking plates was 2 cm, which caused a maximum tilting of $\sim 5^\circ$. (a/b) Schematic cross sections through the centre of the experiments showing the final structures. (c/d) Top view of the cumulative x-displacement d_{xcum} at the end of the experiment (23 h). Blue colours reflect rightward and red colours leftward movement. (e/f) Incremental x-displacement (d_x) Evolutionary plots of the cumulative x-displacement d_{xcum} for experiment Tb and Tv, respectively, extracted from a central profile (shown in (c) and (d)) and plotted against experimental time. (g/h) Averaged cumulative x-displacement $\overline{d_{xcum}}$ for the left (blue) and the right (green) half of the box. The grey line presents the uplift of the flanks, which was equal for both sides. Note that shown results of Tv are shorter than in other experiments, since no data are available after 23 h.

(Fig. 6c and e) similar to the displacement patterns in Eb. With the presence of a viscous layer in ETv, the inward movement of the cover layer towards the subsiding graben affected the entire region of the uplifted flanks (Fig. 6d). Inward movement was at a maximum close to the basal normal faults ($\sim 110 \text{ mm}$ in 32 h) and reflects the downslope gliding of the cover and silicone. The

evolutionary map shows that the downslope movement appears to behave been activated at roughly half of the syn-extensional phase (~ 1 cm of subsidence of the graben and ~ 1 cm uplift $\sim 2.5^\circ$ tilting of the flanks) (Fig. 6f). During the post-extensional phase, downslope gliding continued. The averaged cumulative x-displacement $\overline{d_{x\text{cum}}}$ increased strongly towards the end of the syn-extensional phase in ETv and still increased linearly during the post-extensional phase (Fig. 6h). which This implies that the downward movement of the cover layer continued at a relatively high velocity.

4.2 Flat vs. tilted flanks

285 4.2.1 Displacement patterns

~~The effects of inclined flanks on the structural development is evaluated by comparing experiments E2 and ET2 (Fig. 7a and b). Figure 7e and d show that the total amounts of x-displacement $d_{x\text{cum}}$ were significantly higher and deformation affects a larger region, if the flanks were tilted contemporaneously with the basal extension. Especially during the post-extensional phase, widespread downward movement can be observed in the entire silicone basin in ET2, whereas deformation is still limited to regions close to the basal graben in E2 (Fig. 7e and f). In both experiments, the region affected by downward movement adjacent to the graben becomes wider after extension stopped. Furthermore, the total amount of d_x and, therefore, the displacement rates are considerably higher in the post-extensional phase (12 h) than in the syn-extensional phase (20 h). The evolutionary maps (Fig. 7g and h) demonstrate that the onset of inward movement into the central graben began towards the end of the syn-extensional phase in both experiments and did not occur noticeably earlier in ET2.~~

295 The effects of tilted flanks on the structural development is evaluated by comparing experiments Ev (only extension) and ETv (extension + tilting) (Fig. 7a and b). The total amounts of x-displacement d_x were significantly higher and deformation affects a larger region, if the flanks were tilted contemporaneously with basal extension (max. syn- + post-extensional displacement $d_{x\text{cum}} = 82$ mm in Ev and 110 mm in ETv; Fig. 7c and d). Particularly during the post-extensional phase, widespread downward movement occurred in the entire silicone basin in ETv, whereas deformation was limited to regions close to the basal graben in Ev (Fig. 7e and f).

In both experiments, the onset of inward movement above the graben flanks began towards the end of the syn-extensional phase and the regions affected by downward movement adjacent to the graben became wider after extension stopped (Fig. 7g and h). No displacement can be observed above the subsided graben block in Ev (Fig. 7g). In contrast, convergent movement began abruptly in the graben centre at the beginning of the post-extensional phase in ETv (Fig. 7h). This implies that influx into the graben centre is stronger, when flank tilting is applied simultaneously with basal extension.

305 Note that the incremental movement of the graben flanks is not fully continuous and symmetrical during basal extension (Fig. 4g and h). This likely results from mechanical locking of the graben block on the basal faults.

4.2.2 Strain patterns

~~The patterns of the normal strain in x-direction ϵ_{xx} (Fig. 10a) reveal that extensional fault zones developed at an offset of roughly 10 cm from the basal faults ('graben-edge extension') in E2, whereas compressional strain occurred directly above~~

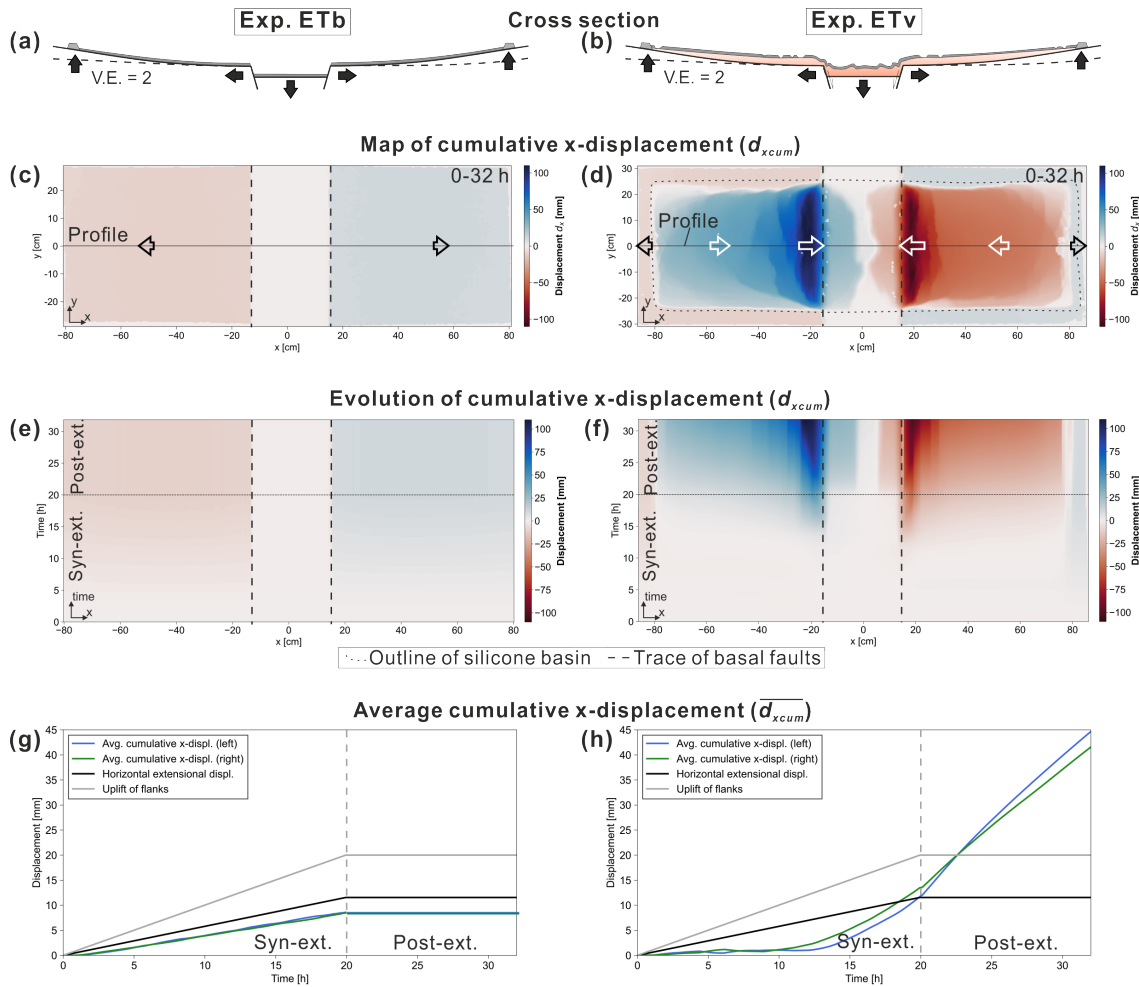


Figure 6. Results of the DIC of experiment [ET1ETb](#) (without viscous substratum) and [ET2ETv](#) (with viscous substratum). In both experiments extension and tilting with rates of 1 mm h^{-1} were applied. The total vertical displacement of the central block and the flanking plates was 2 cm. [The max. tilting of the flanks at the outer edges was \$\sim 5^\circ\$.](#) [\(a/b\) Schematic cross sections through the centre of the experiments showing the final structures.](#) [\(c/d\) Top view of the cumulative x-displacement \(\$d_{xcum}\$ \) at the end of the experiment \(32 h\). Blue colours reflect rightward and red colours leftward movement.](#) [\(e/f\) Incremental x-displacement \(\$d_x\$ \) Evolutionary plots of the cumulative x-displacement \$d_{xcum}\$ for experiment ETb and ETv, respectively, extracted from a central profile \(shown in \(c\) and \(d\)\) and plotted against experimental time.](#) [\(g/h\) Averaged cumulative x-displacement \$\overline{d_{xcum}}\$ for the left \(blue\) and the right \(green\) half of the box.](#)

the basal faults. The extensional fault zones propagated away from the graben during the post-extensional phase (Fig. 10b). In ET2, graben-edge fault zones were located closer to the graben ($< 8 \text{ cm}$). Additional extensional fault zones occurred at the outer margins of the silicone basin ('basin-margin extension') (Fig. 10a). During the post-extensional phase, the graben-edge extension propagated upslope, whereas the basin-margin extension migrated downward. Compressional zones evolved in the

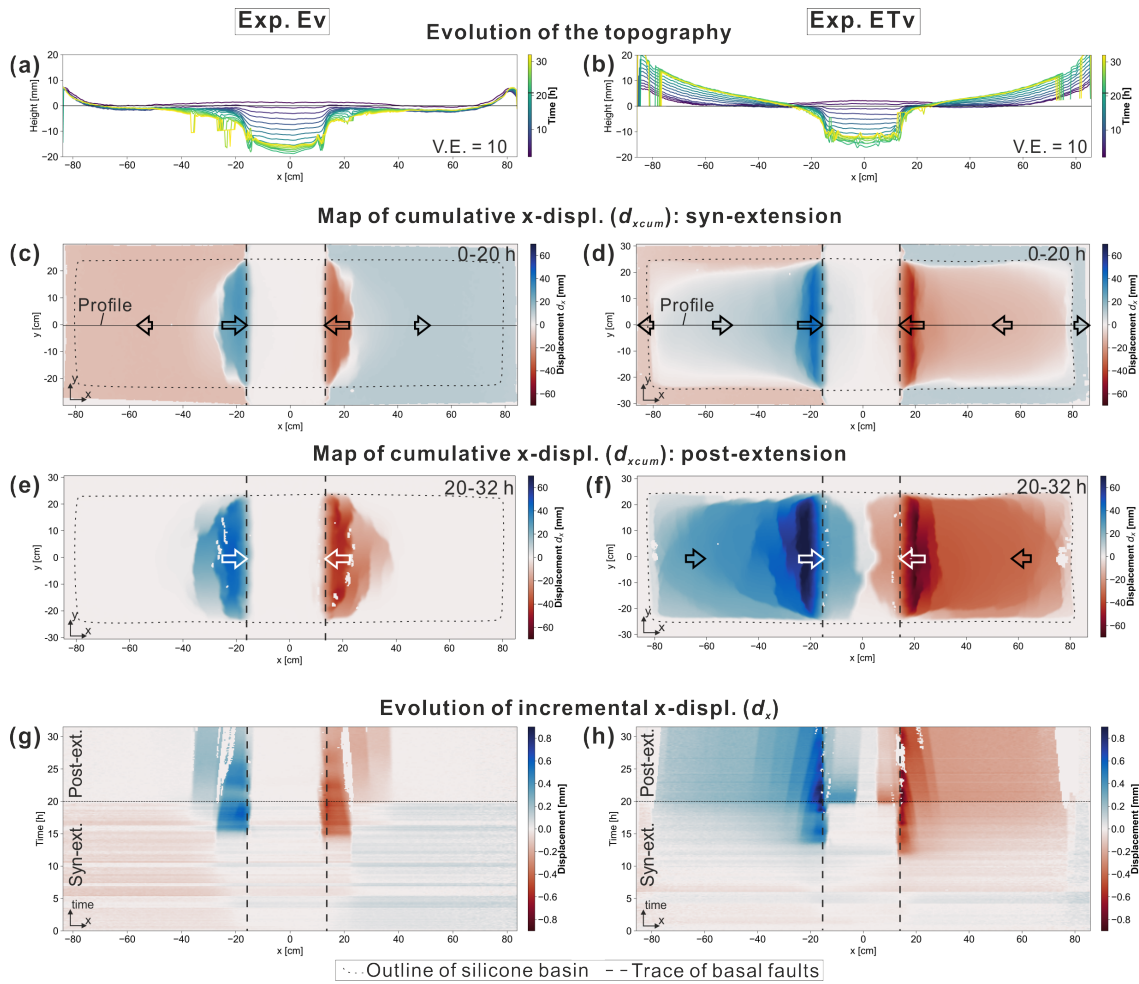


Figure 7. Comparison of displacement patterns between experiment **E2Ev** (only extension) and **ET2ETv** (extension and tilting). (a/b) Evolution of the **surface elevation extracted along a center profile topography**. Along a central profile (see c and d), the elevation of the experimental surface was extracted every 2 h. The artefacts in the lines result from exposed reflective silicone, which disturbs the DIC analysis. (c/d) Cumulative x-displacement $d_{x_{cum}}$ at the end of the syn-extensional phase. (e/f) Cumulative x-displacement $d_{x_{cum}}$ at the end of the post-extensional phase. (g/h) Evolutionary maps showing the cumulative x-displacement plots of the incremental x-displacement d_x extracted along a central profile (shown in (c) and (d)) and plotted against the experimental time. Blue means rightward and red leftward movement.

315 centre of the subsided graben, which reflects the influx of the silicone and downslope gliding of the sand cover into the graben. In contrast, compressional strain in E2 was limited to the regions close to the basal faults. The main difference between experiment E2 and ET2 is that the amount of the extensional strain at the graben-edge fault zones was higher in E2 (~ 2) than in ET2 (~ 1.5) and graben-edge extension was distributed over a wider region in E2 (Fig. 11). It also can be observed that

extended and dissected cover blocks are convexly curved in E2, whereas they strike roughly parallel to the basal faults in ET2

320 (Fig. 11a and c).

The strain patterns are analysed by means of the cumulative normal strain in x-direction e_{xxcum} . Here, extensional strain is defined as positive and compressional strain as negative (Fig. 8). These patterns reveal that extensional fault zones developed above the footwall blocks of the graben ~6 cm offset from the basal faults ('graben-edge extension') in Ev (Fig. 8a). During the post-extensional phase, additional extensional zones were initiated further away from the basal normal faults at offsets of up to

325 ~12 cm (Fig. 8c), which gradually moved towards the central graben (Fig. 8e). Compressional strain, which accommodated the extensional strain, was focused in straight, narrow regions directly above the ramps of the basal faults. These compressional regions remained stable throughout the experiment and only slightly increased in width during the post-extensional phase (Fig. 8e).

In ETv, graben-edge fault zones were located closer to the basal faults during the syn-extensional phase (<5 cm; Fig. 8b) than in Ev. Furthermore, the max. amount of e_{xxcum} in these zones was smaller in ETv (~1.5) than in Ev (~2.5). Until the end of the post-extensional phase, the region affected by graben-edge extension migrated upslope and widened to roughly 13 cm on each side (Fig. 8d). Additional extensional fault zones occurred at the outer margins of the silicone basin ('basin-margin extension'; max. $e_{xxcum} = 1.2$) (Fig. 8b and d). These migrated downslope and changed from straight to curved fault zones during the post-extensional phase (Fig. 8d and f). Similar to Ev, the compressional strain localized above the ramps of the basal

335 faults in ETv during the syn-extensional phase (max. $e_{xx} = \sim -1$; Fig. 8f). However, after basal extension stopped, additional curved compressional zones developed above the hanging wall graben block in ETv (Fig. 8d, f and h), which is different from the post-extensional evolution in Ev (Fig. 8c, e and g).

In summary, syn-extensional flank tilting in ETv has the following effects on the overall deformation patterns: 1. the formation of basin-margin extensional fault zones, which gradually migrate downslope; 2. decreased amounts of extensional strain at the graben-edge fault zones; and 3. the formation of compressional zones (buckle folding) above the hanging wall graben block during the post-extensional phase.

340

4.3 Effects of syn-kinematic sedimentation

In ~~ETS~~ETSv, extension and tilting of the basal parts was accompanied by regular intervals of syn-kinematic sand accumulation (Fig. 9a and b). Similar to ~~ET2~~ETv (equivalent experiment without sedimentation), a downslope gliding of the overburden

345 above the flanks can be observed (Fig. 9c). However, the total amount of cumulative x-displacement d_{xcum} at the end of the extensional phase is smaller in ~~ETS~~ETSv (max. ~10 mm) than in ~~ET2~~ETv (~30 mm). This difference is even higher for the total ~~d_x at the end~~ d_{xcum} of the post-extensional phase (~~ETS~~ETSv: 1.5 mm; ~~ET2~~ETv: 60 mm) (Fig. 9e). This implies that sequentially a sequential filling of the downslope depocentre with syn-kinematic sand significantly reduces downslope gliding of the cover.

350 Graben-edge extension and compressional strain in the graben centre were almost completely suppressed in ~~ETS~~ETSv (max. $e_{xxcum} = 0.2$; Fig. 9d and f). The total amount of strain at the basin-margin extensional fault zones (Fig. 9d and f) was also less (max. $e_{xxcum} = 0.8$) than in ~~ET2~~ETv (max. $e_{xxcum} = 1.2$; Fig. 8b and d). In ETSv, extensional strain is was

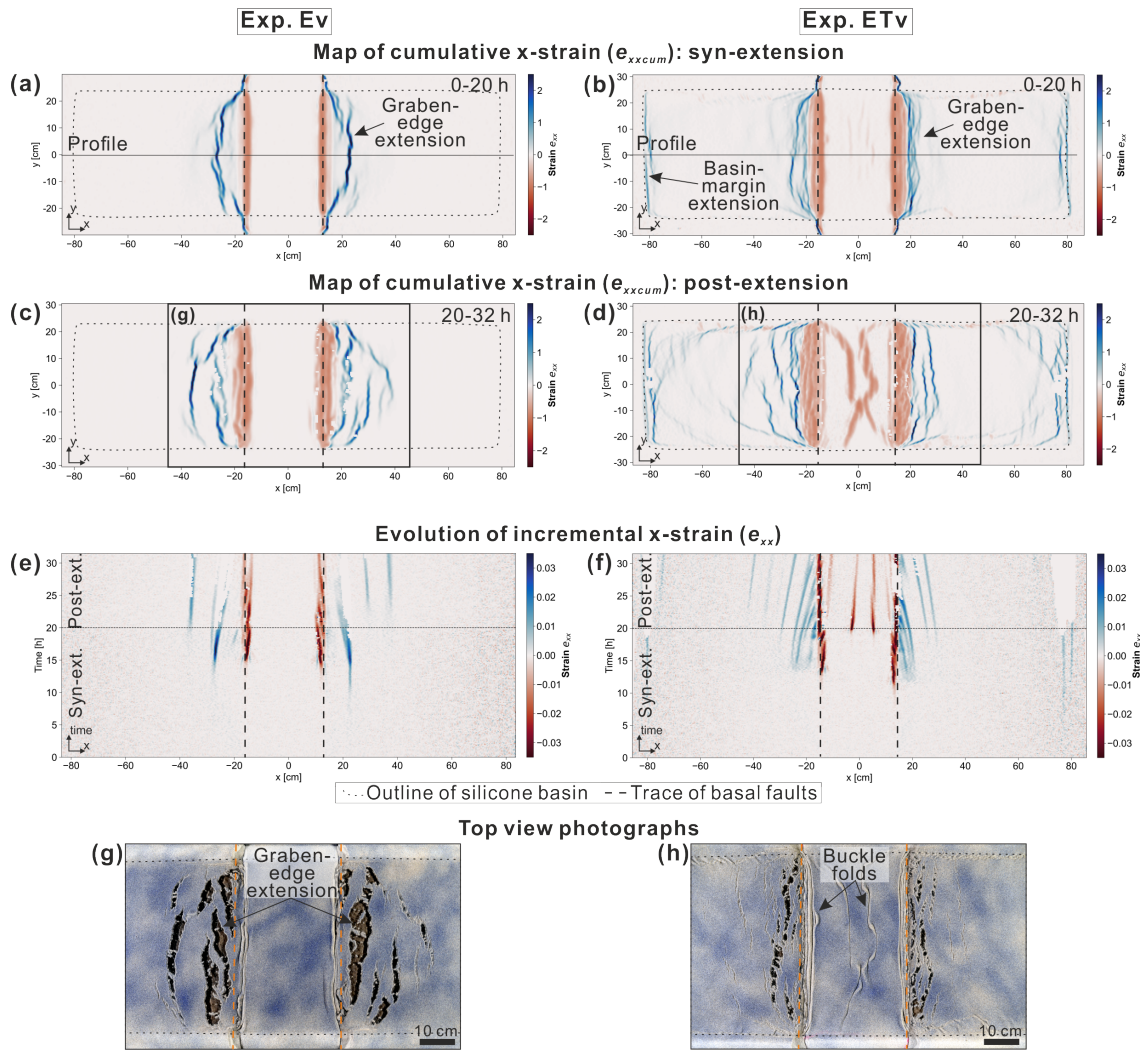


Figure 8. Comparison of strain patterns between **E2Ev** (only extension) and **ET2 ETv** (extension and tilting).

(a/b) Cumulative normal strain in x-direction e_{xxcum} at the end of the extensional phase of the syn-extensional phase. (c/d) Cumulative x-normal strain e_{xxcum} at the end of the post-extensional phase e_{xxcum} of the post-extensional phase. (e/f) Evolutionary maps showing plots of the incremental x-normal strain e_{xx} extracted along a central profile (shown in (c) and (d)) and plotted against the experimental time. Blue means extensional and red compressional strain. (g/h) Top-view photographs of deformation structures close to the central graben at end of the experiment.

focused on a1–2 discrete zones in **ETS1** and, hence, more localized compared to the widely distributed strain observed in **ET2 ETv** (Fig. 8f). During the post-extensional phase, downslope gliding was restricted to the upslope regions of the graben flanks (Fig. 9e). Diffuse, widespread compressional strain occurred in the central parts of the flanks (Fig. 9f) and not above the graben centre as in **ET2 ETv** (Fig. 8d).

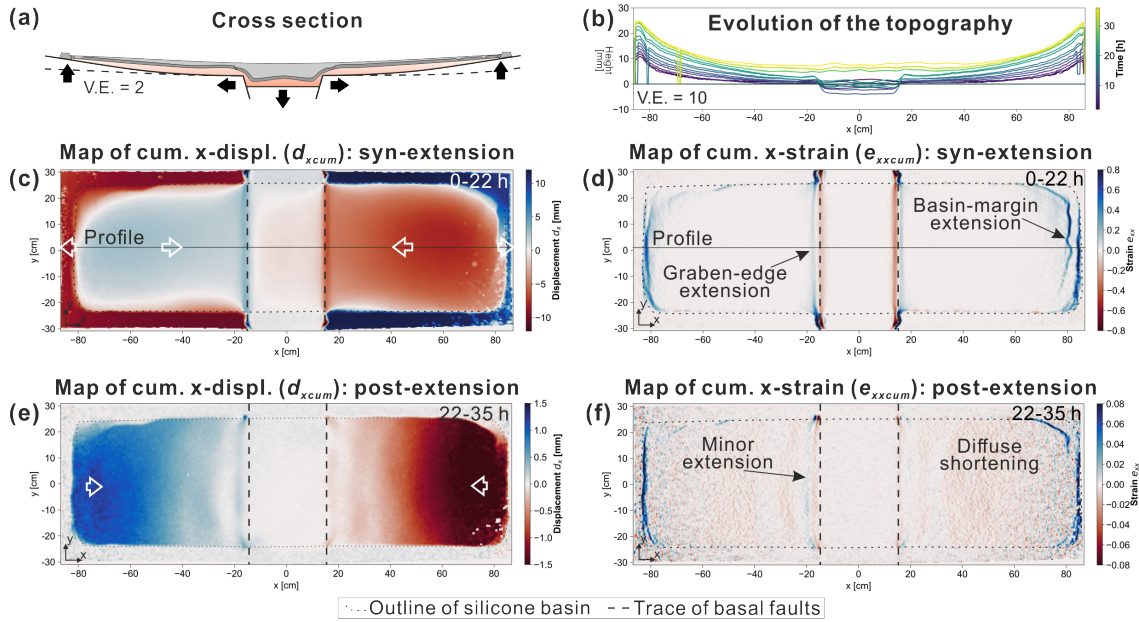


Figure 9. Displacement and strain patterns of **experiment ETS1ETSv** (extension + tilting + sedimentation). **(a)** Schematic cross section through the centre of the experiment showing the final structures. **(b)** Evolution of the topography. Along a central profile, the elevation of the experimental surface was extracted every 2 h. The artefacts in the lines result from exposed reflective silicone, which disturbs the DIC analysis. **(c)** Cumulative x-displacement d_{xcum} and (d) cumulative x-normal strain e_{xxcum} at the end of the syn-extensional (20 h) and post-extensional phase (32 h): of the syn-extensional phase. **(e)** Cumulative x-displacement d_{xcum} and (d) cumulative x-normal strain e_{xx} of the post-extensional phase. **(g)** Evolutionary map showing the incremental d_x and (h) the incremental e_{xx} plotted against the experimental time. **(e)** Averaged cumulative and (f) averaged cumulative for the left (blue) and the right (green) half of the box.

The evolutionary plots in Fig 10 show that **incremental displacement and strain reached their maximum at the end of the extensional phase and only slowly continued after extension and tilting was stopped.** most of the deformation occurred during the basal extension and rapidly decreased during the beginning of the post-extensional phase. This is also in contrast to **ET2ETv**, where maximum amounts of displacement and strain were achieved during the post-extensional downslope gliding (**Fig. 8f**).

Cross sections cut through the final experiment reveal that the structures of basin-margin extensional fault zones in the sand cover are characterized by synthetic rollovers and symmetric grabens (**Fig. 11a** and **b**). Strongest deformation occurred directly at the pinch out of the silicone layer. Aside from minor extensional thinning of the pre-kinematic layer, no deformation structures **can be observed** are visible close to the basal faults (**Fig. 11c**). This indicates that strain of the graben-edge extensional zone is insufficient to produce significant faulting.

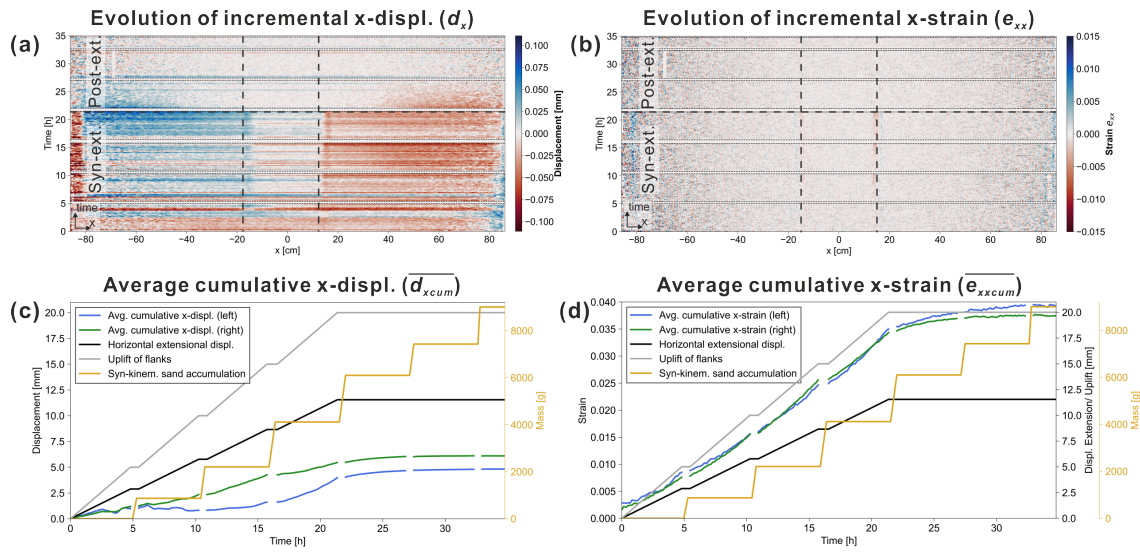


Figure 10. Displacement and strain evolution of ETSv (extension + tilting + sedimentation). (a) Evolutionary plots of the cumulative incremental x-displacement d_x and (b) the cumulative incremental x-strain e_{xx} extracted along a central profile (shown in Fig. 9c and d) and plotted against the experimental time. (c) Averaged cumulative x-displacement $\overline{d_{xcum}}$ and (f) averaged cumulative x-strain $\overline{e_{xxcum}}$ for the left (blue) and the right (green) half of the box.

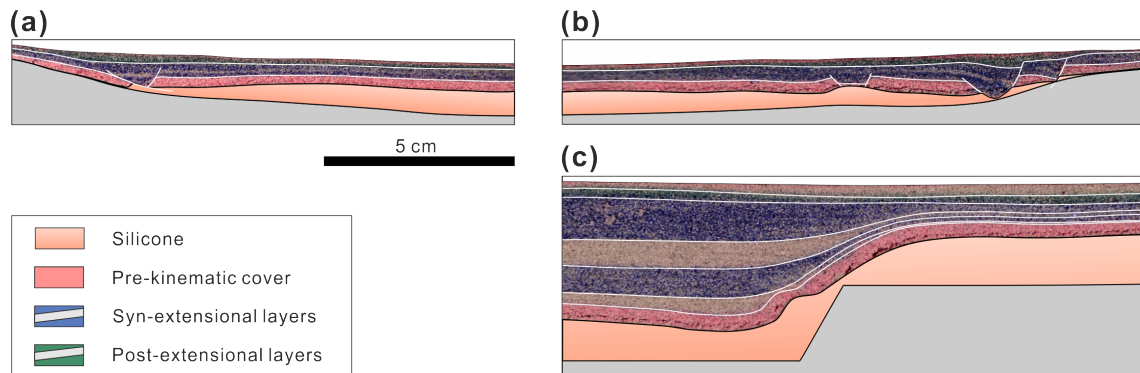


Figure 11. Cross sections cut through the centre of Exp. ETSv in which syn-kinematic sand accumulation was applied during and after basal extension. (a) Cross section of the extensional fault zone at the left basin margin. (b) Cross section of the extensional fault zone at the right basin margin. (c) Cross section of the syn-kinematic strata close to the right basal fault. Structures above the left basal fault were similar.

5 Discussion

5.1 Experimental ~~procedures~~structures and kinematics

The presented experiments on gravity-driven deformation in salt-bearing rift basins (SBRB) provide first insights into the influence of flank tilting on basin-wide deformation patterns. The comparison between experiments with and without tilting of the flanks illustrate how local deformation structures due to a sub-salt graben are affected by regional-scale flow of the viscous material.

When graben flanks remain flat (E_v), deformation is restricted to the area in the vicinity of the basal graben ~~even though thin-skinned extensional zones are decoupled from the basal normal faults (Fig. 8a)~~. (Fig. 12a). Distributed extensional fault zones in the cover developed offset from the basal normal faults above the footwall blocks of the graben. This implies that deformation in the cover layer was effectively decoupled from basal extension as it has been demonstrated by previous analogue modelling studies (Nalpas and Brun, 1993; Vendeville et al., 1995; Ge et al., 1997; Withjack and Callaway, 2000; Dooley et al., 2005).

When only tilting is applied (T_v), localized extension occurs at the basin margins, while diffuse shortening affects the cover in downslope regions (Fig. 12b). Such deformation patterns are similar to those in early post-salt phases of passive margin basins affected by slight, progressive tilting (Ge et al., 2019b).

~~Due to tilting of the flanks in experiment ET2, downward displacement of the brittle overburden is induced over the entire area underlain by a viscous layer (Fig. 12c). Total amounts of syn-extensional and post-extensional displacement were considerably larger compared to those in the experiment without tilting (E_2) (Fig. 7c and d). Nevertheless, extensional strain in the graben-edge fault zones is less in ET2 than in E_2 (Fig. 8a and b). These observations indicate that the amount of graben-edge extension is decreased by downslope gliding of the cover layer. In terms of kinematic balancing and cross-section restoration, this implies that there is a deficit in the balance between sub- and supra-salt extension close to the graben centre. Due to simultaneous basal extension and flank tilting (ET_v), the local (close to the graben) deformation is overprinted by a regional (entire basin scale) downward flow of the viscous material and gliding of the brittle cover (Fig. 12c). Total syn-extensional displacement was larger than in the experiment with only basal extension (E_v) (Fig. 7c and d). Nevertheless, extensional strain in the graben-edge fault zones is less in ET_v than in E_v (Fig. 8a and b) indicating that the amount of graben-edge extension is decreased by downslope gliding of the cover layer. In terms of kinematic balancing and cross section restoration, this implies that the amount of sub-salt extension is larger than that of supra-salt extension in the region close to the central graben and is partly balanced by thin-skinned extension at the basin margins. Thus, we suggest that the occurrence of basin-margin extensional fault zones and reduced supra-salt extension or even shortening close to the basal normal faults are diagnostic indicators for syn-extensional gravity gliding in natural rift basins.~~

~~During the post-extensional phase, downslope gliding continues across the entire graben flanks in ET2, which enhanced extension at the basin margin and caused shortening of graben-edge extensional fault zones as well as fold and thrust faults in the graben centre. We suggest that basin-margin fault zones and shortened graben fault zones, in particular if they are active~~

subsequent to the rifting, are diagnostic indicators for the influence of gravity gliding on the structural dynamics in natural rift
400 basins.

During the post-extensional phase in ETv, downslope gliding continued across the entire flanks, which enhanced extension
at the basin margins. The basin-margin extensional domain, which has started as 1-2 localized extensional zones, began to
expand and migrated downslope (Fig. 8d and f). This evolution is similar to that in analogue models of passive margins in
which the extensional domain gradually moves basinward during progressive basin tilting (e.g. Ge et al., 2019b). Furthermore,
405 the displacement of the cover into the basin centre was significantly larger in ETv than in Tv (only tilting) and in Ev (only
extension). Therefore, it can be suggested that a substantial gravity potential is only produced, if graben subsidence and flank
tilting are combined.

Shortening structures above the hanging wall graben centre only occurred in ETv (Fig. 8d), which lead to buckle folding
and thrusts that are similar to those in downslope compressional domains of passive margin salt basins (e.g. Brun and Fort,
410 2004). Since this shortening was first initiated at the beginning of the post-extensional phase (Fig. 8f), it can be suggested that
during basal extension, enough space is created in the basin centre to accommodate the displacement of the downward gliding.
Although inward flow into the graben centre occurred in Ev as well, shortening was limited to the region close to the basal
faults. Hence, supra-salt shortening in the graben centre as observed in ETv might be a general indicator for post-extensional
gravity gliding in SBRB. However, this is speculative, because the degree of gravity gliding and resulting extensional and
415 compressional strain are highly dependent on the ratio between basal extension rate and tilting rate (among other parameters)
as it can be inferred from passive margin analogue experiments with various tilting rates (e.g. Ge et al., 2019b). Furthermore,
supra-salt shortening in rift basins can also be a result of post-rift basin inversion (e.g Dooley and Hudec, 2020). Therefore, the
exact timing of shortening has to be determined when relating it to gravity gliding.

~~The experiment with syn-kinematic sedimentation (ETS1) demonstrates that the ability of downward gliding is reduced
420 by sediment accumulation in the basin centre. In particular during the post-extensional phase, gravity driven deformation
decreased rapidly after the first post-extensional step of sand accumulation. This suggests that gravity gliding is not a compelling
processes in SBRB with tilted flanks and might depend on the amount and accumulation rate of sediments in the basin centre.~~

In the experiment with syn-kinematic sedimentation (ETSv), amounts of strain and displacement were considerably less
(Fig. 12c and d) than in the equivalent experiment without sedimentation (ETv) (Fig. 7d and f). No shortening took place in the
425 graben centre in ETSv. In contrast, gravity gliding was accommodated by reduced extension at the graben edges and by diffuse
shortening above the downslope flanks (Fig. 9d and f). This demonstrates that the ability of downward gliding is significantly
reduced by sediment accumulation in the basin centre. Such a dynamic resistance is suggested to act in passive margin salt
basins as well, where the gravity potential has to be large enough to overcome the strength of the downslope overburden in
order to induced downward gliding and spreading (Gemmer et al., 2004; Rowan et al., 2004). After basal extension and flank
430 tilting was stopped in ETSv, gravity-driven deformation decreased rapidly. In natural rift basins, however, thermal subsidence
and flank tilting continues after rifting ceased such that gravity potential is further increased and gliding might last longer.

In ETSv, extensional strain at the basin margins and the graben edges occurred at 1 to 2 discrete fault zones, whereas it was
distributed over several branched fault zones in ETv (without sedimentation) (Fig. 12c and d). This implies that sedimentary

loading localizes the strain and causes subsidence in discrete extensional fault zones, whereas the initiation of new extensional
435 zones is suppressed. Similar effects of syn-kinematic sedimentation on the strain patterns has been observed in analogue models
of rift evolution (Zwaan et al., 2018) and passive margin salt basins (Ge et al., 2019a). In passive margin basins, syn-kinematic
sedimentation during gravity gliding initially leads to symmetrical graben structures, which progressively transform into listric
growth faults (e.g. Mauduit et al., 1997; Krézsek et al., 2007). Such structures also developed at the tilted basin margins in
ETSv (Fig. 11) and, therefore, represent clear indicators for gravity gliding in SBRB.

440 Experiment ~~ET2~~ETv (without sedimentation) and ~~ETS4~~ETSv (including sedimentation) might represent end-member types
of SBRB of an underfilled basin (low sedimentation rates) and ~~overfilled~~filled up basin (high sedimentation rates). Gravity
gliding is favoured in underfilled rift basins or basins experiencing fast flank tilting relative to sedimentation rates, respectively.
In such a case, the kinematic evolution and structural patterns are similar to passive margin basins in which basin tilting is
the main driving factor for salt tectonics (e.g. Mauduit et al., 1997; Fort et al., 2004; Brun and Fort, 2011; Adam et al., 2012;
445 Ge et al., 2019a,b). High sediment accumulation rates, in contrast, might effectively prevent gravity gliding which suggests
that gravity gliding is not a compelling processes in SBRB with tilted flanks, i.e. with a noticeable gravity potential. This is
a major difference to passive margin basins, where fast sediment supply from the continent accelerates salt tectonic processes
as it has been shown by numerous analogue modelling studies (e.g. Ge et al., 1997; Mauduit et al., 1997; McClay et al.,
2003; Vendeville, 2005; Krézsek et al., 2007; Brun and Fort, 2011). However, ~~such a conclusion require a more comprehensive~~
450 ~~parameter study in which various ratios of sedimentation vs. tilting rates are tested as it has already been undertaken for passive~~
~~margin setting~~ the exact conditions under which gravity gliding is completely blocked or still possible in SBRB have not been
investigated in our study and require a more comprehensive parameter study in which various ratios of sedimentation vs. tilting
rates are tested.

5.2 Limitations of the apparatus and the procedure

455 In our experimental setup, basal extension is focused on two parallel faults, which is a simplification of a natural rift basin.
Natural rift basins and the process of crustal extension include greater complexities than presented in our experimental setup.
Rift basins usually consist of overlapping fault systems with normal faults, fault intersection, ramps, strike-slip segments and
intrabasinal ridges (e.g. Fazlikhani et al., 2017). Such diverse fault patterns in the sub-salt basement determines the distribution
of the original salt, its post-depositional redistribution and localization of salt structures (e.g. Koyi et al., 1993; Dooley et al.,
460 2005; Dooley and Hudec, 2020). Similarly, the flanks in our model setup have an even base, whereas in nature ~~small-offset~~
~~faults occur in these regions due to, for instance, widening of the graben through time.~~the sub-salt base is often perturbed by
an irregular relief resulting from, for instance, fault steps and volcanic edifices (e.g. Dooley et al., 2017; Ferrer et al., 2017).
Basal discontinuities influence the strain localization in the viscous and brittle layers mainly because of abrupt changes of the
thickness of the ~~ductile~~viscous layer (Gauillier et al., 1993; Dooley et al., 2017; Ferrer et al., 2017; Pichel et al., 2020) so that
465 basin-margin extension zones are distributed above such basal steps rather than at the marginal pinch-out of the viscous layer.

Natural rift basins often possess asymmetrical half graben shapes characterized by a large-offset normal fault on one side
and a smoothly tilted, slightly faulted flank on the other. In such half grabens, decoupled extension, probably caused by gravity

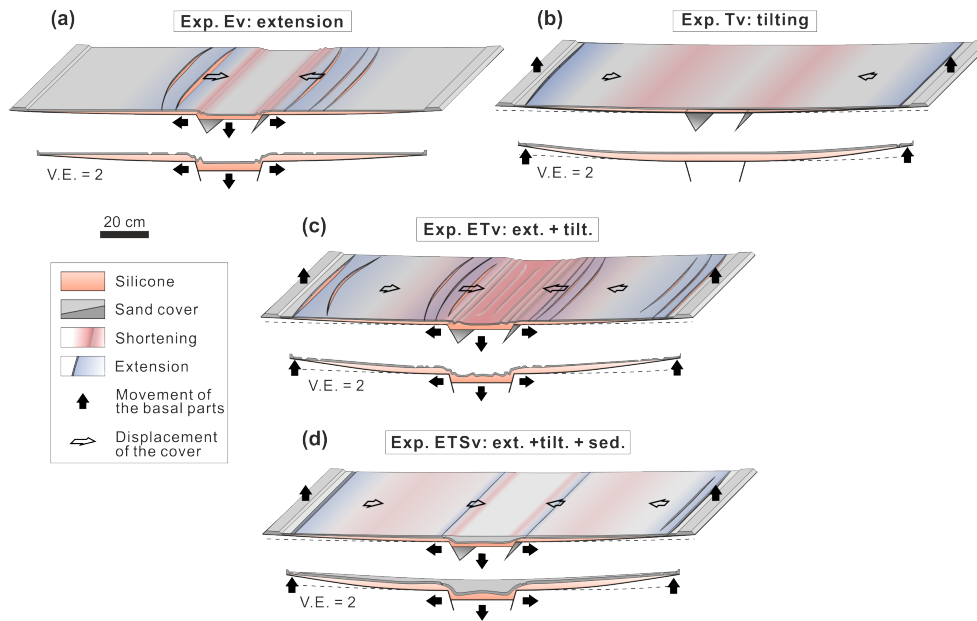


Figure 12. Sketches summarizing the main results of this experimental study. (a) Exp. Ev: In the experiment with only basal extension, widespread extension occurred in the cover layer above the footwall blocks, whereas shortening was limited to the regions close to the tips of the basal normal faults. (b) Exp. Tv: When only flank tilting was applied, discrete fault zones developed at the upslope basin margins, which were balanced by diffuse shortening further downslope. (c) Exp. ETv: Due to simultaneous graben extension and flank tilting, extensional strain at the graben edges was overprinted by basin-wide gravity gliding. This led to shortening in the graben centre and thin-skinned extension at the basin margins. (d) Exp. ETSv: Syn-kinematic sedimentation during basal extension and flank tilting caused localization of extensional strain at the graben edges and the basin margins. Furthermore, diffuse shortening occurred above the ramps of the basal faults and the downslope graben flanks.

gliding, can occur above the slightly tilted flank (e.g. Coward and Stewart, 1995; Thomas and Coward, 1996; Dooley et al., 2003; Jackson and Larsen, 2009; Ge et al., 2017; Lymer et al., 2018; Roma et al., 2018). Such alternative basin shapes **cannot** be modelled here, but should be considered and discussed when interpreting the modelling results have not been modelled in this study. Our experimental approach is to initially simplify the geometrical boundary conditions, and in later experimental studies, further complexities can then be introduced, such as static structural discontinuities can be inserted in the sub-silicone base to examine lateral variation of the strain patterns depending on initial thickness of the viscous layer.

A technical issue can be observed in Fig. 7g and h showing the incremental lateral movement of the flanks. Especially at the beginning of the extensional phase, this movement is not fully continuous and symmetrical on both sides, which likely results from mechanical locking of the graben block on the basal faults. This effect cannot be avoided even by using lubricants on the faults. Finally, the cumulative displacement of the extensional movement over longer time periods is smooth and regular as observed in Fig. 7e and d Fig. 4e and f.

5.3 Application to nature and outlook

480 In many SBRB worldwide (Tab. A1), the basin floor at the flanks is inclined and thin-skinned extensional faults occur in the supra-salt cover at the basin margins. The seismic examples shown in Fig. 13 resemble the basin-margin extensional fault zones observed in our experimental cross sections of experiment ETSv (Fig. 11). In our study such structures only occurred in experiments applying tilting. Thus, it can be proposed that these natural examples were also formed partly by gravity-driven extension. If so, our experimental results suggest that these natural faults should have been activated in the late stage of the rifting in the main graben and ~~show continued active faulting post-rifting~~ might have remained active after rifting ceased. Such temporal relationships have to be evaluated in detail by cross section balancing for specific rift basins.

~~The tectonic and sedimentary conditions in the basins of the shown examples were very different. One of the most prominent examples of overlapping gravitationally and tectonically driven supra-salt extension can be found at the western flank of the northern Central Graben (North Sea). Rollover structures and separated overburden blocks ('rafts') developed above the tilted outer platforms of the graben contemporaneously with the Triassic rifting phase of the graben (Fig. 13a) (Hodgson et al., 1992; Buchanan et al., 1996; Penge et al., 1999; Stewart and Clark, 1999). At the North Iberian rifted margin, rollovers detached on pre-rift Triassic evaporites during hyperextension in the Cretaceous (Fig. 13b) (Jammes et al., 2010; Lagabrielle et al., 2010). At the eastern flank of the Red Sea basin basinward listric faults developed on top of the mid-Miocene salt detachment coeval to the development of a fold belt and salt canopies in the basin centre (Fig. 13c) (Heaton et al., 1995). The amount of tectonic stretching and subsidence were relatively large in these three examples — the Iberian margin and the Red Sea Basin translated to early stage passive margins after mobilization of the salt. This probably resulted in steep slopes of the basin floor early on and a basin centre underfilled by sediments. We suggest that these geological conditions promoted intense gravity gliding. Nevertheless, clear evidence for syn-gliding contractional structures are lacking, except for the Red Sea Basin (Heaton et al., 1995). Extensional fault zones developed above gently dipping flanks of the Polish Basin from Late Triassic to Jurassic, in particular along the NE flank located above the East European Craton (Fig. 13f) (Krzywiec, 2006a; 2012). At the northern edge of the Norwegian-Danish Basin, triangular diapirs of Permian salt and tilted blocks in the overburden are supposed to be triggered by an inclined basin floor during the Triassic (Fig. 13e) (Geil, 1991; Clausen et al., 2015). In these basins, however, reactive diapirs and deeply subsided minibasins evolved in the basin centre coeval to basin-margin extension. Thus, hints for compressive structures due to gravity gliding are difficult to find here as well. Geil (1991) proposed a conceptual model, which stated that loading-driven upward directed salt flow in the basin centre took place contemporaneously to gravity-driven downward salt flow at the inclined graben flanks. In the Sole Pit Basin, decoupled extensional grabens and listric faults overlie a slightly faulted salt base at the basin margins (Fig. 13f). These extensional structures coincided with the formation of compressive, salt-cored anticlines in the basin centre during the Triassic (Hughes and Davison, 1993; Coward and Stewart, 1995; Stewart and Coward, 1995). In the Ems Trough (NW-Germany) and the south-eastern North Sea, the pre-kinematic overburden (lowermost Lower Triassic) was separated into individual blocks potentially due to raft tectonics (e.g. Thieme and Roekenbauch, 2001; Mohr et al., 2005; Vackiner et al., 2013). Mohr et al. (2005) determined that this~~

~~coincided with the formation of salt anticlines further downslope. In these examples, upslope extension is partly balanced by downslope compressional structures.~~

515 One of the most prominent examples of overlapping gravitationally and tectonically driven supra-salt extension can be found at the flanks of the northern Central Graben (central North Sea) (Fig. 1b). Rollover structures and separated overburden blocks ('rafts') developed above the tilted outer platforms of the graben (Fig. 13a) contemporaneously with compressional structures, such as salt anticlines and squeezed diapirs, in and near the graben centre (Gibbs, 1984; Hodgson et al., 1992; Bishop et al., 1995; Buchanan et al., 1996; Penge et al., 1999; Stewart and Clark, 1999; Stewart, 2007). It is suggested that major gravity-driven salt movement coincided with the main period of rifting during Jurassic to Early Cretaceous times (Bishop et al., 520 1995; Buchanan et al., 1996; Stewart and Clark, 1999). However, it has also been postulated that gravity gliding affected the post-salt depositional patterns already during the Late Triassic initial rifting phase (Penge et al., 1999). In the Pyrenean rift system and the Basque-Cantabrian basin (North Iberian margin), Early to Late Cretaceous 'hyperextension' caused gravity-driven extension of the cover of the Late Triassic evaporites (López-Mir et al., 2015; Saura et al., 2016; Teixell et al., 2016; Lagabrielle et al., 2020). Rollovers and listric growth faults are documented in outcrops and seismic data (Fig. 13b) and indicate 525 large-offset, thin-skinned extension at the southern basin margin (McClay et al., 2004; Jammes et al., 2010; Lopez-Mir et al., 2014; Saura et al., 2016; Labaume and Teixell, 2020). It is suggested that in the basin centre previously formed passive diapirs and the supra-salt cover was folded and overthrust partly accommodating upslope extension (López-Mir et al., 2015; Labaume and Teixell, 2020). The amount of tectonic stretching and subsidence were relatively large in these two examples – the Iberian margin translated to early stage passive margins after mobilization of the salt. This likely resulted in steep slopes of 530 the basin floor and a basin centre underfilled by sediments, which promoted gravity gliding. Hence, the style of salt tectonics in these basins was similar to that in many passive margin salt basins and as observed in our exp. ETv.

The Sole Pit and Silverpit Basin (southwestern North Sea) are confined by extensional fault zones, which overlie a slightly faulted salt base (Fig. 13c) (e.g. Hughes and Davison, 1993; Allen et al., 1994; Coward and Stewart, 1995; Griffiths et al., 1995; Stewart and Coward, 1995; Stewart et al., 1996). These extensional structures simultaneously developed with compressive 535 buckle folds in the basin centre during the Early Triassic–Jurassic rifting. Observed supra-salt extensional displacement is larger than that of the sub-salt faults at the basin margins and, therefore, can be related to syn-rift gravity gliding (e.g. Jenyon, 1985; Allen et al., 1994; Coward and Stewart, 1995; Griffiths et al., 1995; Stewart et al., 1996). Similar relationships between thin-skinned extension at basin margins and buckle folding and salt pillow growth further downslope are documented in the western (e.g. Thieme and Rockenbauch, 2001; Mohr et al., 2005; Vackiner et al., 2013) and the north-eastern part of the North 540 German Basin (Ahrlichs et al., 2020). The examples show that gravity gliding can also occur in less extended rift basins, whereas the amount of gravity-driven extension is minor as well.

The South Viking Graben and the Egersund Basin (northern North Sea) are salt-bearing half grabens in which the hanging wall graben centre was rotated during rifting (Geil, 1991; Thomas and Coward, 1996; Jackson and Larsen, 2009; Tvedt et al., 2013; Stewart, 2014; Jackson and Lewis, 2016). At the upslope edges of the rotated grabens, thin-skinned extensional growth 545 faults (Fig. 13d and e) are supposed to be triggered by gravity gliding (Jackson and Larsen, 2009; Tvedt et al., 2013; Jackson and Lewis, 2016). In the South Viking Graben, thin-skinned extension was initiated at a tilting angle of $<2^\circ$ and slightly delayed

from the Middle Jurassic main rifting phase. At the edge of the Egersund Basin (Sele High), gravity-driven supra-salt extension occurred during the Early Cretaceous and, therefore, also delayed from the Late Jurassic rifting (Jackson and Lewis, 2016). These observations of the timing of gravity gliding are supported by our experiments (ETv, ETSv) in which gravity gliding started later than and continued after the basal extension.

550 Extensional fault zones developed above gently dipping flanks of the Polish Basin from the Late Triassic to Jurassic, in particular along the NE flank located above the East European Craton (Fig. 13f) (Krzywiec, 2006a, 2012). In this basin, however, reactive diapirs and deeply subsided minibasins evolved in the basin centre so that it is difficult to evaluate the effects of gravity gliding. We can only speculate that the basin-margin extension was balanced by reduced supra-salt extension in the basin centre as it was observed in exp. ETSv (Fig. 9).

Listric growth faults developed in the post-salt strata on the western flank of the Horn Graben (Southern North Sea) (Fig. 13g) (Best et al., 1983; Baldschuhn et al., 2001). However, the underlying sub-salt base is only slightly inclined and faulted. Thus, it is suggested that supra-salt fault zones are rather detached from large-offset sub-salt normal faults in the nearby graben structure and purely tectonically driven (Nalpas and Brun, 1993; Stewart and Coward, 1996). This example demonstrates that there are structural similarities between tectonically and gravitationally driven thin-skinned fault zones underlining the necessity to find distinct structural and kinematic characteristics of both drivers.

560 These examples imply that there are three different scenarios of the interaction between gravity gliding and crustal extension: (1) Gravity gliding took place during ~~strong rifting~~the main phase of rifting but no or minor shortening occurred in the basin centre, since here ~~sub-salt faulting created sufficient space to accommodate and hide shortening~~ sub-salt extension was sufficient to accommodate and balance upslope, supra-salt extension (probably in the northern Central Graben), (2) Gravity gliding is superimposed with minor rifting and caused basin-margin extension balanced by basin-centre shortening (e.g. in the Sole Pit Basin), and (3) no or only minor gravity gliding occurred, and the thin-skinned extensional structures were mainly tectonically driven (e.g. in the Horn Graben). In order to investigate these scenarios and their governing parameters and processes, a comprehensive analogue modelling study has to be performed, modifying variables such as extension rate, tilting rate, layer thickness and timing of tilting in relation to extension.

6 Concluding remarks

Our new experimental apparatus successfully reproduced the overlapping influence of tectonic extension and gravity gliding on salt flow and salt structure evolution in salt-bearing rift basins. The apparatus is suitable for simulating crustal extension within a symmetric graben structure and thermal subsidence, which is represented by vertical uplift of the graben flanks. Since basal extension and flank tilting can be controlled separately, the effects of both processes can be investigated separately or combined. ~~Furthermore, we applied mixtures of quartz sand and silicate cenospheres as analogue for brittle overburden sediments. This enabled us to vary the ratio between overburden and viscous layer density so that the process of density inversion due to compacting overburden sediments can be produced.~~

The preliminary experimental study shown here reveals primary structural and kinematic differences depending on the presence of flank tilting [in salt-bearing rift basins](#). When the graben flanks remain flat during extension, deformation is concentrated in regions close to the basal graben structure. Extensional fault zones develop in the footwall, while overburden above the basal normal faults slides into the graben. When the graben flanks become tilted simultaneously to graben extension, similar deformation patterns occur close to the graben. However, they are overprinted by gravity gliding inducing widespread basin-wide downward movement of the overburden into the central graben. Consequently, additional extension localizes at the upslope basin margins, which is enhanced after basal extension and tilting have stopped. This extension is accommodated in the graben region by reduced amounts of extension on the graben-edge fault zones and compressional strain above the graben bounding faults. If syn-kinematic sedimentation is included in the experimental procedure, lateral downward directed displacement is significantly reduced especially during the post-extensional phase. Furthermore, strain in the margin-edge extensional zones is localized on less, discrete faults.

Observed basin-margin extensional structures in the preliminary experiments resemble typical thin-skinned extensional structures occurring at the flanks of many salt-bearing rift basins with inclined sub-salt bases. Such diagnostic structures indicate that gravity gliding might play an important role in the post-salt structural development [of such rift basins](#).

Code availability. No codes are available

Data availability. Data of the experiments presented in this study will be provided in a data publication on pangaea.de. The data was submitted and is currently processed (Nr.: PDI-27047).

Code and data availability. No codes are published along with this manuscript

Appendix A: Examples of salt-bearing rift basins

[In Tab. A1, basic geological parameters of numerous salt-bearing rift basins worldwide are listed, such as age of the salt layer, timing of rifting, layer thicknesses, width of the basin, maximal amount of tilting and maximal displacement of basal normal faults. Average values of these parameters serve as boundary conditions for the setup and the procedure of the presented analogue models \(Fig. 2; Tab. 1\).](#)

Appendix B: Scaling

~~For the analog model to be representative of the natural prototype, models have to be geometrically, dynamically, kinematically and rheologically similar to their natural prototype (Hubert, 1937; Ramberg, 1981; Weijermars and Schmeling, 1986; Davy and~~

Table A1. Literature-based data compilation of worldwide salt-bearing rift basins. The thickness of the salt layer refers to the estimated maximum thickness of the original evaporitic layer. The pre-kinematic layer is defined as the sedimentary layer accumulated after the salt deposition and before the first post-salt rifting phase. * denotes basins in which gravity-driven salt tectonics has been proposed or proofed.

Basin	Age of salt layer	Pre- & syn-salt rifting	Post-salt rifting	Thickness salt [m]	Thickness prekin. cover [m]	Max. offset subsalt faults	Max. tilting of the basin floor [°]	Width of [km]	References
Azraq-Mauldon Basin*	Late Triassic	Late Permian-Early Triassic	Late Jurassic	2700	1000	4000	9	80	Lagabille et al. (2010); Saptury et al. (2019); Labenne and Tevell (2020)
Atlas Basin	Late Triassic	Late Permian-Late Triassic	Early Jurassic	1500	100	3500	3	100	Martin-Martin et al. (2017); Moragas et al. (2017); Tronzi et al. (2017); Vergés et al. (2017)
Basque-Cantabrian Basin*	Late Triassic	Late Permian-Early Triassic	Late Jurassic	2700	500	3000	7	80	Bodego and Aguirreabala (2013); Lopez-Mir et al. (2015); Saura et al. (2016); Camaró (2020)
Central Graben (North)*	Late Permian	Early Permian	Early-Mid Triassic	2000	1000	500	5	180	Hodgson et al. (1992); Holland et al. (1993); Buchanan et al. (1996); Penge et al. (1999); Duffy et al. (2002)
Central Graben (Central)	Late Permian	Early Permian	Mid Triassic	1500	500	1000	2	40	Tanveer and Kongsgård (2009); Ge et al. (2017)
Central Graben (South)	Late Permian	Early Permian	Early Triassic	1500	300	3000	3	100	Geluk (2005); Wong et al. (2007); Van Winden et al. (2018)
Dniepr-Donets Basin	Late Devonian	Late Devonian	Early Carboniferous	2000	300	5000	6	150	Kusznir et al. (1996); Stovba et al. (1996); Stovba and Stephenson (2003)
East Texas Basin*	Mid Jurassic	Permian	mid Upper Jurassic	800	600	1000	4.5	80	Seni and Jackson (1984)
Egersund Basin*	Late Permian	Permian	Mid Jurassic	800	700	2500	5.5	40	Lewis et al. (2013); Tvedt et al. (2016); Jackson and Lewis (2016); Ge et al. (2017)
Emu Trough*	Late Permian	Permian	Mid Early Triassic	900	300	600	2.5	70	Rudale et al. (2005); Baldehahn et al. (2001); Srozyk et al. (2017)
Glauksjøfjorden	Permian	Early Permian	Early Jurassic	2500	800	4000	0.7	200	Maysenko et al. (2005, 2017); Warsitzka et al. (2017)
Halen Terrace	Mid-Late Triassic	Late Permian-Early Triassic	Mid Jurassic	1000	1200	2000	4	80	Pascoe et al. (1990); Dooley et al. (2003); Coleman et al. (2017)
Horn Graben*	Late Permian	Permian	Early Triassic	1500	400	4000	5	100	Vejthak (1990); Best et al. (1983); Kockel et al. (1995)
Lamreczye Trough*	Late Permian	Permian	Late Permian	1400	0	1000	3	50	Van Geet et al. (2011); Srozyk et al. (2014)
Ladonian Basin	Late Triassic-Early Jurassic	Late Triassic	Early Jurassic	1500	0	2000	5	50	Rasmussen et al. (1998); Alves et al. (2002); Penn dos Reis et al. (2017)
Nordkapp Basin	Late Permian	Late Devonian-Late Carboniferous	Late Permian	4000	800	2000	3	80	Jensen and Storenseter (1992); Koyi et al. (1995); Nilsen et al. (1996); Rowan and Lindoo (2017); Røp et al. (2017)
Norwegian-Finnish	Late Permian	Early Permian	Late Triassic-Early Jurassic	1300	300	3000	5	80	Vejthak (1990); Geil (1991); Petersen et al. (1992); Clausen and Pedersen (1999); Jackson and Lewis (2016)
Polish Basin	Late Permian	Late Permian-Early Triassic	Late Jurassic	1000	1000	4000	10	100	Ferrer et al. (2008); Jannes et al. (2009, 2010); Ferrer et al. (2012)
Red Sea Basin*	Mid Miocene	Oligocene	Early-Mid Triassic	2500	0	3500	7	220	Díazlez et al. (1995); Wagner et al. (2002); Krzywiec (2004); Krzywiec et al. (2017)
Sole Pit Trough*	Late Permian	Permian	Late Triassic	300	200	1500	5	20	Mitchell et al. (1992); Heaton (1995); Augustin et al. (2021)
Wiking Graben	Late Permian	Early-Late Permian	Late Triassic	1500	600	500	2.5	200	Wilker and Cooper (1987); Hughes and Davison (1993); Stewart and Coward (1995)
Western Approaches Basin	Late Permian	Early-Mid Triassic	Mid Jurassic	1500	0	6000	7	60	Thomas and Coward (1996); Jackson and Larsen (2009); Jackson et al. (2019)
Average	-	-	-	~1600	~400	~2500	~5	~100	Chapman (1989)
Scaled mode parameters	-	-	-	0.016	0.004	0.025	5	0.001	

605 ~~Cobbold, 1991; Weijermars et al., 1993; Bonini et al., 2012). In the presented models, we simulated the deformation processes of tectonically driven extension and gravity driven gliding and spreading and identified the following characteristic parameters describing the boundary conditions of these processes (cf. Fig. 3): layer thickness (h_b, h_d), densities (ρ_b, ρ_d), cohesion and coefficient of internal friction of the brittle overburden (C, μ), viscosity of the ductile layer (η), extension rate (u_e) and tilting rate (u_t). Extension and tilting rates depend on the geometrical and kinematical scaling and determine the amount of extension and tilting during experimental time. In the presented analogue models, we applied standard scaling procedures to relate geometrical, dynamical, kinematical and rheological model parameters to the natural prototype (Hubbert, 1937; Ramberg, 1981; Weijermars and Schmeling, 1986; Davy and Cobbold, 1991; Weijermars et al., 1993; Bonini et al., 2012). Geometric scaling ensures that corresponding ratios of lengths and angles are comparable between model and nature. The choice of the geometrical scaling ratio is determined by the procedure for dynamical scaling (Pollard et al., 2005) and, in our study, similar to previous model studies (e.g. Vendeville et al., 1995; Koyi et al., 1993; Warsitzka et al., 2015; Dooley and Hudec, 2020):~~

610 In the presented analogue models, we applied standard scaling procedures to relate geometrical, dynamical, kinematical and rheological model parameters to the natural prototype (Hubbert, 1937; Ramberg, 1981; Weijermars and Schmeling, 1986; Davy and Cobbold, 1991; Weijermars et al., 1993; Bonini et al., 2012). Geometric scaling ensures that corresponding ratios of lengths and angles are comparable between model and nature. The choice of the geometrical scaling ratio is determined by the procedure for dynamical scaling (Pollard et al., 2005) and, in our study, similar to previous model studies (e.g. Vendeville et al., 1995; Koyi et al., 1993; Warsitzka et al., 2015; Dooley and Hudec, 2020):

$$\frac{L_m}{L_n} = 10^{-5} \quad (\text{B1})$$

This means that 1 cm in the model represents 1 km in nature and characteristic geometrical relationships, such as the thickness ratio between overburden and salt or the ratio between extensional displacement and flank uplift can be scaled accordingly (Tab. A1).

620 Fundamental principles of dynamical scaling require that trajectories and ratios of forces acting on the material have to be equal and that rheological behaviours of the involved materials are similar (Weijermars and Schmeling, 1986; Pollard et al., 2005). This can be translated to the concept that driving and resisting forces of the relevant processes have to be related properly. In salt tectonics, driving forces are produced by (1) vertical loading of sediments resulting in a gravitational buoyancy forces, and by (2) tectonic stresses resulting in lateral pressure forces (Jackson and Hudec, 2017). Resisting forces are caused by (1) the frictional strength of the brittle overburden (S_F), and by (2) viscous stresses in the ductile layer substratum (S_V) caused by drag at its boundaries (Jackson and Hudec, 2017). Inertial forces can be considered as insignificant, since strain rates during solid rock flow are very low (Ramberg, 1981; Weijermars and Schmeling, 1986).

625

Because deformation style and localization in the materials is mostly influenced by rheological parameters of the ductile viscous layer and its brittle overburden, a characteristic measure for dynamical scaling is the brittle-to-ductile viscous strength ratio or 'Brittle- ductile viscous-coupling' (BDCBVC) (e.g. Sornette et al., 1993; Schueller and Davy, 2008). This is calculated by integrating the differential stress over the layer thicknesses (Schueller and Davy, 2008):

630

$$BVC = \frac{F_F}{F_V} = \frac{\int_0^{h_b} (\sigma_1 - \sigma_3) dz}{\int_0^{h_v} (\sigma_1 - \sigma_3) dz} \quad (\text{B2})$$

The frictional strength of both, brittle sediments and experimental granular materials, obeys the Mohr-Coulomb criterion according to which the mean stress σ_m and the maximum shear stress τ_{max} are related by (Byerlee, 1978; Dahlen, 1990; 635 Jaeger et al., 2007):

$$\tau_{max} = \sigma_m \sin\phi + C \cos\phi \quad (\text{B3})$$

Here, ϕ is the angle of internal friction and C the cohesion. Translated to an expression by the maximum and the minimum principal stresses σ_1 and σ_3 gives:

$$\frac{\sigma_1 - \sigma_3}{2} = \frac{\sigma_1 + \sigma_3}{2} \sin\phi + C \cos\phi \quad (\text{B4})$$

640 In an extensional stress field, the maximal principal stress σ_1 equals the lithostatic pressure σ_L :

$$\sigma_1 = \sigma_{zz} = \rho_b g h_b \quad (\text{B5})$$

where g is the gravitational acceleration and ρ_b the density of the brittle overburden. Furthermore, the minimal principal stress σ_3 is horizontal:

$$\sigma_3 = \sigma_{xx} \quad (\text{B6})$$

645 With the presence of a pore fluid pressure, each pressure has to be corrected by the pore fluid pressure ratio (Hubbert-Rubey coefficient of fluid pressure λ) (Weijermars et al., 1993):

$$\sigma' = \sigma(1 - \lambda) \quad (\text{B7})$$

Assuming that pore fluid pressure is hydrostatic, λ becomes roughly 0.45 for an average overburden density of 2200 kg m^{-3} and a fluid density of 1050 kg m^{-3} . The critical horizontal stress at failure (Dahlen, 1990):

$$650 \quad \sigma_{xx} = \sigma'_{zz} \frac{1 - \sin\phi}{1 + \sin\phi} - 2C \sqrt{\frac{1 - \sin\phi}{1 + \sin\phi}} \quad (\text{B8})$$

Combining Eq. (B4), (B5) and (B8) gives the brittle strength in a tensional stress regime: (Weijermars et al., 1993):

$$S_F = \rho_b g h_b^2 (1 - \lambda) - \rho_b g h_b^2 (1 - \lambda) \frac{1 - \sin\phi}{1 + \sin\phi} + 2C \sqrt{\frac{1 - \sin\phi}{1 + \sin\phi}} \quad (\text{B9})$$

The force required for extensional failure is then derived by integrating S_F over the thickness h_b (Schueller and Davy, 2008):

$$F_F = \int_{h_b} S_F dz = \frac{1}{2} \rho_b g h_b^2 (1 - \lambda) - \frac{1}{2} \rho_b g h_b^2 (1 - \lambda) \frac{1 - \sin \phi}{1 + \sin \phi} + 2C h_b \sqrt{\frac{1 - \sin \phi}{1 + \sin \phi}} \quad (\text{B10})$$

655 The viscous behaviour of a salt layer is characterized by a stress-strain rate relationship. Rock salt (halite), which usually constitutes the main volumetric proportion of a mobilized evaporitic succession (Jackson and Hudec, 2017), deforms by a combination of dislocation and diffusion (solution-precipitation) creep mechanisms depending on grain size, temperature, strain rate, percentage of intercrystalline fluids and confining pressure (e.g. Schléder et al., 2008; Urai et al., 2008). The large-scale deformation behaviour is suggested to be an average of these creep laws (Weijermars and Schmeling, 1986). At relatively
660 low strain rates ($< \sim 10^{-13} \text{ s}^{-1}$), low confining pressure ($< 10^6 \text{ Pa}$), grain sizes of $< 1 \text{ cm}$ and typical temperature range (20–200°C), rock salt deformation is dominated by diffusion creep, which is characterized by a linear stress-strain rate relationship (Newtonian fluid) (Van Keken et al., 1993). At higher strain rates and confining pressure, the deformation is a combination of diffusion and dislocation creep so that the shear stress-strain strain rate relationship is logarithmic (power-law fluid) and salt shows a strain rate softening behaviour (Urai et al., 2008; Van Keken et al., 1993). Silicone putty, which is commonly
665 used in salt tectonic experiments, possesses a strain rate softening behaviour for strain rates $> 10^{-2} \text{ s}^{-1}$ (Rudolf et al., 2016). However, such high strain rates are not reached in the presented models. Thus, for modelling purposes, the salt layer is here approximated as temperature-independent, mechanically uniform ~~ductile~~viscous material with a linear (Newtonian) stress-strain rate behaviour This condition might be true for wet salt ($> \sim 0.05 \text{ wt\%}$ water content) with relatively small grain sizes ($< 0.5 \text{ mm}$) (Urai et al., 2008). Under this assumption, differential (viscous) stress is calculated by (Weijermars et al., 1993;
670 Bonini et al., 2012; Turcotte and Schubert, 2014):

$$S_V = \sigma_1 - \sigma_3 = \eta \dot{\gamma} \quad (\text{B11})$$

Here, η is the dynamic viscosity and $\dot{\gamma}$ the shear strain rate in x-direction, ~~which can be approximated by the shear velocity u_e at the base of ductile layer and the thickness of the viscous layer h_d . Tectonic extension induces a Couette shear flow (Turcotte and Schubert, 2014) at the base of the silicone on the flanks.~~ Integrating Eq. (B11) over the thickness of ~~the ductile layer h_d~~ the viscous layer h_v gives the viscous force:
675

$$F_V = \int_{h_v} S_V dz = \eta \dot{\gamma} h_v \quad (\text{B12})$$

Tectonic extension induces a Couette shear flow (Turcotte and Schubert, 2014) at the base of the silicone on the flanks. Hence, $\dot{\gamma}$ can be approximated by the ratio of the shear velocity v_x at the base of viscous layer and the thickness of the viscous layer h_v , so that Eq. (B12) becomes:

$$680 \quad F_V = \eta v_x \quad (\text{B13})$$

For comparison between nature and model, we inserted a range of natural and experimental values (Tab. B1) in Eq. (B10) and (B13) ~~and calculated BDC (Eq. (B2))~~. To achieve correct dynamical scaling, ~~BDC/BVC~~ should be equal for model and nature. However, natural values of viscosity ($10^{16} - 10^{19}$ Pas; Van Keken et al. (1993); Mukherjee et al. (2010)) and strain rates ($10^{-16} - 10^{-14}$ s $^{-1}$; Jackson and Hudec (2017)) can be highly variable so that F_V can vary over several orders of magnitude (Tab. B1). Nevertheless, the curves in Fig. B1a show that ~~BDC/BVC~~ for model and nature are relatively similar for a wide range of values, ~~especially for intermediate values~~.

Equation (B10) reveals that the thickness of the overburden h_b mainly governs the coupling between brittle and ductile layer due to its quadratic input. Frictional forces are 1–2 orders of magnitude higher than viscous forces, at least at the beginning of the experiment when strain in the viscous layer is distributed widely. During progressive strain localization in later stage of the experiment, higher strain rates and, hence, higher viscous forces arise. Therefore, the BDC can be a measure for strain localization (Sornette et al., 1993; Schueller and Davy, 2008) in future advanced experiments when strain evolution is quantified.

The ratio between brittle and ~~ductile/viscous~~ material properties sets the velocity scaling (~~extension rate, tilting rate, sedimentation rate~~) by relating viscosity of the ~~ductile layers/substratum~~, strain rates and characteristic lengths, e.g. the thickness h_b , between model and nature (Weijermars et al., 1993; Allen and Beaumont, 2012):

$$\frac{u_m}{u_n} = \frac{\left(\frac{\rho_b g h_b^2}{\eta}\right)_m}{\left(\frac{\rho_b g h_b^2}{\eta}\right)_n} \quad (\text{B14})$$

Applying average values of viscosity, thickness and density (Tab. B1), the velocity scaling ratio is $\sim 10^4$. ~~Natural extension rate can vary between 0.1 and 10 mm a $^{-1}$ so that experimental extension rate should be in the range of ~ 0.1 to 10 mm h $^{-1}$. We choose intermediate rates 1 mm h $^{-1}$ for practical reasons so that a simulation duration of several hours to a few days is achieved.~~

Tectonic stresses and gravity are the main driving factors in the geological scenarios modelled here. For scaling purposes we used the 'Argand number' (Ar) ~~from lithospheric scale models of lithospheric brittle-ductile systems~~ (e.g. England and McKenzie, 1982; Schueller and Davy, 2008), which relates tectonic forces to the buoyancy force F_B . The tectonic force is approximated by the integrated strength of the brittle and the ~~ductile/viscous~~ layer and, therefore, the sum of F_F (Eq. (B10)) and F_V (Eq. (B13)). Thus, the Argand number is calculated by:

$$Ar = \frac{F_B}{F_F + F_V} \quad (\text{B15})$$

The buoyancy force F_B results from isostatic sinking of the overburden into the ductile viscous substratum, which is calculated by the weight of the overburden depending on thickness and density and the weight of the equivalent volume of the viscous substratum (e.g. Zoback and Mooney, 2003; Schueller and Davy, 2008):

$$710 \quad F_B = \int_{\Delta h_v} \rho_v g z dz - \int_{h_b} \rho_b g z dz \quad (\text{B16})$$

where $\Delta h_a \Delta h_v$ is the amount of subsidence of the overburden into the ductile viscous layer. Under isostatic conditions, $\Delta h_a \Delta h_v$ is:

$$\Delta h_v = \frac{\rho_b h_b}{\rho_v} \quad (\text{B17})$$

After integration, Eq. (B16) becomes:

$$715 \quad F_B = \frac{1}{2} g \rho_b h_b^2 \left(\frac{\rho_b}{\rho_v} - 1 \right) \quad (\text{B18})$$

~~The gravitational buoyancy force is mainly influenced by the ratio between overburden and salt density. Salt density is constant with depth, whereas the density of other sediments increases with depth due to compaction. During early stage of the basin history, the overburden density ρ_b is smaller than the salt density ρ_v so that F_B is negative. It changes to positive values when ρ_b exceeds ρ_v . Thus, Ar also changes from negative to positive when inserting a range of potential values (Tab. B1) into Eq. (7) (Fig. B1b). The diagram in Fig. B1b shows that the curve shapes of Ar are similar between model and nature meaning the general behaviour of the influence of buoyancy force, which changes the dynamics of flow in the viscous layer, is adequately approximated in the models.~~

When inserting a range of possible values (Tab. B1) into Eq. (B15), the curve shapes of Ar (Fig. B1b) are similar between model and nature meaning the general behaviour of the influence of buoyancy force is adequately approximated in the models.

725 **Appendix C: ~~Density stratification~~ Density stratification in model and nature**

~~Whereas the average density of a salt layer is assumed to be nearly constant with depth ($\sim 2150 \text{ kg m}^{-3}$), other~~ The average density of a salt layer is nearly constant with depth ($\sim 2150 \text{ kg m}^{-3}$). In contrast, sedimentary rocks consolidate during burial (Fig. C1a and b). Due to chemical mineralization and mechanical compaction the porosity ϕ of these sediments decreases with depth leading to an increasing density. The most common porosity-depth-relationship is given by Athy's law (Athy, 1930):

$$730 \quad \phi = \phi_0 e^{-cz} \quad (\text{C1})$$

Table B1. List of parameters and values used for the scaling procedure, i.e. to calculate BVC (Eq. (B2)) and Ar (Eq. (B15)).

Parameter	Sign	Dimension	Model	Nature	Scaling factor
Thickness (brittle layer)	h_b	[m]	0.001–0.03	100–3000	10^{-5}
Thickness (<u>ductile</u> <u>viscous</u> layer)	h_v	[m]	0.005–0.03	500–3000	10^{-5}
Extension rate	u_x	[mm h ⁻¹] / [mm a ⁻¹]	0.1–10	0.1–10	$\sim 10^4$
Strain rate	$\dot{\epsilon}$	[s ⁻¹]	10^{-6} – 10^{-4}	10^{-16} – 10^{-13}	10^{-6} – 10^{-10}
Density (brittle layer)	ρ_b	[kg m ⁻³]	800–1050	1850–2300	~ 0.45
Density (<u>ductile</u> <u>viscous</u> layer)	ρ_v	[kg m ⁻³]	970	2150	~ 0.45
Coefficient of internal peak friction	μ	–	0.55–0.65	0.5–0.8	~ 1
Cohesion	C	[Pa]	20–60	$1-10 \times 10^6$	0.5×10^{-6} – 2×10^{-5}
Dynamic viscosity	η	[Pa s]	2.2×10^4 – 2.8×10^4	10^{17} – 10^{19}	10^{-13} – 10^{-15}
Lithostatic pressure	σ_L	[Pa]	10–300	$5-40 \times 10^6$	$\sim 7 \times 10^{-6}$
Brittle strength	S_F	[Pa]	30–300	$2-40 \times 10^6$	0.7×10^{-5} – 1.5×10^{-5}
<u>Ductile strength</u> <u>Viscous stress</u>	S_V	[Pa]	0.006–0.08	30–30000	2×10^{-6} – 2×10^{-4}
Frictional force	F_F	[N m ⁻¹]	0.03–5	10^8 – 10^{11}	7×10^{-11} – 16×10^{-11}
Viscous force	F_V	[N m ⁻¹]	10^{-4} – 10^{-2}	10^5 – 10^9	10^{-9} – 10^{-11}
Buoyancy force	F_B	[N m ⁻¹]	-10^{-3} – 10^{-1}	-10^7 – 10^{10}	$\sim 5 \times 10^{-11}$
Brittle- <u>ductile</u> <u>viscous</u> coupling	BVC	–	10^{-1} – 10^3	10^{-2} – 10^5	0.04–7
Argand number	Ar	–	-0.07–0.07	-0.11–0.09	0.3–1.5

where ϕ_0 and c stand for the initial porosity and the depth coefficient. The corresponding density $\rho_b(z)$ at each depth can then be calculated by using the grain density ρ_G (Fig. C1b) and the density of the pore fluid ρ_f :

$$\rho_b(z) = \phi \rho_f + (1 - \phi) \rho_G \quad (C2)$$

Integrating over the entire thickness of the cover h_b gives the average density $\bar{\rho}_b$ of the cover:

$$\bar{\rho}_b = \rho_G + \frac{(\rho_G - \rho_f) \phi_0 (e^{-ch_b} - 1)}{c h_b} \quad (C3)$$

Because of this density increase, the bulk density of overburden sediments $\bar{\rho}_b$ is smaller than that of salt at shallow depths and exceeds the average salt density at a depth between 2200 and 3800 m (e.g. Jackson and Talbot, 1986; Hudec et al., 2009) (Fig. C1b). Exceeding this thickness, the overburden is able to autonomously subside into the salt layer and squeeze the salt upwards (Warsitzka et al., 2018).

740 Compaction of natural sediments cannot be directly simulated in analogue models using the common granular materials. However, it is not the density increase itself, but the increase of the lithostatic pressure at the base of the overburden that is the crucial parameter determining differential loading and the pressure gradient in the ductileviscous layer. In order to simulate an appropriately scaled increase of the lithostatic pressure, we increased the density of the sand mixture (ρ_{bmi}) during each step

of syn-kinematic sand accumulation, similar to the approach by Dooley and Hudec (2020). This results in an inverse density stratification with less dense material at the bottom and denser material at the top of the cover (Fig. C1c) and a stepwise increase of the bulk density of the entire sand cover. To compare natural with experimental density-depth curves, the standardized density ρ' was calculated by dividing all density values through the density of the salt or the silicone, respectively. Likewise, the depth was standardized by dividing the depth through the maximal thickness of the cover (here: 5 cm in the model and 5 km in nature).

The increase in density of each new sand layer has to be set in such a manner that the standardized bulk density resembles the exponential increase of the standardized bulk density in nature (Fig. C1d). During the experiment, we first estimated the maximum thickness of the next syn-kinematic sand layer, which roughly equals the amount of vertical subsidence of the central graben block. Then, the standardized density of this layer, which depends on the total thickness of the sand column in the graben centre, is given by the correlation shown in Fig. C1d. Based on that, the bulk density of the next syn-kinematic layer can be calculated and the appropriate mixture of quartz sand and silicate cenospheres can be produced. In advance, we determined bulk densities of various mixing ratios. The ratios in exp. ETSv ranged from 3:8 to 3:2 (quartz sand : silicate cenospheres in volumetric fractions), which produced bulk densities of 850 to 1160 kg m⁻³.

Author contributions. MW invented the concept of this modelling study, designed the experimental apparatus, performed the experiments and composed main parts of the manuscript. Furthermore, data processing and analysis was undertaken by MW. PZ advised the designing and construction of the apparatus and assisted during the experiments. Theoretical input for planing the project, the model design and the concept of the experiments came from FJK. PK supported the comparison of the experiments with the natural case study of the Polish Basin. FJK, PK and PZ helped with finalizing the manuscript.

Competing interests. The authors declare that they have no conflict of interest.

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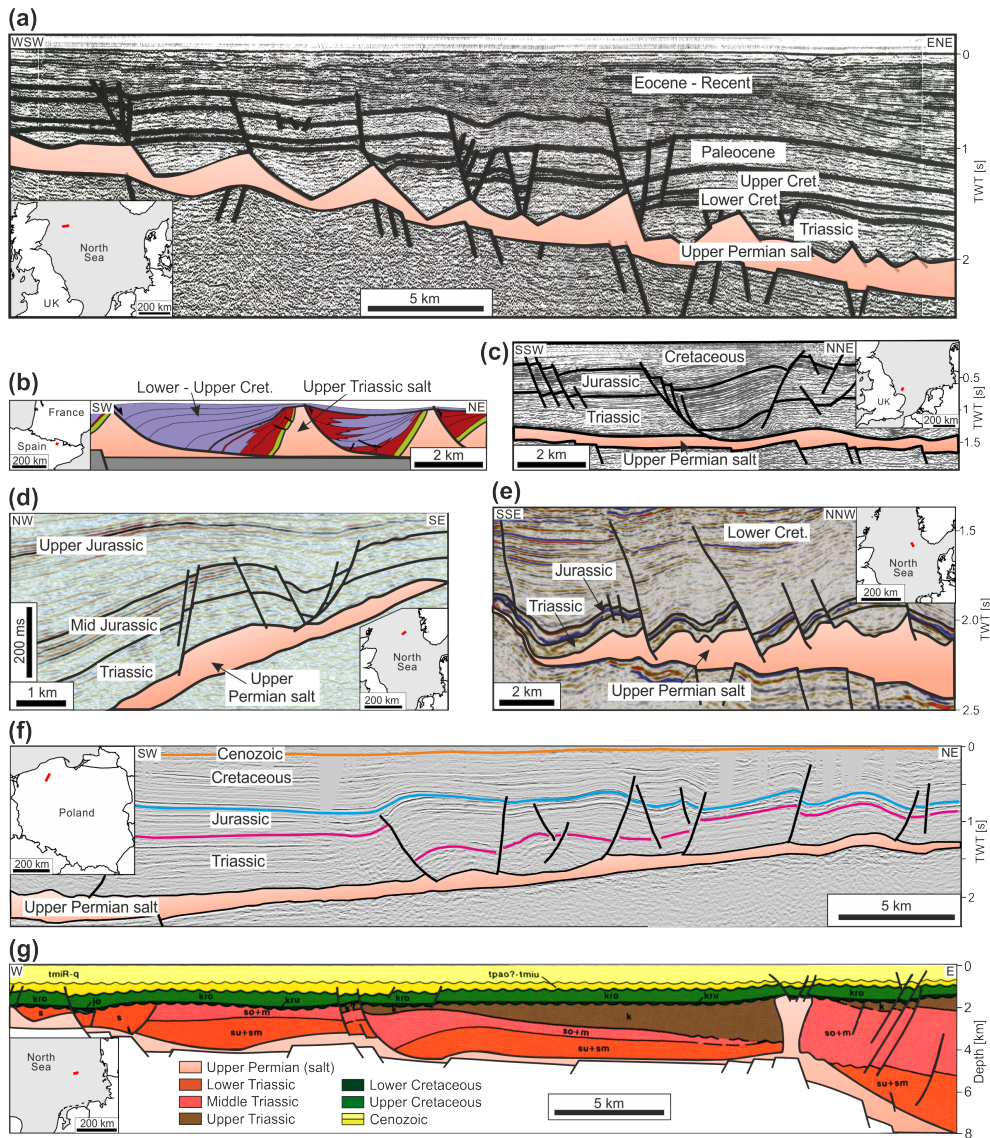


Figure 13. Examples of thin-skinned extensional structures recognized above inclined flanks at the margins of salt-bearing rift basins. **(a)** North-western Central Graben (central North Sea): Evenly spaced listric growth faults above a gently dipping, faulted sub-salt base (modified from Bishop, 1996). **(b)** Pyrenean rift system (northern Spain): Sketched cross section based on geological field data showing thin-skinned listric growth faults and reactive diapirs of Upper Triassic salt (modified from Lopez-Mir et al., 2014). **(c)** Sole Pit Basin (southwestern North Sea): Thin-skinned graben structure ('Dowsing Graben System') above the Permian salt detachment was partly attributed to gravity gliding into the bowl-shaped basin centre (modified from Stewart et al., 1996). **(d)** South Viking Graben (northern North Sea): Supra-salt fault zone at the margin of a tilted half graben (modified from Jackson and Larsen, 2009). **(e)** Egersund Basin ('Sele High Fault System'; central North Sea): Supra-salt fault zones above a slightly faulted salt base (modified from Jackson and Lewis, 2016). **(f)** North-eastern Polish Basin: Peripheral extensional fault zones characterized by several supra-salt, tilted blocks and listric faults with increased thickness of Jurassic sediments. During the Late Cretaceous, the structure was inverted, and this inversion tectonic was associated with thin-skinned reverse faulting and folding (modified from Krzywiec, 2006b, 2012). **(g)** Horn Graben (south-eastern North Sea): Listric growth faults in the Triassic strata overlie a relatively flat sub-salt base and are detached from the major graben structure (modified from Baldschuhn et al., 2001).

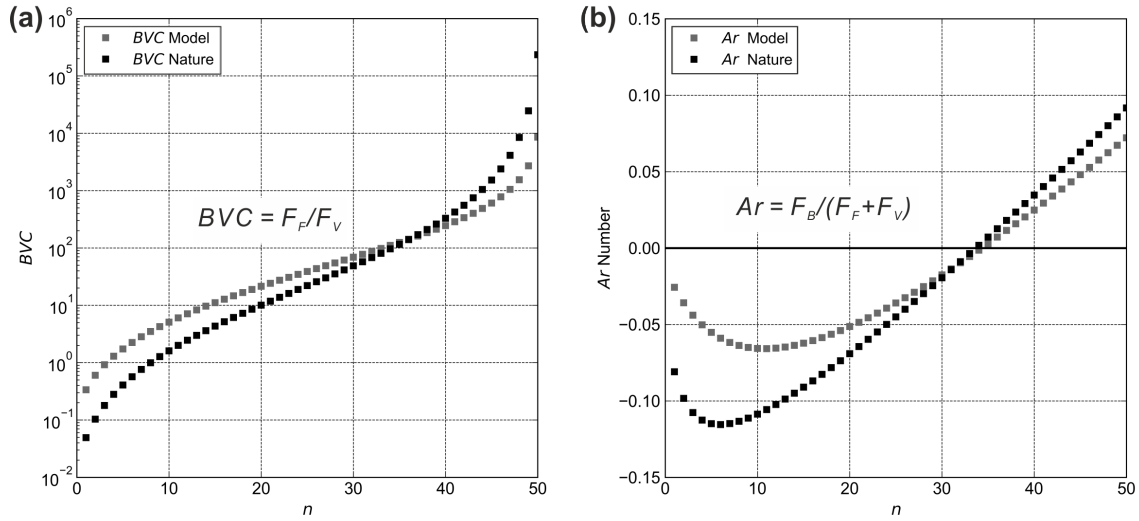


Figure B1. Diagrams of the (a) brittle-~~ductile~~viscous-coupling (BVC) and (b) the Argand number (Ar) for natural and experimental values. To calculate the values of both parameters, a range of n ($= 50$) values of the input parameters (Tab. B1) arrays of each input parameter were inserted in Eq. (B2) and Eq. (B15). For each parameter, the array contains 50 equidistant values ranging between the minimum and maximum values given in Tab. B1). Ar - Argand number, BVC - brittle-~~ductile~~viscous coupling, F_B - buoyancy force, F_F - frictional force in the brittle layer, F_V - viscous force in the ~~ductile layer~~substratum.

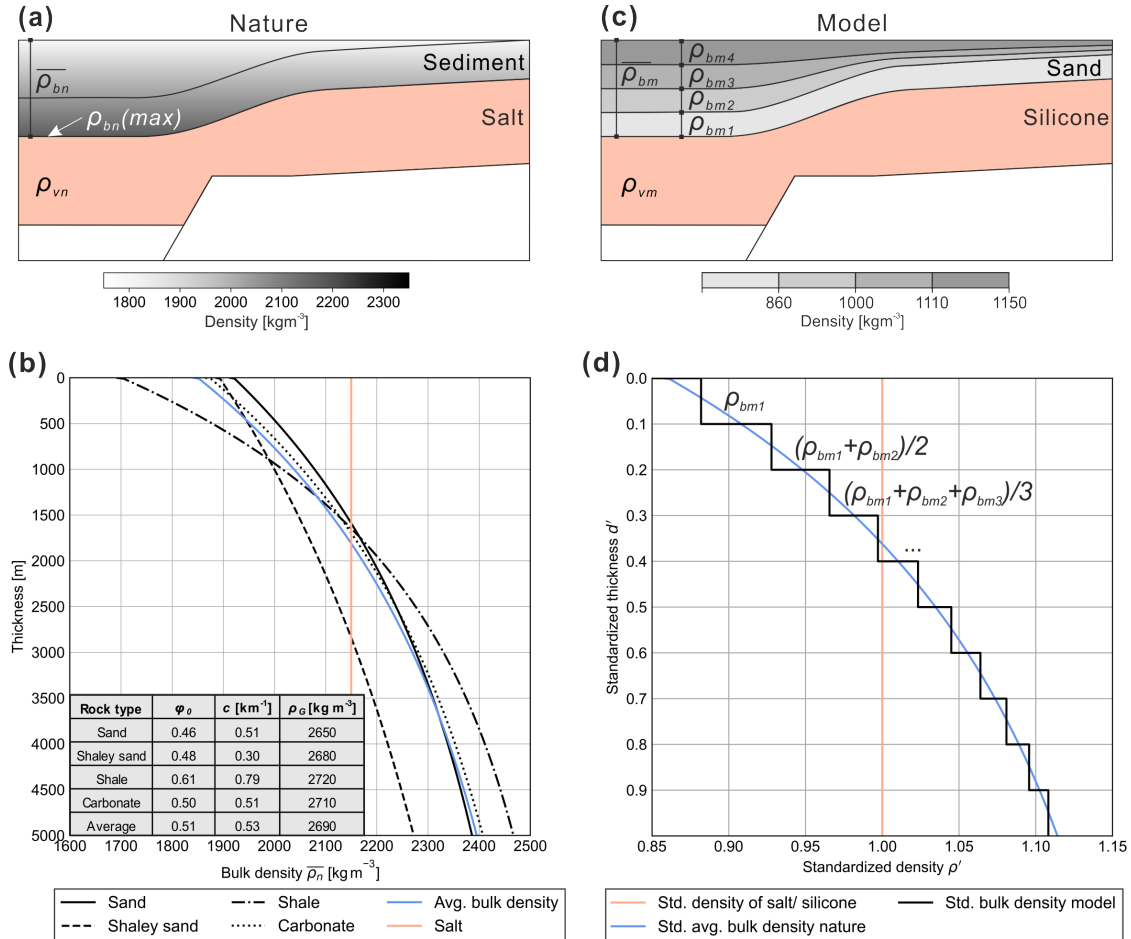


Figure C1. Density configuration in the natural prototype and in the analogue model. (a) Natural sediments compact during burial leading to an increasing density ρ_{bn} with depth according to Athy's law (Athy, 1930). (b) Curves of the bulk density $\bar{\rho}_n$ for different sediment types related to the thickness of the overburden. Due to physical compaction according to Athy's law (Athy, 1930), the density increases exponentially with depth. The blue curve shows the bulk density calculated by using average (avg.) values listed in the table. (c) Inverse density stratification in the analogue model presented here. Layers with increased density are sieved in leading to an increasing bulk density during subsidence. (d) Standardized (Std.) thickness and density values for nature and model. Thickness was standardized by the (presumed) maximum thickness of the overburden (nature: 5 km, model: 5 cm). Density values were standardized by dividing the density of the cover (nature: ρ_{bn} , model: ρ_{bm}) by the density of the viscous layer (nature: ρ_{vn} , model: ρ_{vm}). Note that the density of each sieved layer in the models increases upwards. This procedure is applied to increase the bulk density of the subsiding overburden in a similar manner as in nature. ϕ_0 – initial porosity, c – depth coefficient, ρ_G – grain density.