

Dear referee,

We are glad to receive the review report and would like to express our sincere thanks to you and the editors. Without the constructive comments from you, the quality of this manuscript cannot be significantly improved. All the comments and suggestions have been carefully considered to revise the manuscript. Detailed reply to all comments and the associate manuscript modifications are given below.

Reply on RC2

1. Please revise the language using a professional editor or native speaker. There are many areas that may cause confusion as they are currently written.

Reply: *We appreciate your advice and have revised the manuscript carefully to improve the language, with the help of all the co-authors.*

2. Definitions: one of the major problems with this work at the moment is the lack of a definition of "ridge suction". Plate drag is reasonably well explained as the frictional force imposed upon the sub-lithospheric plume material, but ridge suction seems to be simply anything that causes plume material to travel toward the ridge. Currently, I have to infer this definition since no clear description is given and the quantitative assessment of ridge suction is a fractional number looking at the volume of plume material flowing toward and away from the ridge. If ridge suction is only assessed in this way, then this is inconsistent with the literature and should be retermed as ridgeward flow or something similar. Better, I believe the authors need to reassess their model results with a more consistent definition of ridge suction that can be quantitatively assessed.

Reply:

Thanks for your suggestions. The concept of "ridge suction" refers to Niu (2004), who suggest that the spreading ridge sucks the material from depths due to the pressure gradient between the ridge center and deep hot material. The buoyant

mantle plume from deep is overpressured, while the ridge center is in the state of underpressure. Therefore, the ridge suction we termed in the manuscript indicates the dynamic pressure gradient between the plume and ridge, which drives the plume material flowing to the ridge. In any cases, we reworded most of the occurrence of “ridge suction” in the manuscript, because it is only one of the driving mechanisms for ridge-ward plume flow.

On the other hand, the active gravitational spreading of plume also contributes to the ridge-ward plume flow. We discussed and compared the gravitational gradient and ridge-ward dynamic pressure gradient of these two mechanisms in the main text (see section 3.3)

3. In the abstract, the authors claim that plate drag has not been studied very much. However, I think this is perhaps an overly strong statement. There are several studies that incorporate the affect of plate motion (and the consequent drag) on plume spreading including Ribe et al. (1995); Ribe (1996); Ribe and Delattre (1998); Ito et al. (1997), Hall et al., 2003; 2004, etc. Each of these works (and others) incorporates the affects of plate drag on plume spreading in their calculations. I think the authors should be clear about what aspect of their work contributes something these other authors do not.

Reply: We appreciate this comment. We revised and improved the description of the motivation for our study and its novelty in the abstract and introduction.

4. Model Comments

4.1 In my opinion, for the scope of this work, the model used is overly complicated and in some respects inaccurate for a mid-ocean ridge setting. For example, the authors are examining the flow of mantle material beneath a lithosphere and have included a 1.5 km thick sediment layer across the model, but near ridge (especially fast ridges) there is little to no sediment. In fact, even along the slow spreading MAR, 1.5 km of sediment does not occur along the ridge axis and, indeed not for a reasonable distance away.

Reply: We appreciate this comment. In our models, we imposed an average 1.5 km sediment layer vertically following the typical oceanic lithospheric structures. Horizontally, we use a uniform thickness lithosphere. We agree to the reviewer that the sediment near the ridge is negligible. Indeed, the processes of erosion/sedimentation are also considered in our simulations. We plot the thickness of sediment in the reference models (Figs. S1, S4). We are confident that the additional sediment layer thickness does not affect the conclusions of our study.

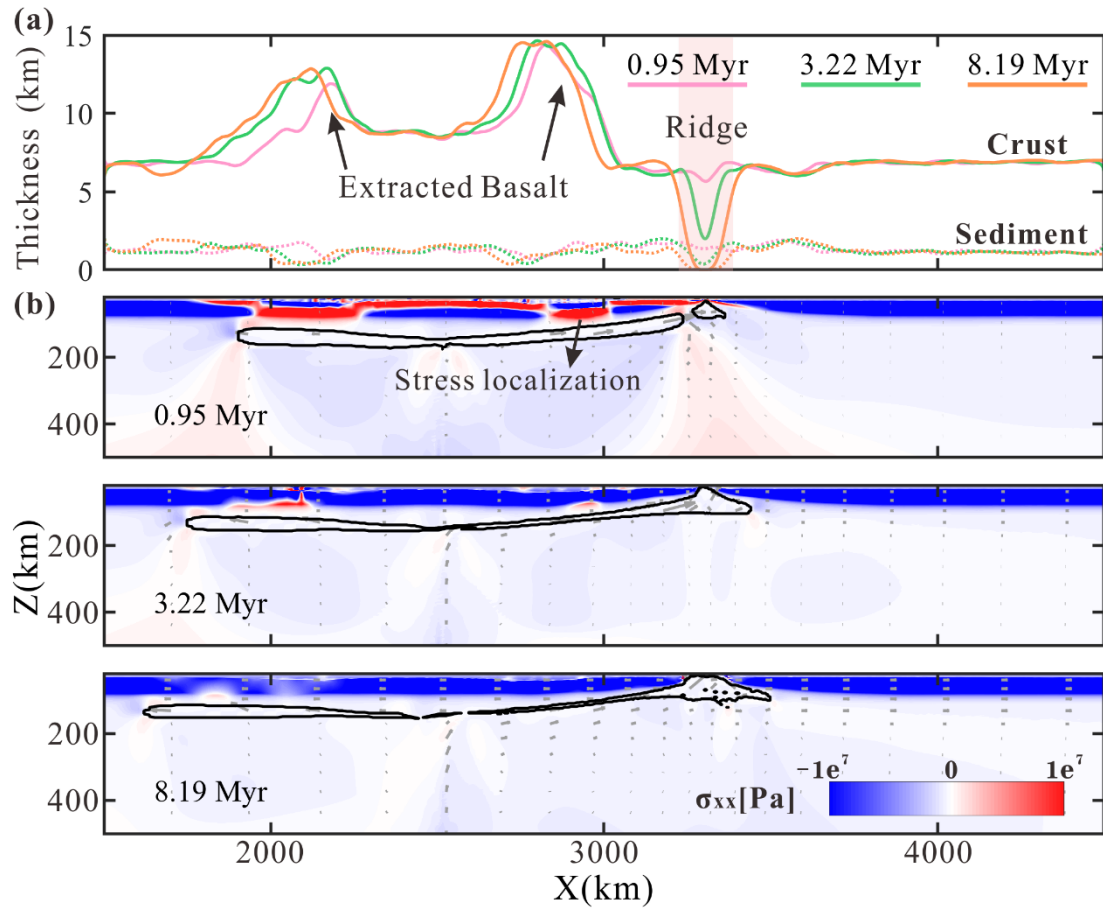


Figure S1. Reference model (M12, see Table S1, same as Fig. 3) evolution of ridge-ward plume flow shown by (a) crust and sediment thickness, (b) normal stress. The mantle plume weakens the overlying oceanic plate and changes the stress state of the overlying oceanic plate. Molten plume material beneath the lithosphere is extracted to the crust.

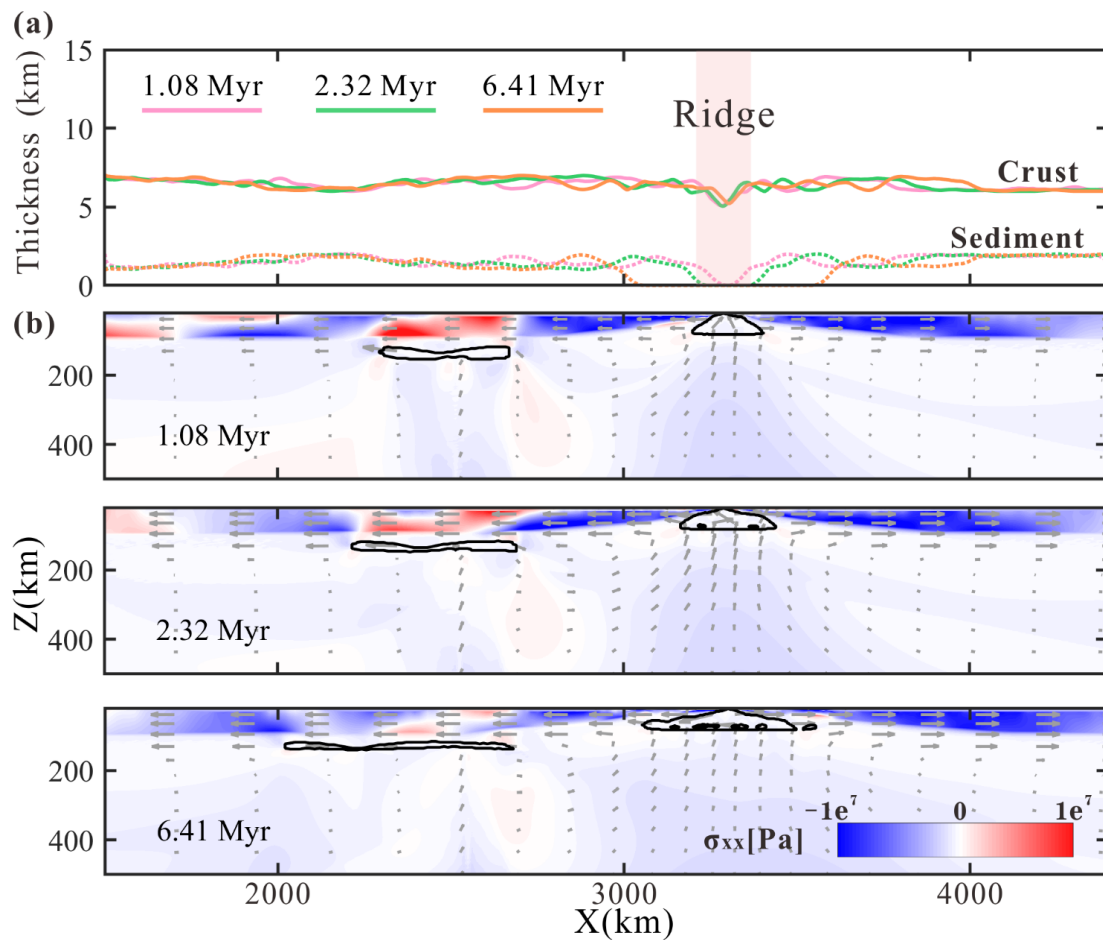


Figure S4. Reference model (M77, see Table S1, same as Fig. 4) evolution of trench-ward plume flow shown by (a) crust and sediment thickness, (b) normal stress. The mantle plume weakens the overlying oceanic plate and changes the stress state of the overlying oceanic plate. Molten plume material beneath the lithosphere is extracted to the crust.

4.2 Next, why is melting and heat flux useful for this study? Given the stated goal of the study to assess plume flow, I do not see (and it was not stated) why melting was useful or necessary. It is also unclear how the movement of melt throughout the system does or does not violate conservation of mass since you are working with an incompressible material and claiming to add material beneath the crust after removing it from another location. Please justify the use of melting and melt movement/accumulation. Also, clearly state any affect this melt has on your model (viscosity? temperature structure? density?, etc.)

Reply: *We appreciate your comments. In the model, melting of plume and asthenosphere are taken into account. The molten plume material is transformed to mafic magma and added to the crust, forming a thickened crust. We discussed the thickened crust between the plume and ridge to indicate the formation of oceanic plateau near the ridge during plume-ridge interaction. Besides, the melt extraction in the model is mass-conservative, as has been applied in previous modeling works (Gerya et al., 2015; Gülcher et al., 2020). We have modified the description of the melt extraction and its mechanism in the method section (lines 134-158). We also agree that the heat flux may not be that important for our study, so we remove the discussion of heat flux in the revised manuscript.*

4.3 Why do you need a plastic rheology? Given the scale of the problem you are working on, is the added focusing of the ridge axis to a smaller number of grid cells necessary? I don't believe that the current models can answer this question given the problem I mention next.

Reply: *Thanks for your comments. we employ viscoplastic rheology in the models. However, it is true that the stress and strain of ridge area usually are not strong enough to exceed the yielding criterion. The lithosphere does not exhibit plastic deformation in the model. Since plastic deformation does not occur in a widespread manner in our models, it indeed does not seem to be critical to be included in the model. On the other hand, including it for completeness (as we do here) does not change the main results of our study.*

Secondly, we added smaller grid cells (grid size decreases linearly from 20km at the edges to 2 km at the ridge axis). Indeed, we set denser grids in the middle model domain in order to better simulate the interaction between plume and ridge. Of course, considering the scale of our research problem, the size of model grid does not have a great influence on the results. According to initial test cases, using smaller grids helps to make the computation more stable and does not fail to converge due to numerical perturbations.

4.4. Another issue is the lack of adjustment of the lithosphere for plate spreading rate. In other words, the ridge and lithosphere in the “fast spreading, ridge drag dominated” cases do not appear to be in equilibrium before the plume is introduced. Looking at the compositional slices and temperature contours of the model in Figure 4, it appears that the sub-axial lithosphere is flattening out and a new, flatter lithosphere is forming without the initial half-space cooling structure (or with one that is in equilibrium with the faster spreading rate). This will alter the mantle flow field, the upslope topography of the ridge, and, potentially the location of the spreading ridge.

Reply: *We appreciate this comment and further checked our models. According to the model viscosity and temperature structure (Figs. S5, S6), we think that the flattening of the 1300 °C isotherm beneath the MOR is caused by the latent heat consumption during asthenospheric melting beneath the ridge. The effect of latent heating is not included in the initial temperature structure. Thus, while it is true that the 1300 °C isotherm flattens, the lithosphere does not flatten. It behaves as expected (see plots of the viscosity filed; Figs.S3, S6). Accordingly, we are confident that our predicted temperature and viscosity structure beneath the MOR is realistic.*

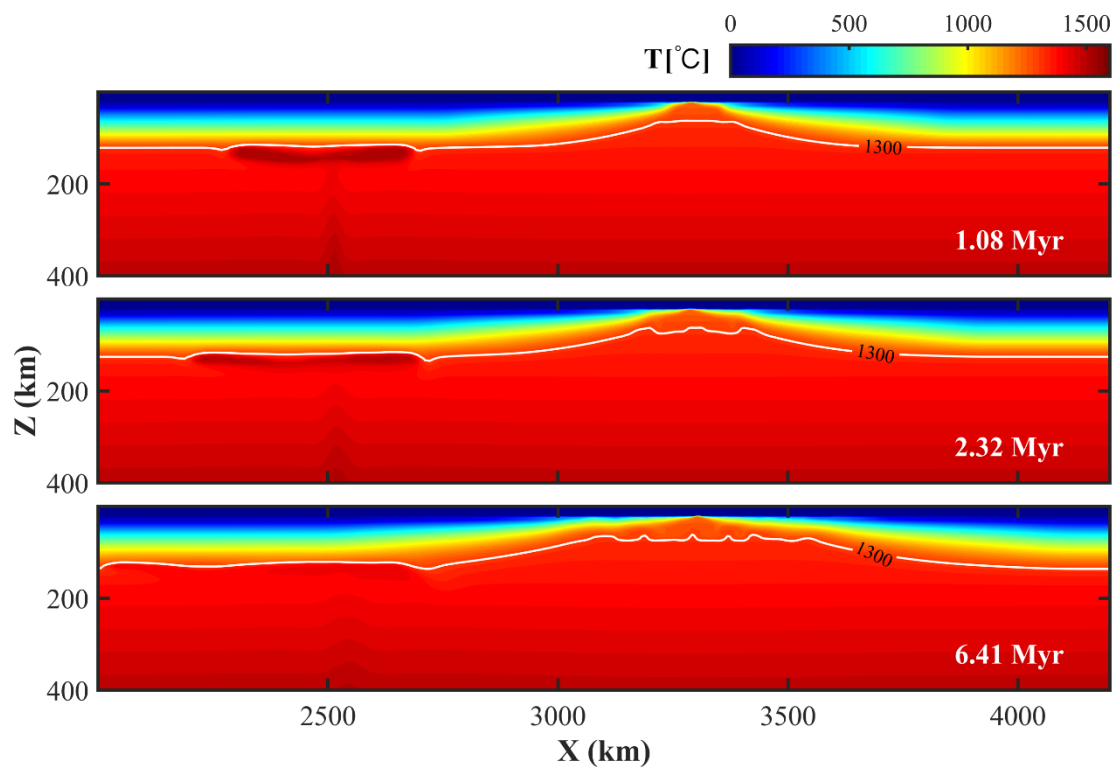


Figure S5. Temperature evolution of reference trench-ward plume flow model (M77, see Table S1, same as Fig. 4).

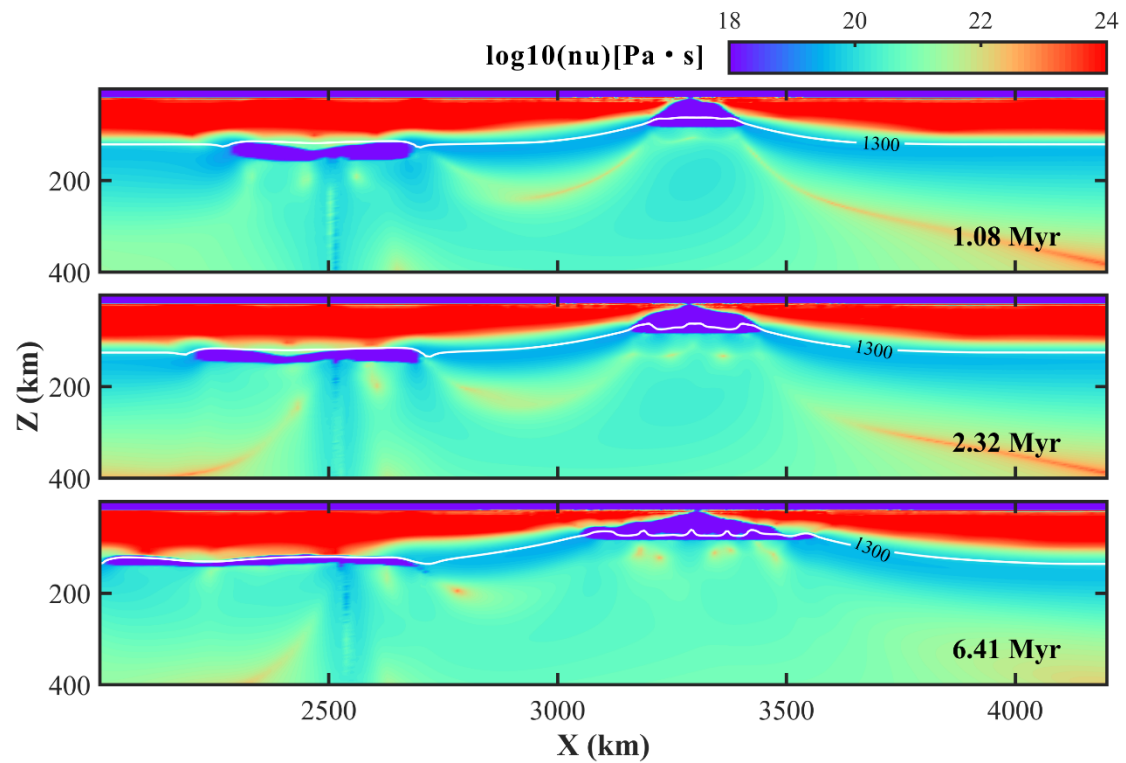


Figure S6. Viscosity evolution of reference trench-ward plume flow model (M77, see Table S1, same as Fig. 4).

4.5. Thermal structure

4.5.1 It is not clear to me how you arrive at a bottom boundary condition of 2513K when the base of the lithosphere has a $T_{\text{max}} = 1573$. Since the base of the lithosphere is at ~ 100 km depth (Figure 1) and there is an imposed 0.5 K/km adiabatic temperature gradient, the max temperature at 660 km depth should be $1573 \text{ K} + 560\text{km} \cdot 0.5\text{K/km} = 1853\text{K}$. This is a big discrepancy that might imply a much hotter mantle than is realistic, which would likely have significant impacts on the results of this study.

Reply: The model sizes in our study are set as 6600(width) and 1200(depth). The temperature of 2513K refers to the temperature at the bottom of model, that is, the temperature at 1200 km. Based on adiabatic temperature gradient of 0.5 K km^{-1} , the temperature at 660 km in all models is set initial to $1573 \text{ K} + (660-120) \text{ km} \cdot 0.5\text{K/km}$

=1843K.

4.5.2 How is the plume tail maintained? This is not clear or perhaps I missed it

Reply: *Thanks for your comments. The mantle plume imposed in our model is activate by given excess temperature at the beginning. Indeed, this temperature anomaly is only an initial condition, but not a long-term boundary condition. We do not impose a hot patch at the base of model, which means the mantle plume can uprise or erupt only once and does not create a long plume tail self-consistently. Considering that hotspots erupt periodically on earth, at intervals of millions of years. Therefore, in our study, we simplified the formation of mantle plume only considered one pulse of plume uprising. Such plumes are also widely used in similar studies (Baes et al., 2016; Gerya et al., 2015; Gülcher et al., 2020).*

5. Results/Interpretation

5.1 The images in Figure 4 demonstrate a factor that may explain the affect of plume head size on ridgeward flow – the erosion of the lithosphere by the plume head. As pointed out by Kincaid et al., 1995 in their laboratory experiments, the formation of lithospheric levees can act to block plume flow. This appears to be happening here. Small plume heads eat into the lithosphere a bit, effectively create ridges (or levees) that are the same thickness as the plume material and halt its motion. Then, as the plate moves the plume material has no choice but to flow with the plate. In contrast, large plume heads push the lithosphere out of the way all the way to the ridge. Despite the significant ridgeward flow, I would argue that this has nothing to do with ridge suction, but the lithosphere rheology and plume buoyancy forces.

Reply: *We appreciate this comment and give the following reply. As mentioned in our reply to comment 2, we suggest that the dynamic pressure gradient from the high-pressure plume to the low-pressure ridge actually contributes to the ridge-ward plume flow. The pressure gradient from the plume head to spreading center is notable (Fig. 5). Such pressure gradient varies with plume buoyancy force and different off-axis*

distances. An increase in the plume head size, indeed, enhances the dynamic pressure and promotes the plume flowing ridge-ward (Fig.S8a).

Primarily, the plume flows to the ridge owing to the pressure gradient at first, which decrease gradually when plume get closer to the ridge. Then, the gravitational spreading of plume starts to drive ridge-ward flow. As the reviewer points out, the gravitational gradient is related to topography at the base of the lithosphere. See section 3.3 for a detailed discussion.

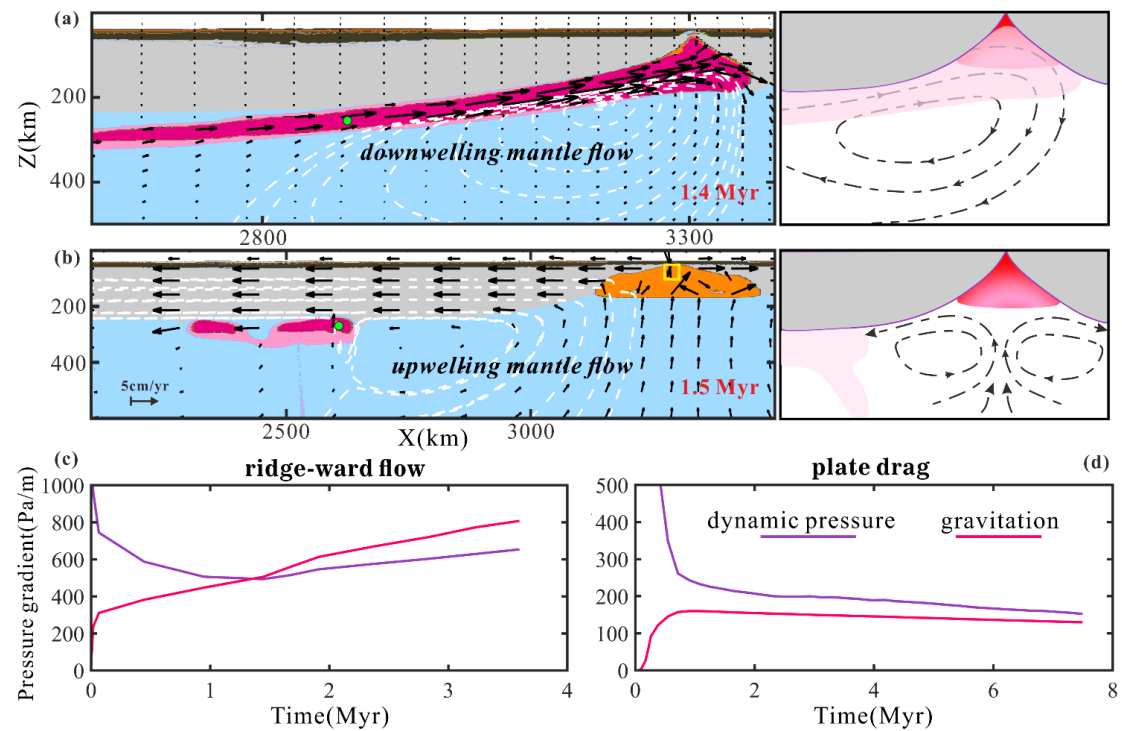


Figure 5. Comparison between models with ridge-ward vs. plate-drag plume flow. **(a)** Ridge-ward flow with downwelling beneath the MOR (results from case M12 as in Figure 3). White dashed lines are streamlines; black arrows visualize the flow field. Schematic of flow in the sub-panel on the right-hand side. **(b)** Plate-drag flow with upwelling mantle corner flow beneath the MOR (results from case M77 as in Figure 4). **(c)** The dynamic pressure and gravitational gradient of plume marker (i.e. green circle in (a)) over time. The yellow box in (b) marks the location for the computation of average dynamic pressure at the ridge, needed for the calculation of the dynamic

pressure gradient (see text). **(d)** The dynamic pressure and gravitational gradient of plume marker (i.e. green circle in (b)) over time.

5.2. The claim of “tension cracks” seems to be based on the stresses in the model. These stresses reach maximums of + or – 3×10^{-7} Pa (Figure 5), much too small to actually fracture of rock - especially near the surface, which typically has yield strengths many orders of magnitude larger. Is this a typo? If this should be + or – 3×10^7 Pa (i.e., 30 MPa) that seems very large and so I am left to question how tension cracks are justified here. However, I would note that I don't think these are essential for the results of this paper and fall into the over complication of the model for the state purpose of the modeling.

***Reply:** Thanks for your comments. Yes, the magnitude of normal stress here is a typo. The correct magnitude of these stresses should be $\pm 3 \times 10^7$ Pa in the models. Actually, the “tension cracks” used in the main text may be inaccurate. There are no brittle fractures in the lithosphere and normal faults near the surface in the models. While there are horizontal extensional stresses, there is in fact no yielding. As a result, we removed the description of the “cracks” in the main text.*

5.3. The role of ridge suction vs plate drag. I think the authors have glossed over some of the factors likely to contribute to the plume flow including the slope of the lithosphere and its role in guiding plume material up the slope, the buoyancy flux of the plume stem since this is not described by the plume radius definition here (which seems to describe the size of the plume head).

***Reply:** We appreciate this comment which triggers us a lot of thinking. We think the effect of the lithospheric slope is relative to the half spreading rate. The slope of the lithosphere varies with the spreading rate of the mid-ocean ridge. The base of the lithosphere would be flatter at the fast-spreading ridge. Consequently, a slow*

spreading ridge imposes smaller shear force on plume head and forms a steeper lithosphere base which benefits to the ridge-ward plume flow. Besides, we now evaluate the plume gravitational gradient, considering the local lithospheric slope, and ridge-ward dynamic pressure gradient (Fig. 5). For a detailed analysis, see section 3.3.

5.4. Much of the interpretation of these results hinge around spreading of the plume head, not the plume after it has established itself beneath the lithosphere. Many plume have been active for 10 Myr or more and the plume head will have greatly diminished or completely spread away by that time. Yet, these plume tails can still interact with ridges since ridges migrate and often approach plumes. How does the long term interaction look - after the plume head has disappeared?

Reply: *According to our results, the plume tail (which follows the plume head at model times 1~8 Myr) bends either toward or bend away from the ridge in ridge-ward flow and plate-drag flow dominated models, respectively. Taking the long-term evolution of the model, the degree of plume tail tilts towards or away from ridge increases with time, and ends up maintaining a relative steady state. Please also see our previous reply to 4.5.2.*

5.5. Related to 4., I don't think the authors should be claiming to assess plume radius, as this commonly is used to refer to the radius of the plume stem. Instead, I think the manuscript would be much clearer if the authors would state that they were varying the plume head radius.

Reply: *We appreciate this comment and made revisions. In this study, we varied the plume buoyancy flux by changing its radius. The plume radius in our model only refers to the initial size of mantle plume. The width of plume stem decreases dramatically when the plume rising to the plate and plume head spreading out, which is different from the columnar plume stem radius detected by the geophysical tomography. Thus, as you suggested, we use the “initial plume head radius” to replace “plume radius” in the main text.*

Other minor specific comments:

•Line 47-48 – I’m not clear as to what this statement has to do with the EPR sucking in plumes so that they do not appear near the ridge.

Reply: Thanks for your comments. We removed this sentence and reword this paragraph in introduction section.

•Line 52 – the use of the work “push” is inappropriate here and should be changed to “drag” or similar

Reply: Thanks for your advice. We replace the “push” with “convey” in the text (line 54-55).

•Line 58 - The authors should reconsider how they phrase things – for example, “slow spreading rate, short distance (small plume-ridge distances??), and large plume radii promote ridge suction,...” is an inaccurate statement – really, I think what the authors are trying to say is that these factors favor plumes being pulled toward ridges by ridge suction

Reply: Thanks for your advice. We rephrased all these inaccurate phrases in the main text.

•Line 59 – maybe try a more careful wording – it is the fast plate motions associated with fast-spreading ridges that exert strong drag forces on plumes

Reply: Thanks for your advice. We revised this paragraph to clarify our goal of this study in the main text.

Figure 2 – this does not look like a half-space cooling model. Is this a plate cooling model or some modified half-space model? The half-space cooling model does not flatten like this.

Reply: We appreciate this comment and made revisions in figure 2 and the main text. Indeed, the initial temperature distribution of the oceanic plate is prescribed by the half-space cooling model and thermal equilibrium structure. The half-space cooling

model is used to describe the oceanic plate younger than 50Myr, and the thermal equilibrium structure is used to describe older oceanic parts. We set up the models in this way because we consider that the theoretical half-space cooling model has a good match with some geophysical observations when the plate is young, but the fit becomes poor when the age is greater than 60/70 Ma (Turcotte and Schubert, 2014; Stein and Stein, 1994). Therefore, we set a half-space cooling model with a maximum age at 50Ma, and the thermal equilibrium thickness of the older lithosphere is constant (i.e., ~ 100 km; corresponding to a thermal age of 50 Ma).

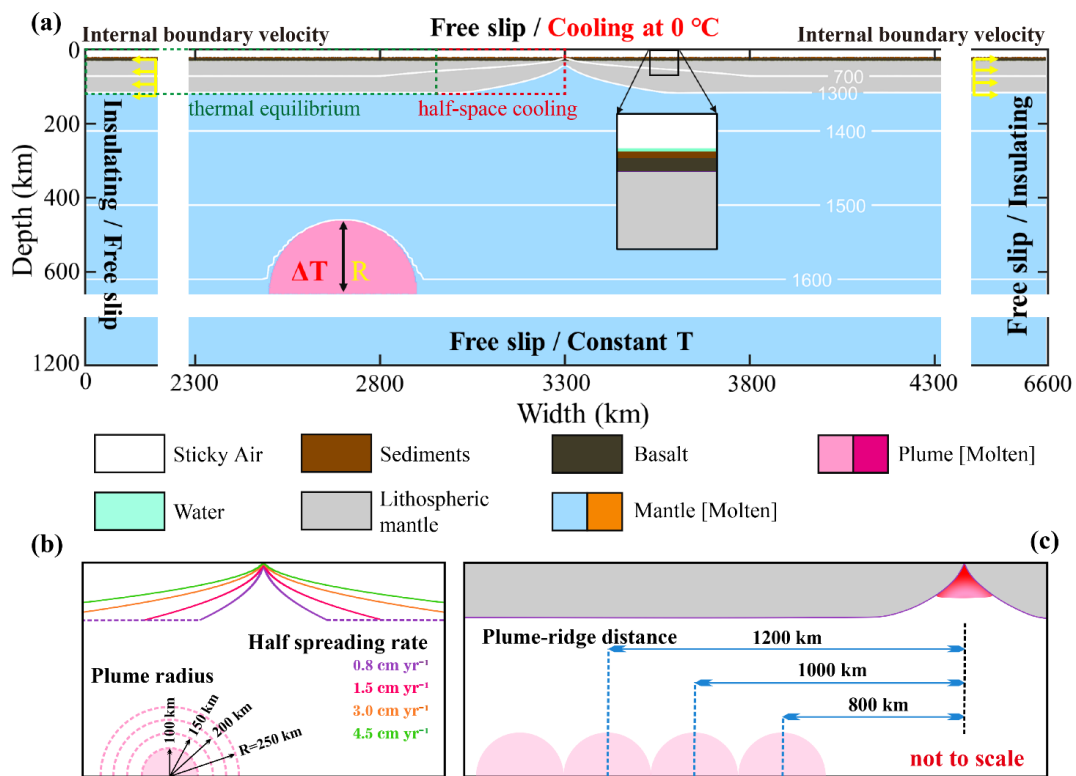


Figure 2. Model setup. **(a)** Initial composition and boundary conditions. The oceanic plate consists of half-space cooling part and the thermal equilibrium part. A 50-Myrs-old mid-ocean ridge sets in the middle of the model based on half-space cooling temperature structure. A thermal and chemical anormal mantle plume locates at 660 km. Different colors indicate the initial rock types and corresponding newly formed molten rock types. Yellow arrows are the half-spreading rates imposed internal in the lithosphere (i.e., from 20 km to 120 km in depth) to simulate ridge spreading. **(b)**

Initial tested ridge and plume configurations. (c) Initial tested plume-ridge distances.